

# Towards more accurate and reliable predictions for nuclear applications

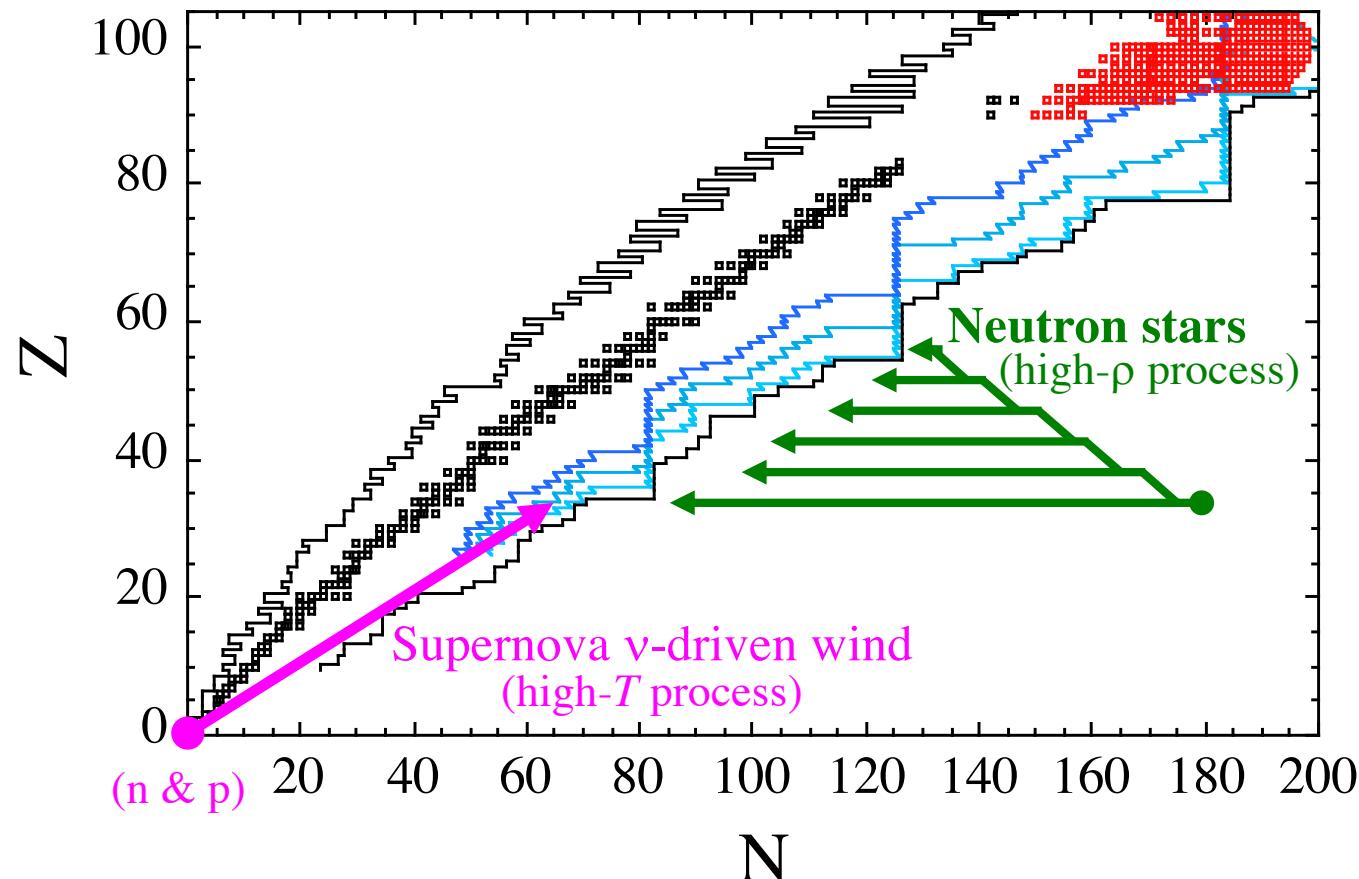
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# The r-process nucleosynthesis responsible for half the elements heavier than iron in the Universe

one of the still unsolved puzzles in nuclear astrophysics  
... the r-process site remains to be confirmed ...

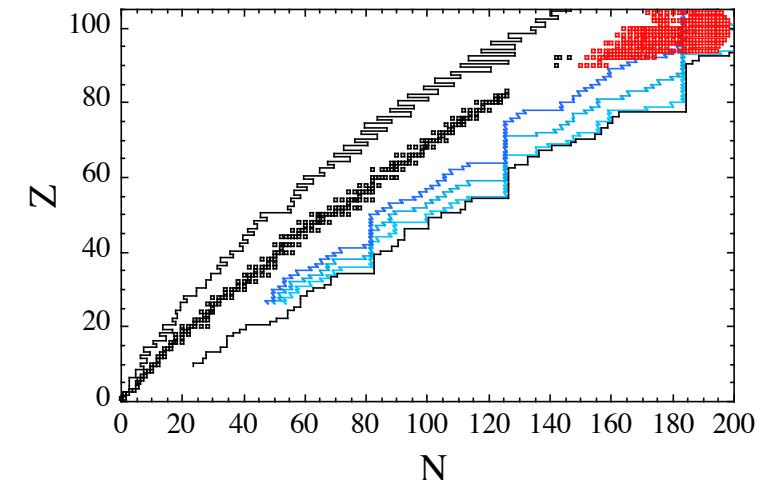
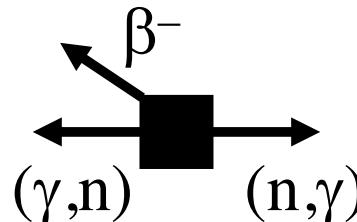


# Nuclear physics input for r-process applications

(n, $\gamma$ ) – ( $\gamma$ ,n) –  $\beta^-$  competition & Fission recycling

Main needs

- $\beta^-$ -decay
- (n, $\gamma$ ) and ( $\gamma$ ,n) rates
- Fission (nif, sf,  $\beta$ df) rates
- Fission Fragments Distributions



Nucleosynthesis requires RATES for some 5000 nuclei !

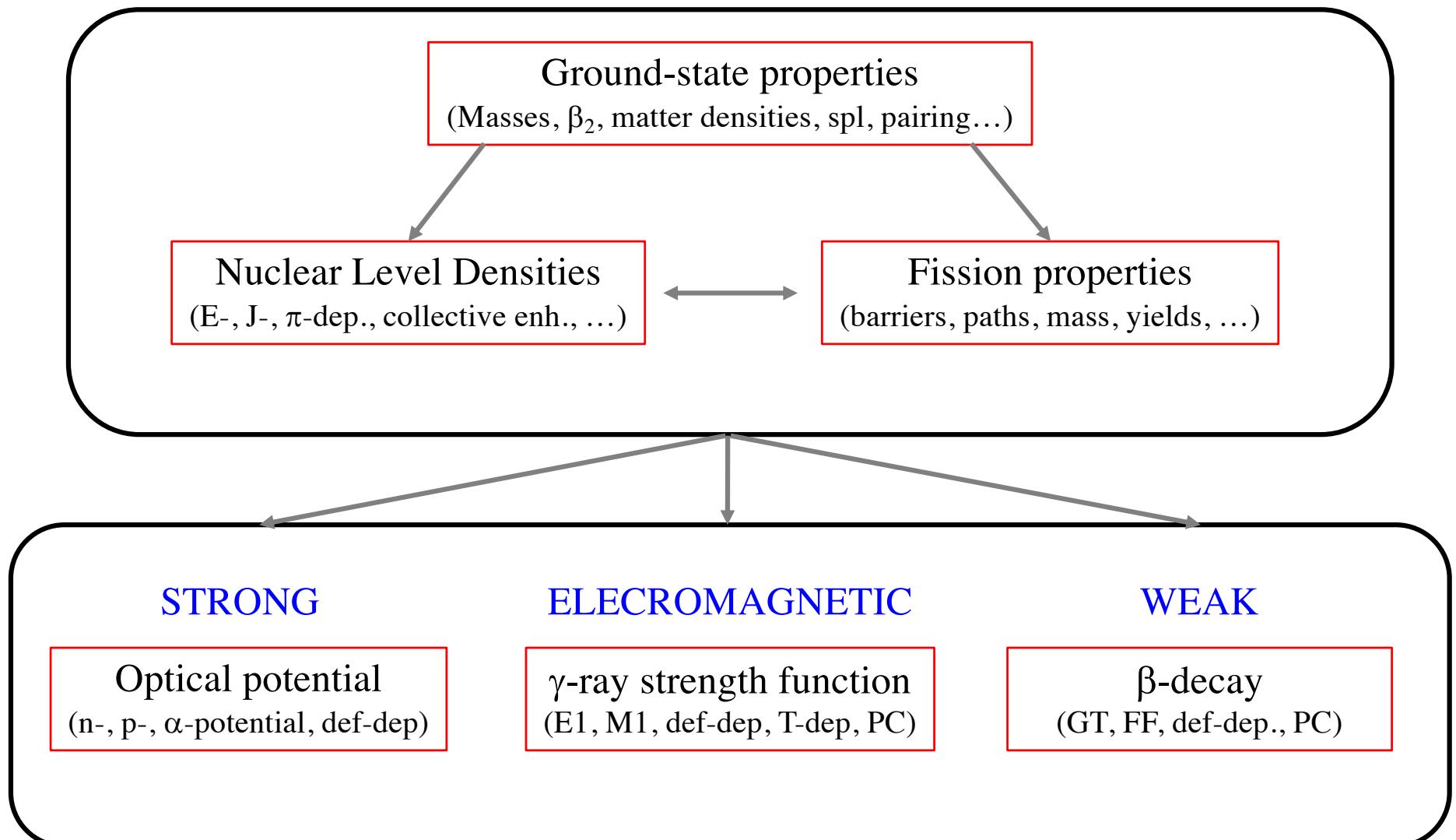
(and not only masses or  $\beta$ -decay along the oversimplified so-called “r-process path”)

→ simulations rely almost entirely on theoretical predictions

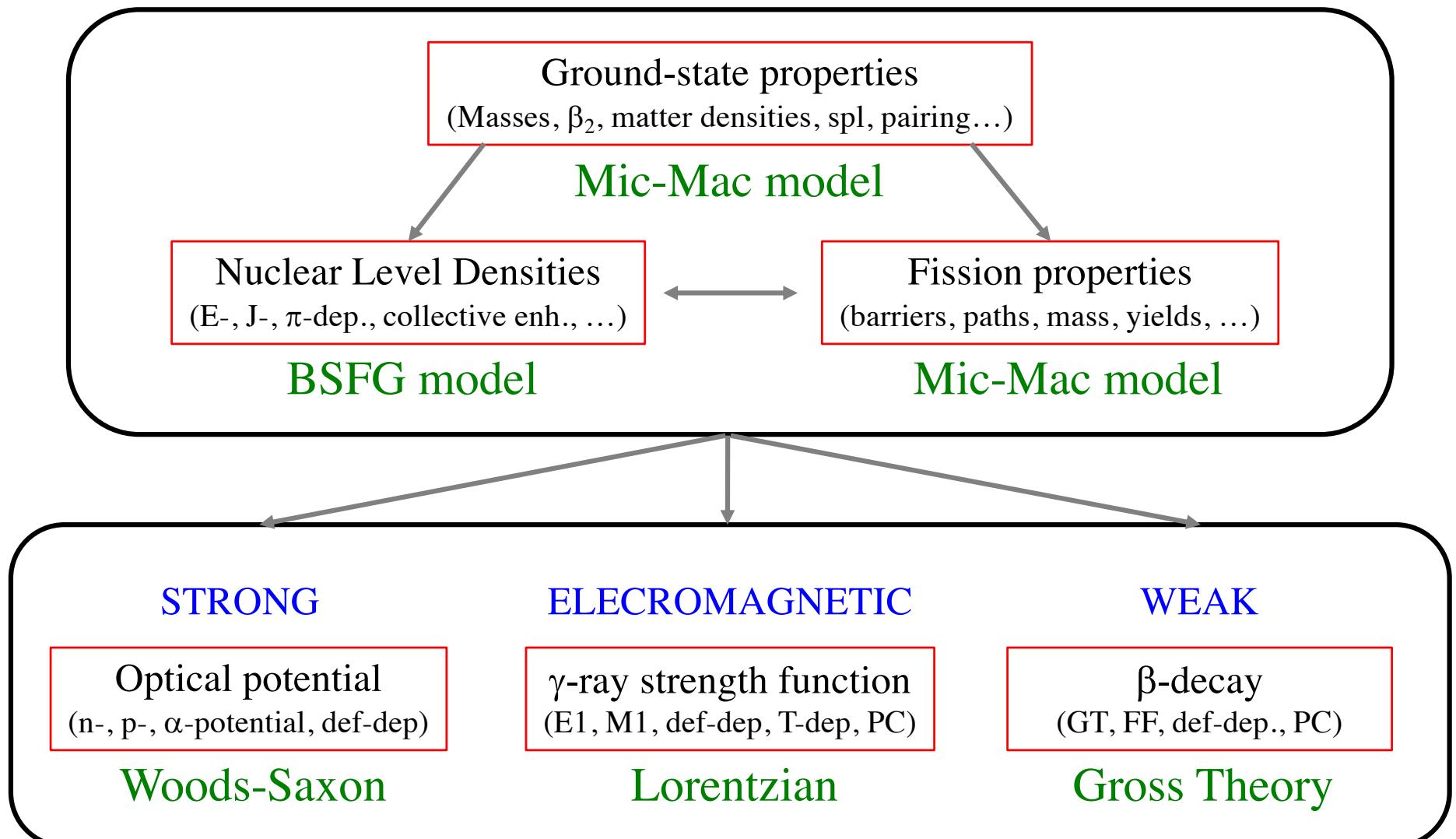
In turn, theoretical models are tuned on available experimental data

Ongoing progress on both theoretical and experimental sides

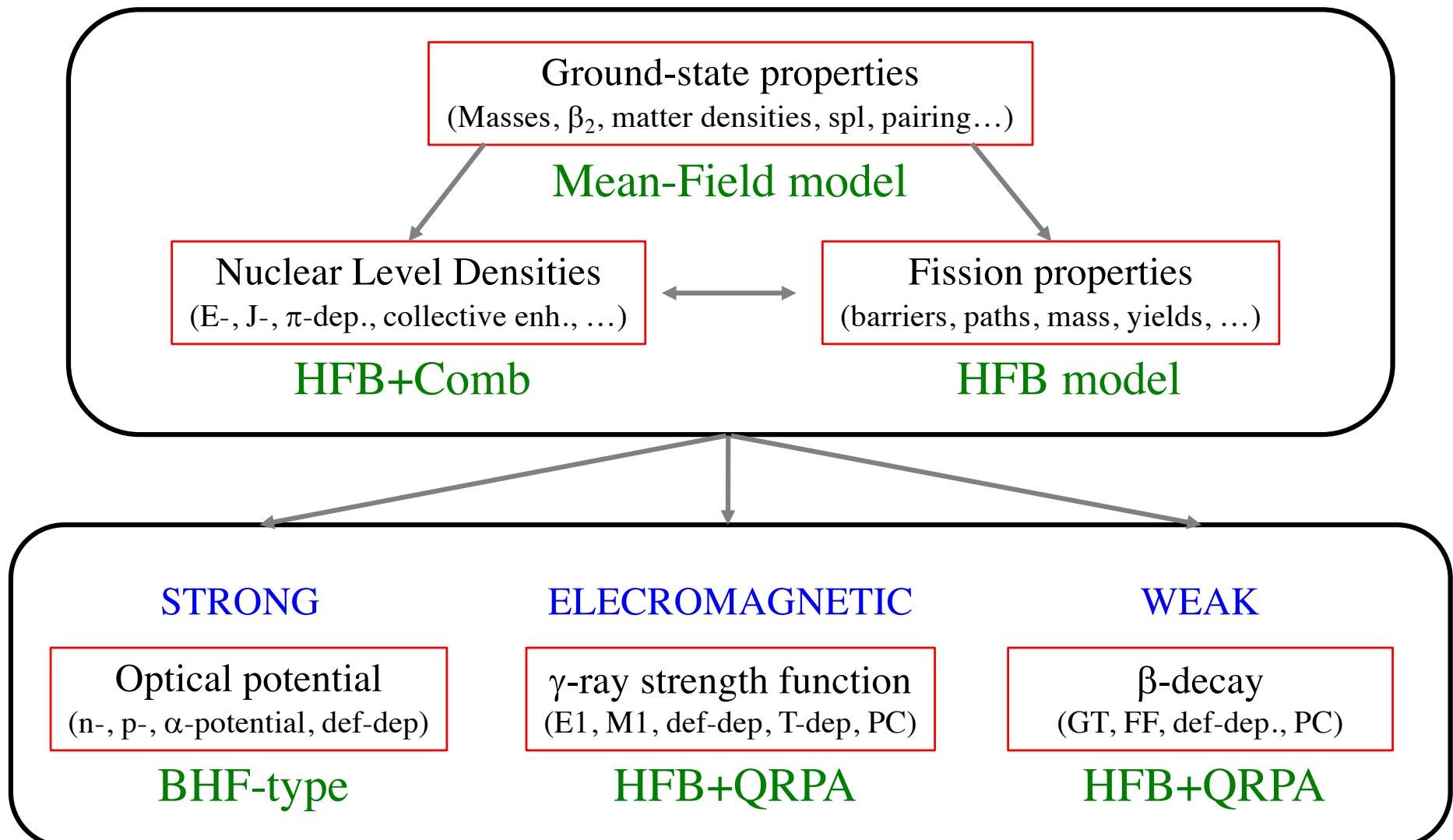
# Nuclear inputs to nuclear reaction & decay calculations



# Nuclear inputs to nuclear reaction & decay calculations

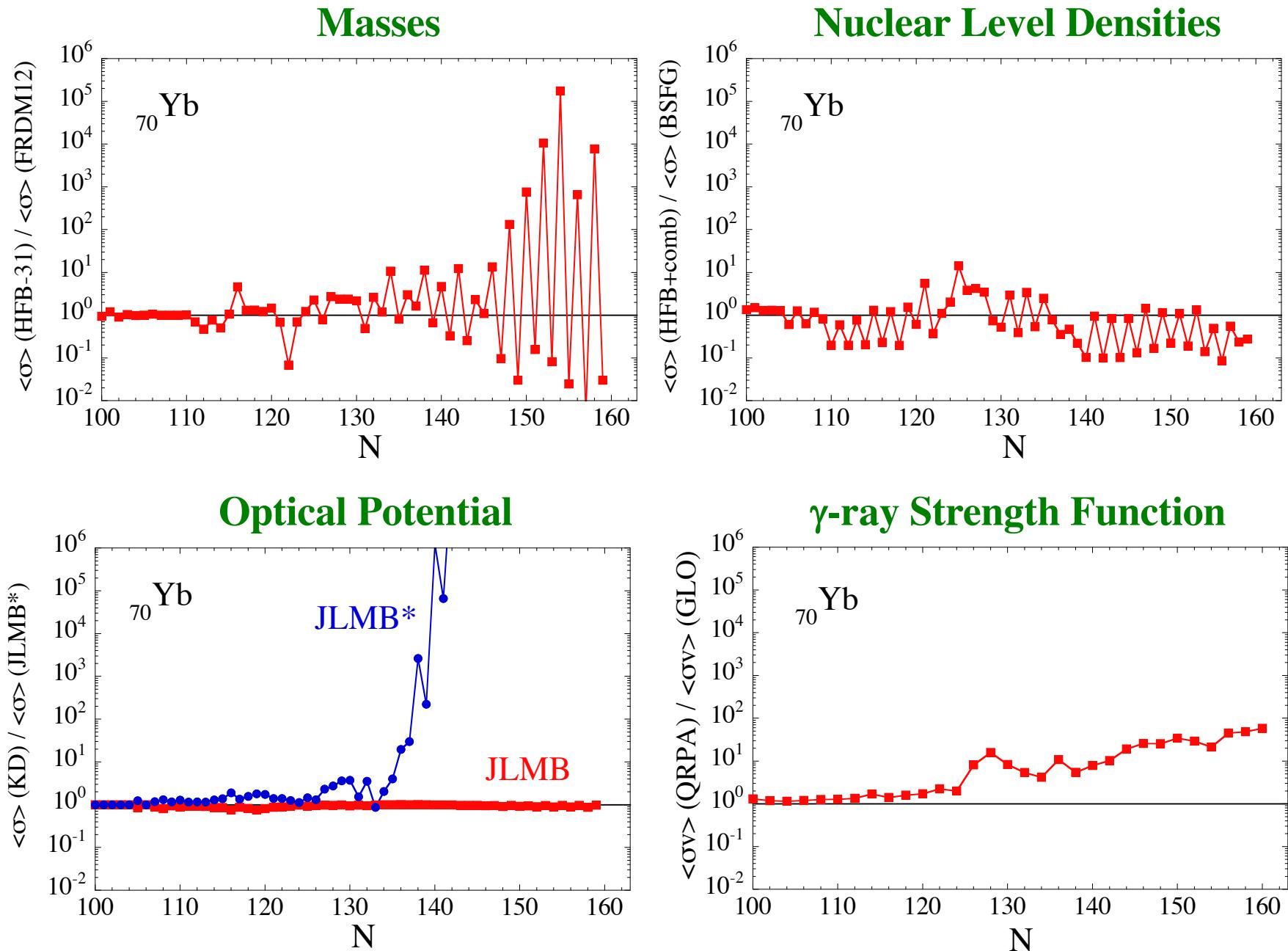


# Nuclear inputs to nuclear reaction & decay calculations



“Microscopic” approach is a necessary but not a sufficient condition !  
”(Semi-)Microscopic” models must be competitive in reproducing exp. data !

# Impact of the various ingredients on the radiative neutron capture ( $E_n \sim 100\text{keV}$ )



# Mean Field mass models

$$E = E_{MF} - E_{coll} - E_W - E_{b\infty}$$

$E_{MF}$ : HFB or HF-BCS (or HB) main Mean-Field contribution

$E_{coll}$ : Quadrupole Correlation corrections to restore broken symmetries and include configuration mixing

$E_W$ : *Wigner* correction contributes significantly only for nuclei along the  $Z \sim N$  line (and in some cases for light nuclei)

$E_{b\infty}$ : Correction for infinite basis

Skyrme-HFB

Gogny-HFB

Relativistic MF

rms  $\sim$  0.5-0.8MeV

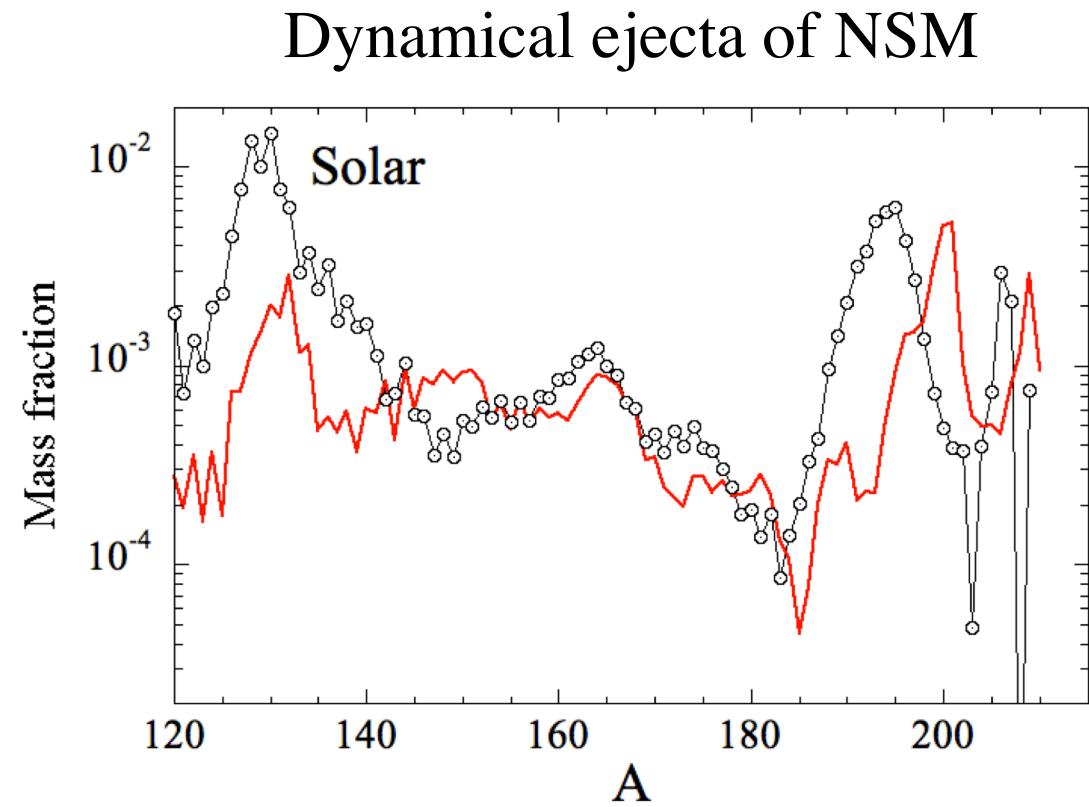
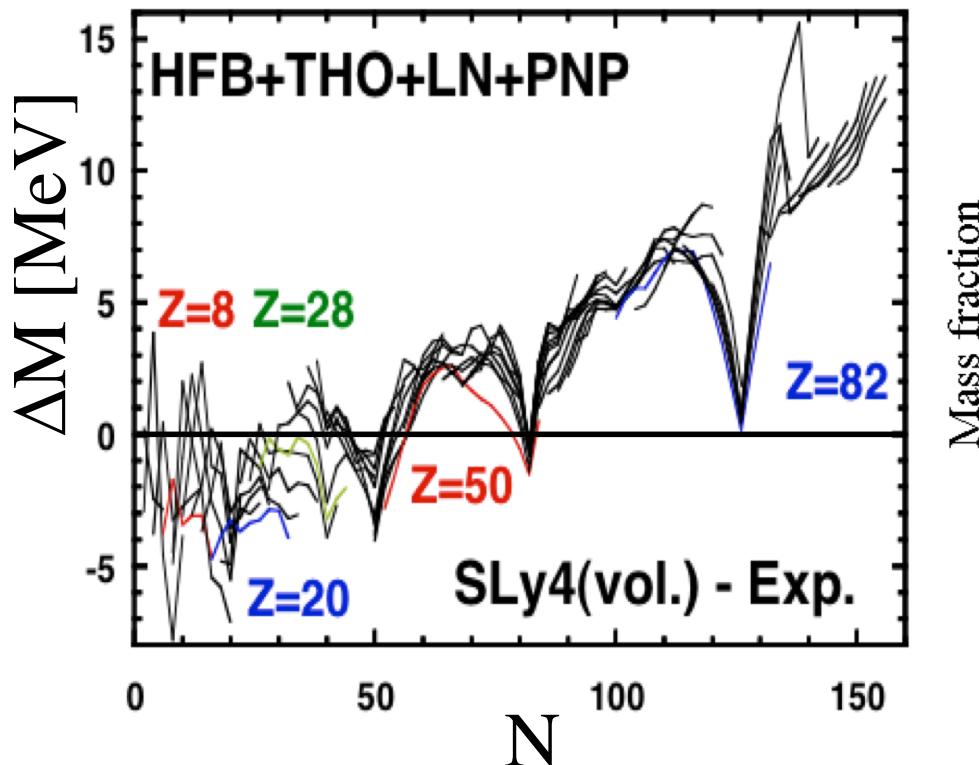
rms  $\sim$  0.8MeV

rms  $>$  1.1MeV

# Nuclear masses and their impact on the r-abundance distribution

Nuclear masses driving the  $(n,\gamma) \longleftrightarrow (\gamma,n)$  competition

SLy4 :  $\sigma_{\text{rms}}=5.1\text{MeV}$  on 570 e-e nuclei

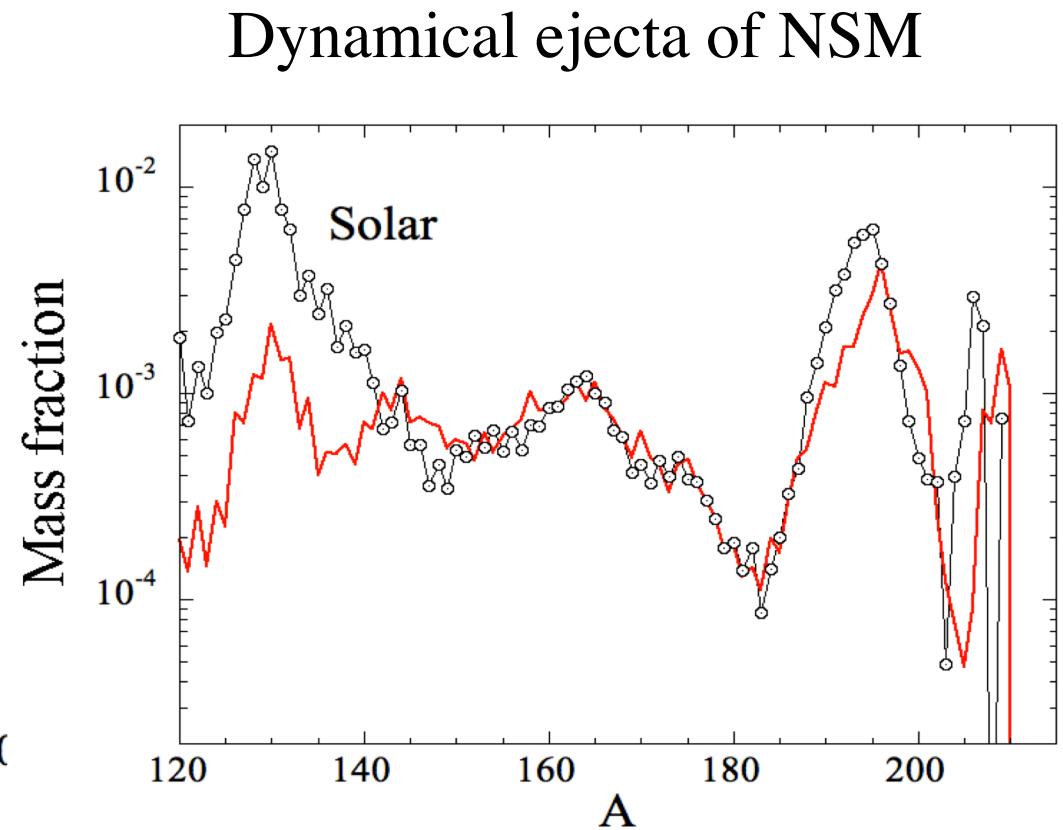
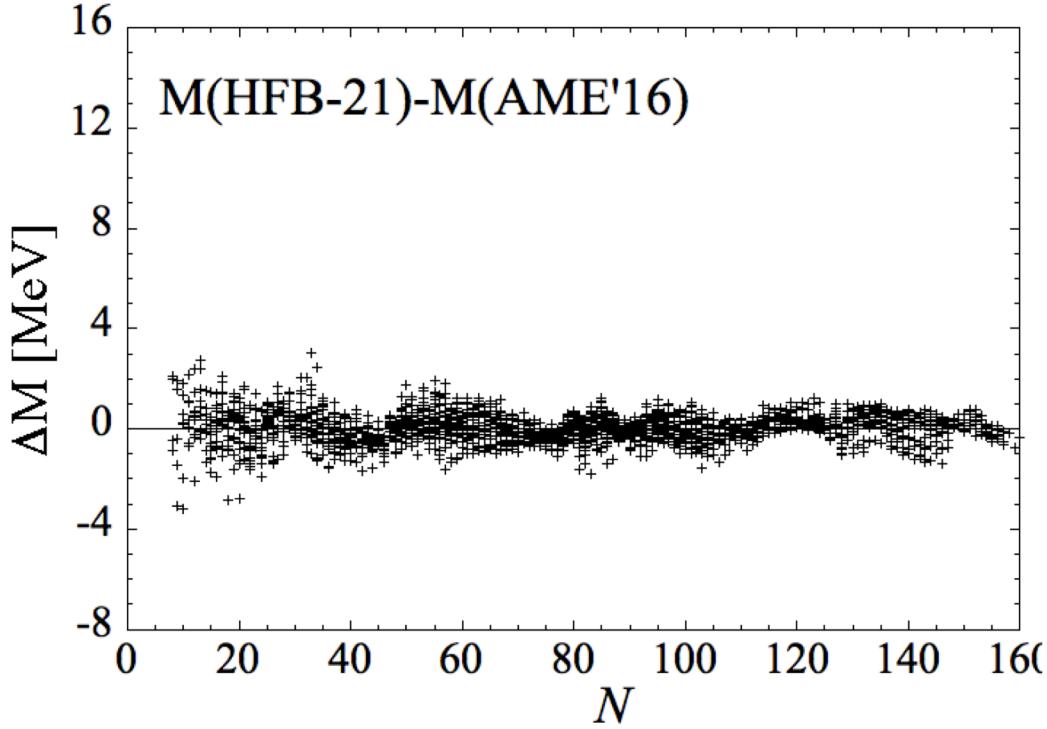


Stoitsov et al. 2003

# Nuclear masses and their impact on the r-abundance distribution

Nuclear masses driving the  $(n,\gamma) \longleftrightarrow (\gamma,n)$  competition

HFB-21 :  $\sigma_{\text{rms}}=0.59\text{MeV}$  on 2408 nuclei



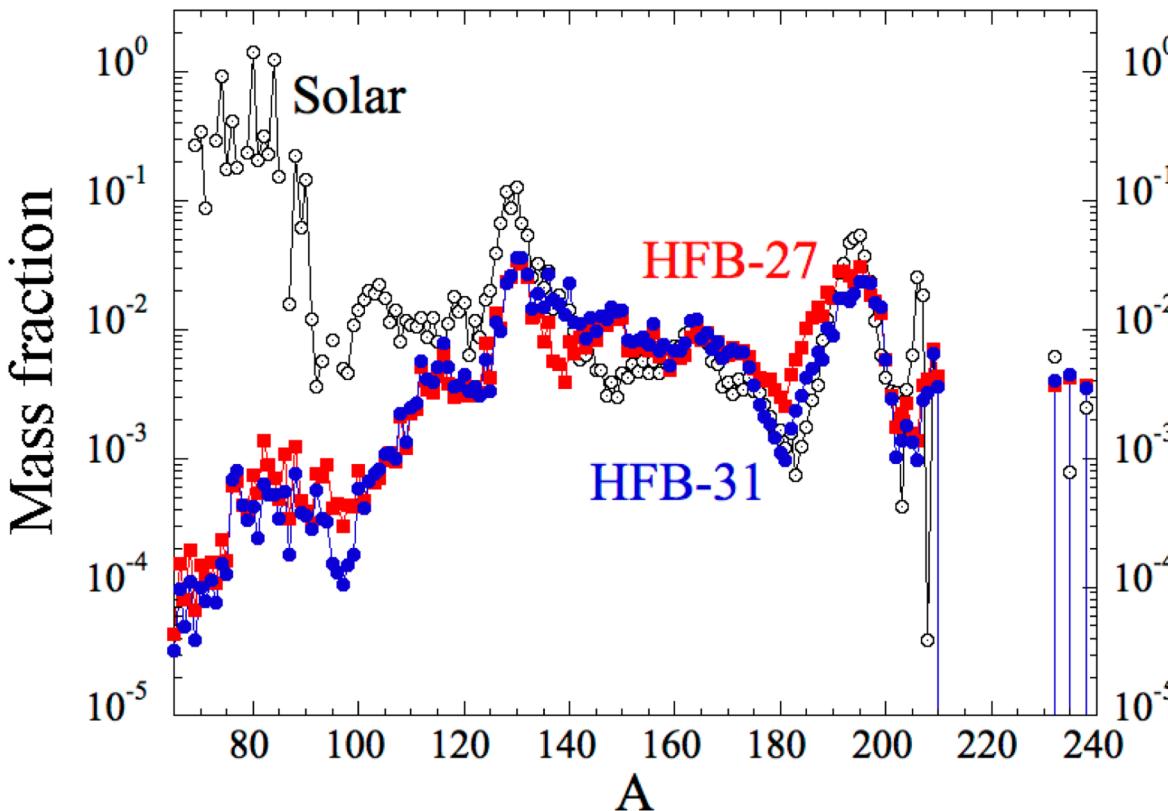
Good mass fit ( $\text{rms} < 0.8\text{MeV}$ ) is a necessary condition, but not a sufficient one

# Nuclear masses and their impact on the r-abundance distribution

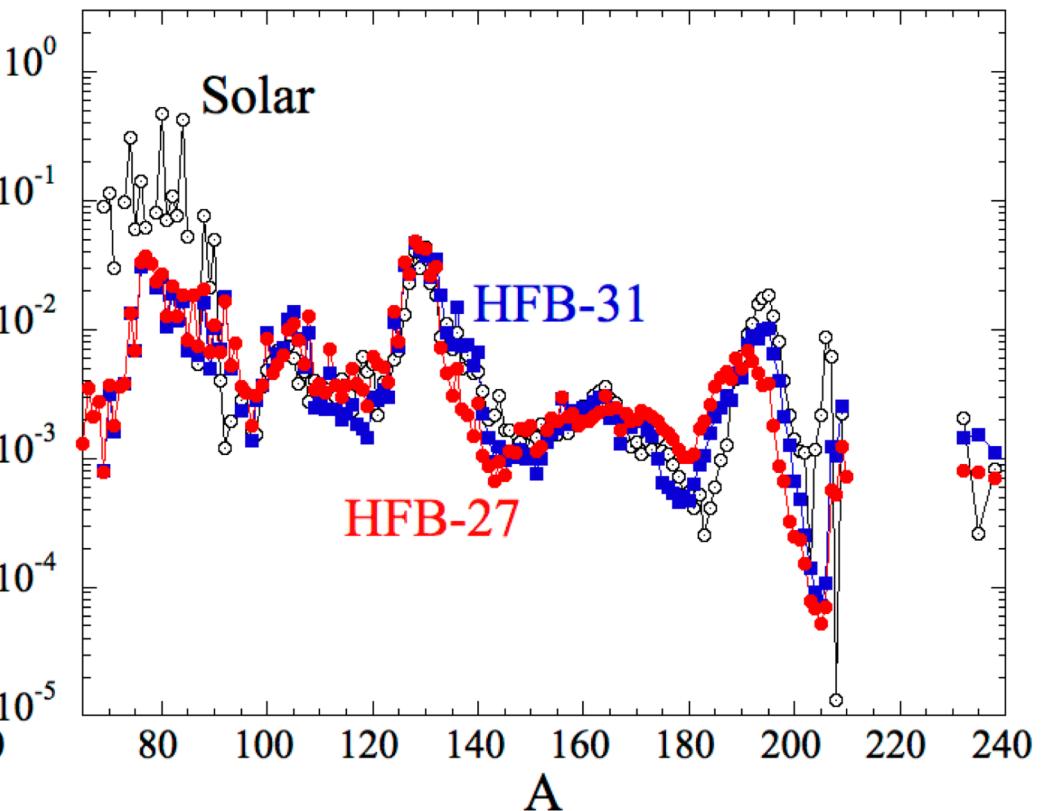
Nuclear masses driving the  $(n,\gamma) \longleftrightarrow (\gamma,n)$  competition

2 HFB mass models: HFB-27: Standard Skyrme  $\sigma_{\text{rms}}=0.50\text{MeV}$   
HFB-31: Generalized Skyrme  $\sigma_{\text{rms}}=0.56\text{MeV}$

Dynamical ejecta

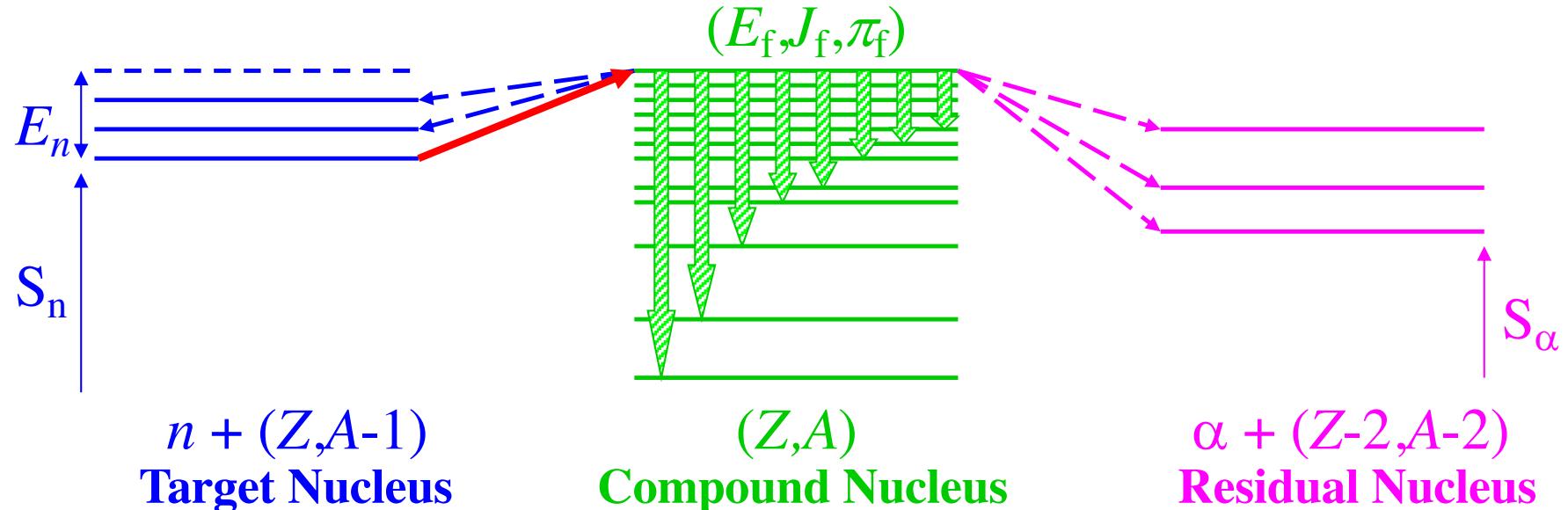


Wind ejecta



→ Need further “microscopic” calculations beyond mean-field

# Hauser-Feshbach model for radiative neutron capture reactions

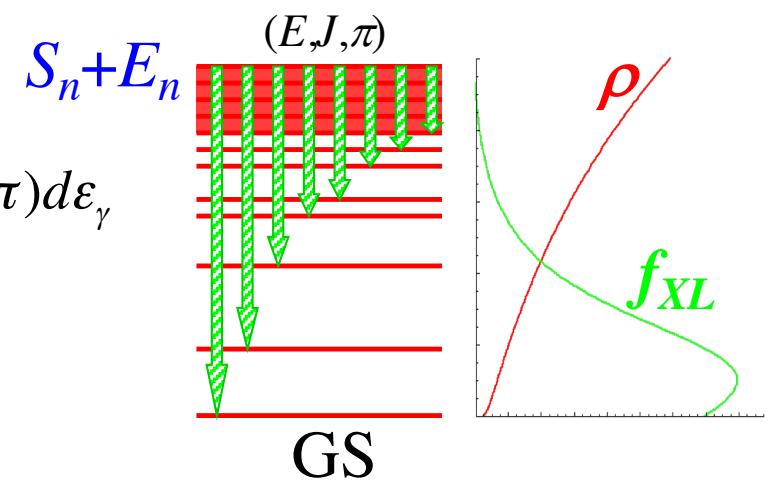


$$\sigma_{(n,\gamma)} \propto \sum_{J,\pi} \frac{T_n(J^\pi)T_\gamma(J^\pi)}{T_n(J^\pi) + T_\gamma(J^\pi)} \approx \sum_{J,\pi} T_\gamma(J^\pi) \quad \text{since } T_n(J^\pi) \gg T_\gamma(J^\pi): E_n \sim \text{keV}$$



$$T_\gamma = \sum_{J^\pi XL} \int_0^{S_n + E_n} 2\pi \varepsilon_\gamma^{2L+1} f_{XL}(\varepsilon_\gamma) \rho(S_n + E_n - \varepsilon_\gamma, J, \pi) d\varepsilon_\gamma$$

Nuclear astrophysics apps require  
NLDs & GSF for  $\sim 8000$  nuclei



# Mean Field + QRPA $\gamma$ -ray strength function

*Large-scale E1 Mean-Field + QRPA calculations*

Skyrme-HFB + QRPA

Gogny-HFB + QRPA

RMF +QRPA

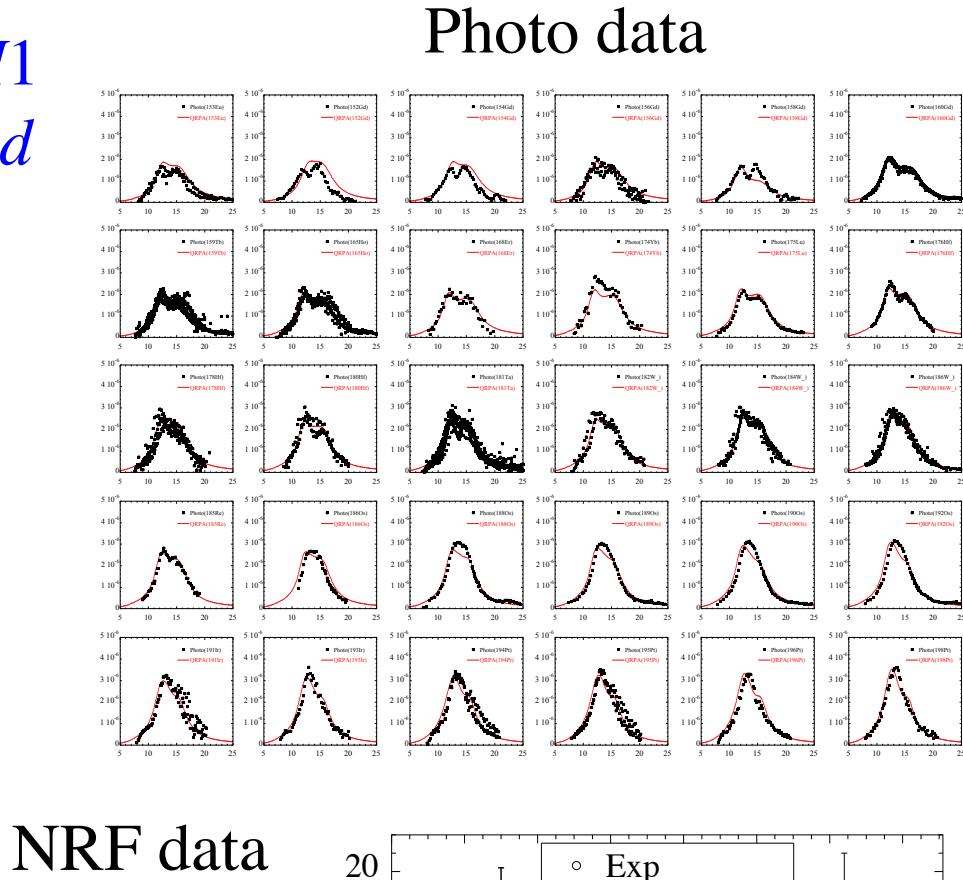
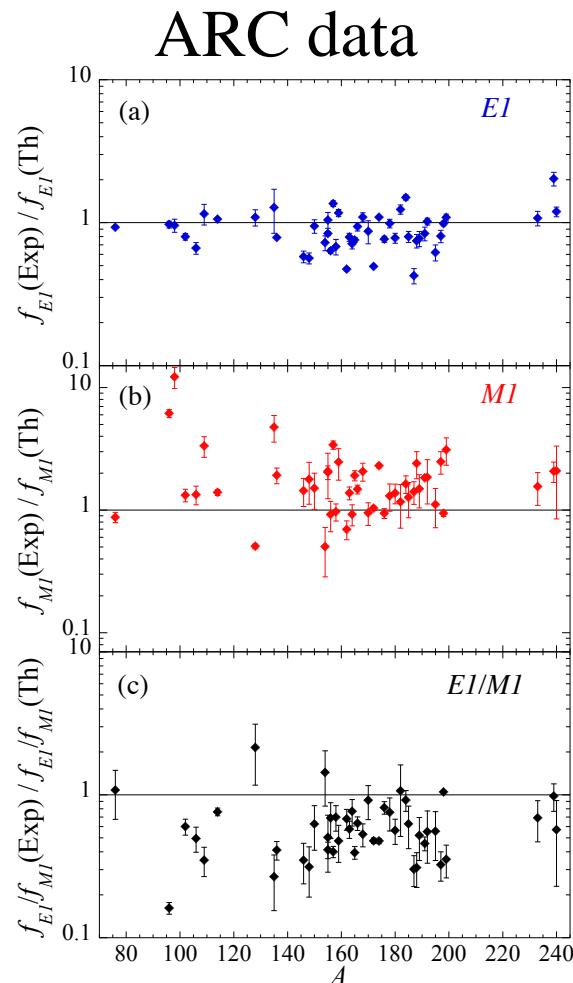
QRPA calculations can accurately reproduce experimental data,  
provided *empirical corrections* are made, *i.e.*

- Empirical Energy shift (beyond 1p-1h excitations and phonon couplings)
- Empirical damping of collective motions → broadening
- Empirical deformation effects for *spherical calculations*

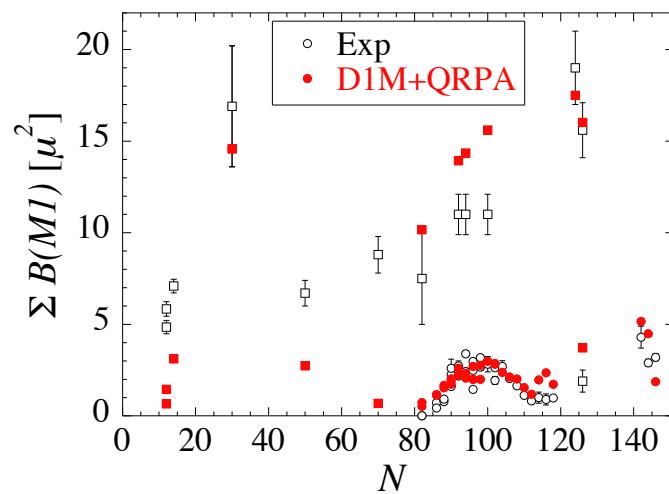
Recent large-scale axially-deformed Gogny-HFB + QRPA  
calculations with D1M interaction  
for  $\sim 2000$  e-e nuclei with  $8 \leq Z \leq 110$   
(interpolation for odd-A & odd-odd nuclei)

# Gogny-HFB + QRPA $E1$ and $M1$ strength functions

QRPA calculations of the  $E1$  &  $M1$   
strengths *empirically renormalized*  
on exp. data



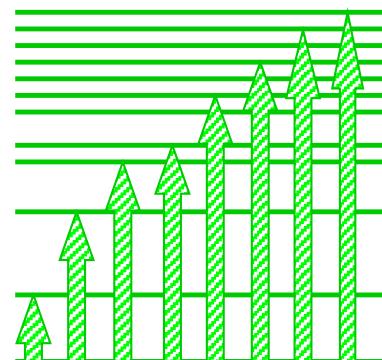
NRF data



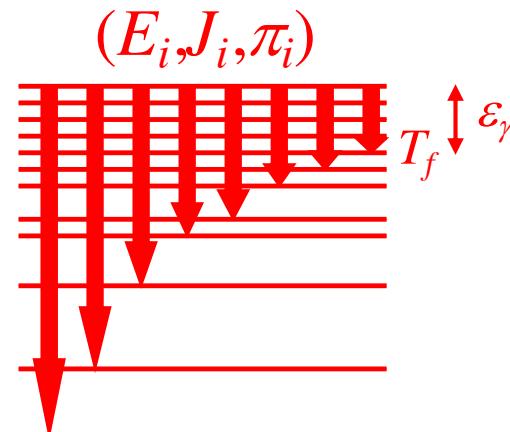
# Presence and impact of a possible low-energy $M1$ upbend of the de-excitation strength function

Violation of the Brink hypothesis

$$\begin{array}{c} \text{---} \\ \leftarrow \\ \vec{f}_{E1}(\varepsilon_\gamma) \neq \vec{f}_{E1}(\varepsilon_\gamma) \\ \curvearrowright \\ \vec{f}_{E1} = \vec{f}_{E1}(\varepsilon_\gamma, T_f) \end{array}$$



$(E_0, J_0, \pi_0)$



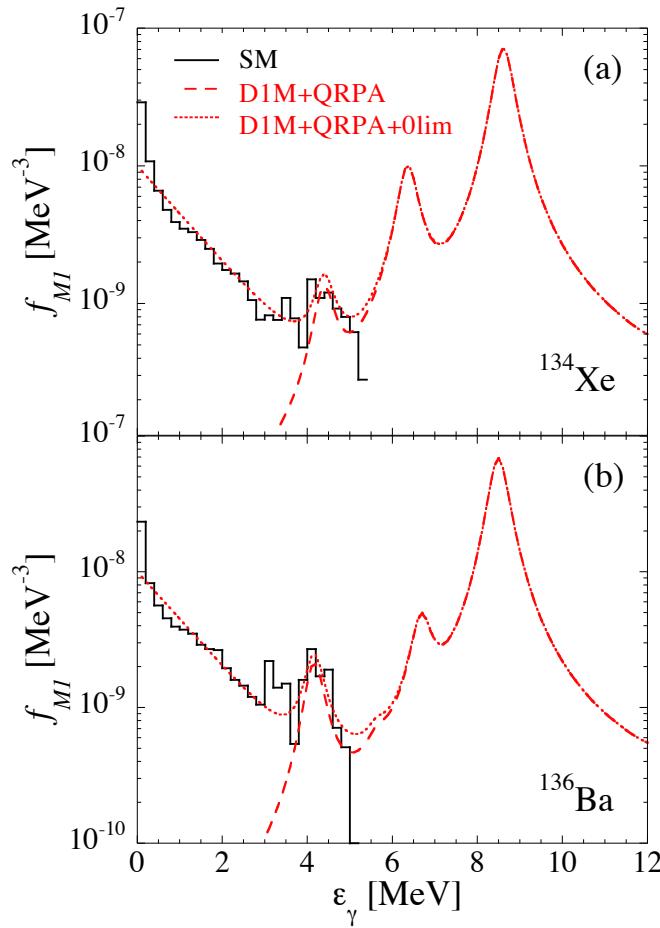
# SM-inspired low-energy correction of the de-excitation strength

$$f_{E1} = f_{E1}^{QRPA} + f_{E1}(\varepsilon_\gamma \rightarrow 0)$$

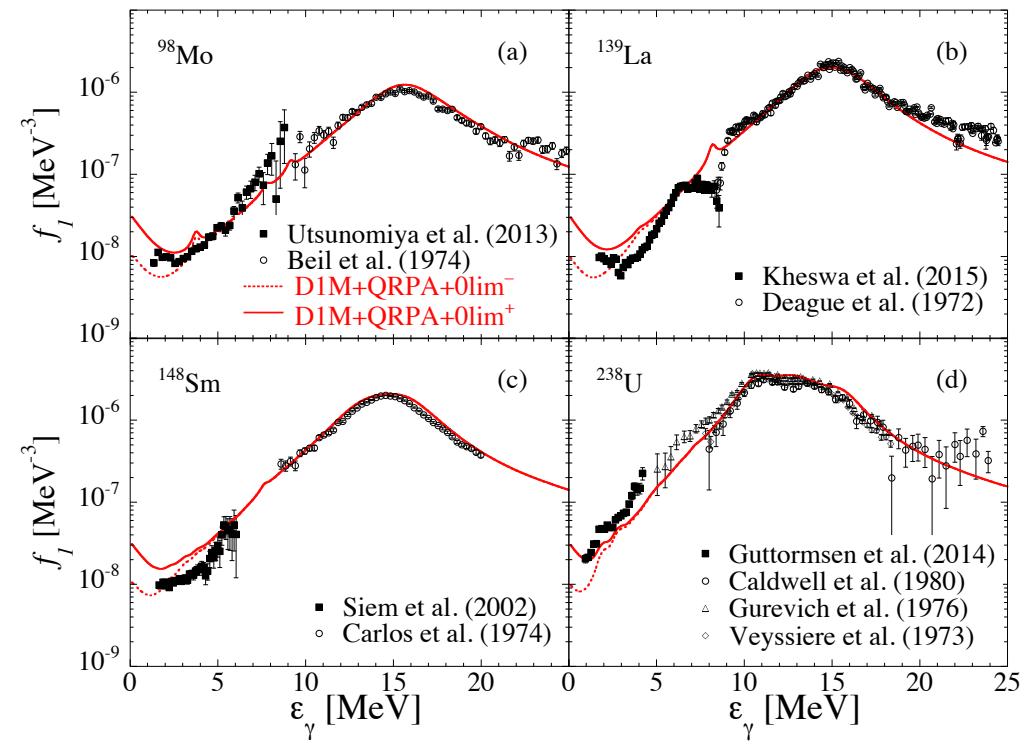
Non-zero limit of the  $E1$  strength at  $\varepsilon_\gamma \rightarrow 0$

$$f_{M1} = f_{M1}^{QRPA} + f_{M1}(\varepsilon_\gamma \rightarrow 0)$$

Upbend of the  $M1$  strength at  $\varepsilon_\gamma \rightarrow 0$



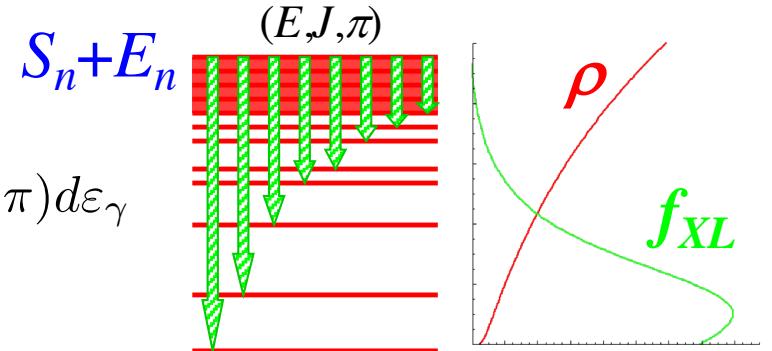
Oslo data



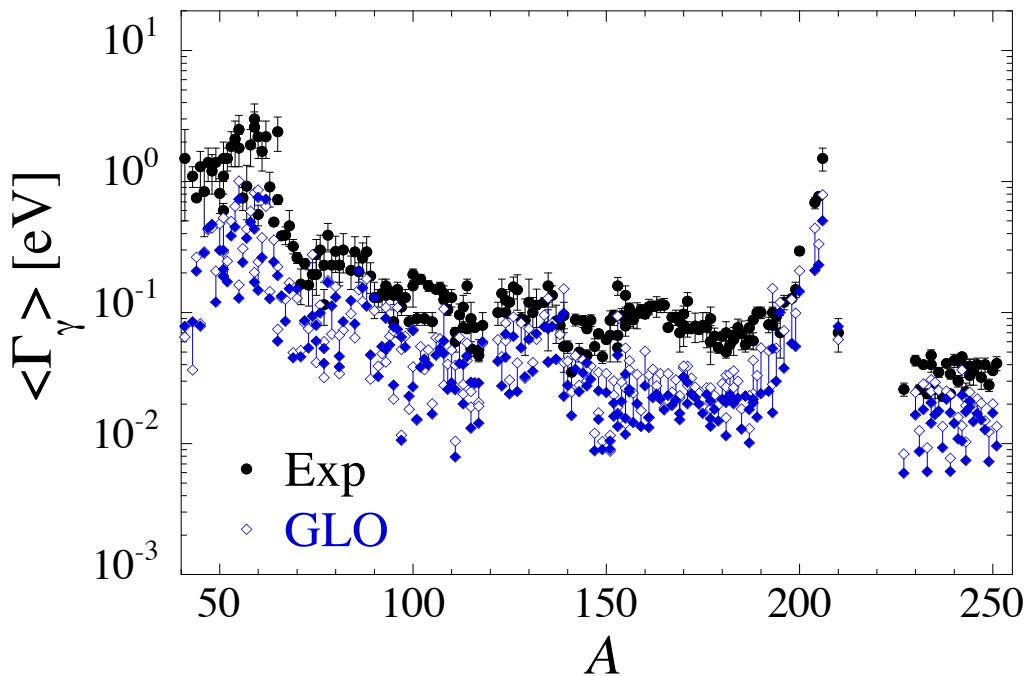
See also Litvinova talk (this workshop)

# The long-standing problem of the average radiative width $\langle \Gamma_\gamma \rangle$

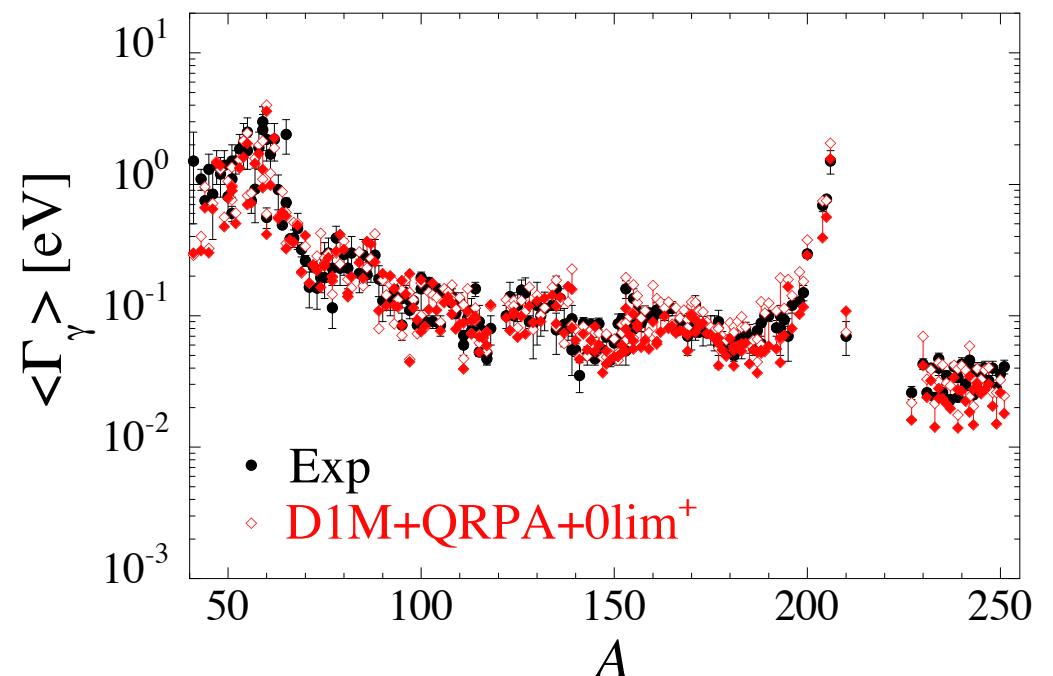
$$\langle \Gamma_\gamma \rangle = \frac{D_0}{2\pi} \sum_{X,L,J,\pi} \int_0^{S_n + E_n} T_{XL}(\varepsilon_\gamma) \times \rho(S_n + E_n - \varepsilon_\gamma, J, \pi) d\varepsilon_\gamma$$



Standard GLO

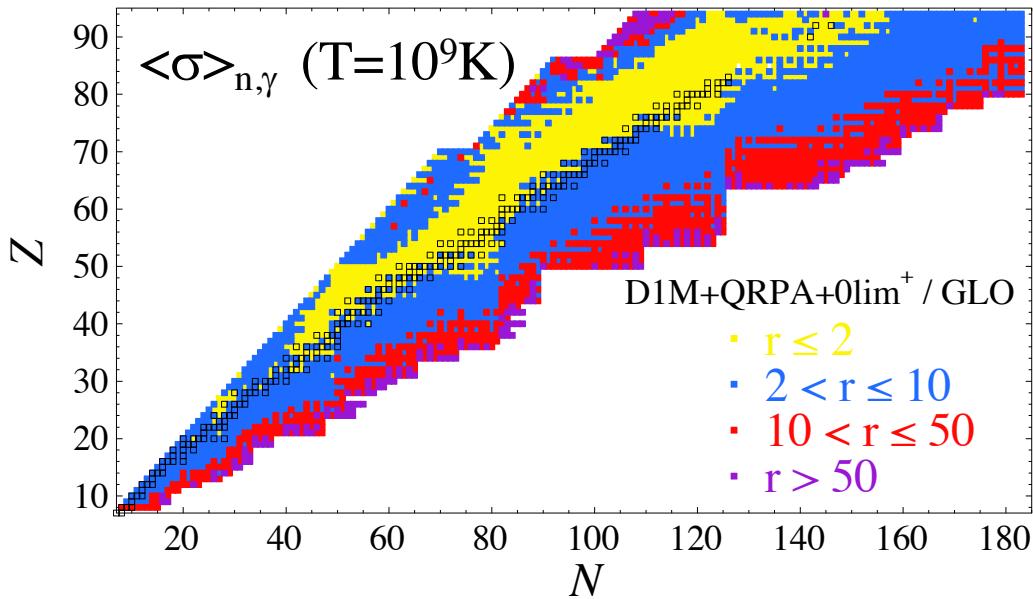


New D1M+QRPA+0lim



Theoretical error bars correspond to different NLD prescriptions

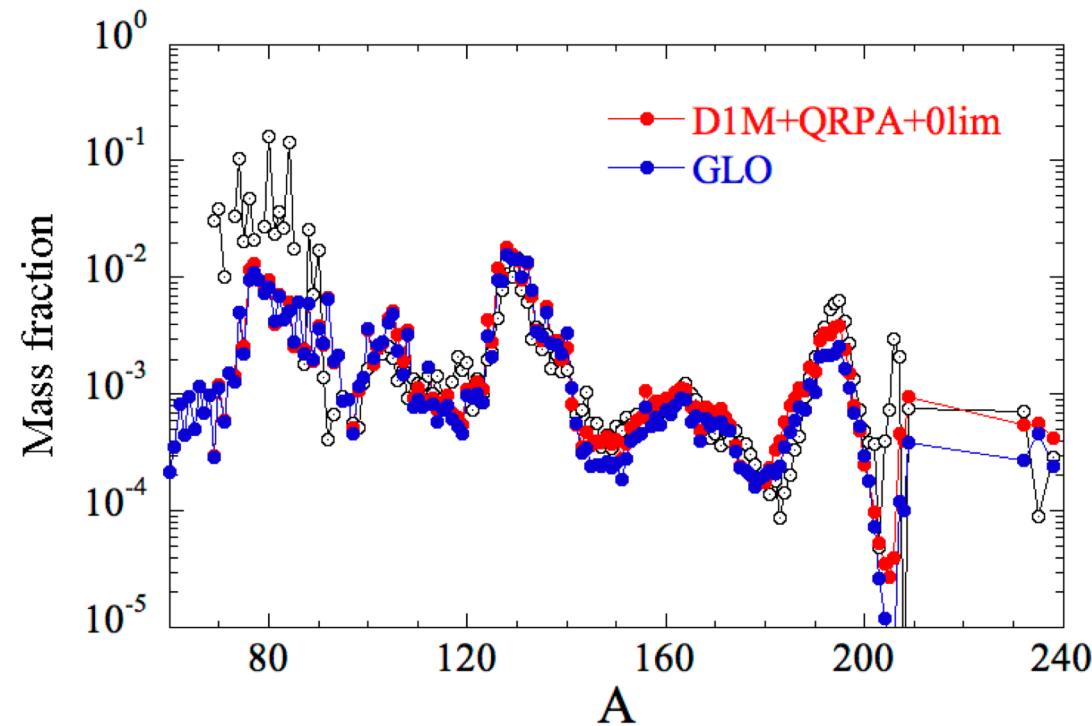
# Impact of the new E1/M1 strength on the radiative n-capture rate of astrophysical interest



Increase of  $\langle\sigma\rangle$  due to low-energy

- $E1$  QRPA low-E strength
- $M1$  scissors & upbend

# Impact of the new E1/M1 strength on the r-process nucleosynthesis in the NSM disk ejecta



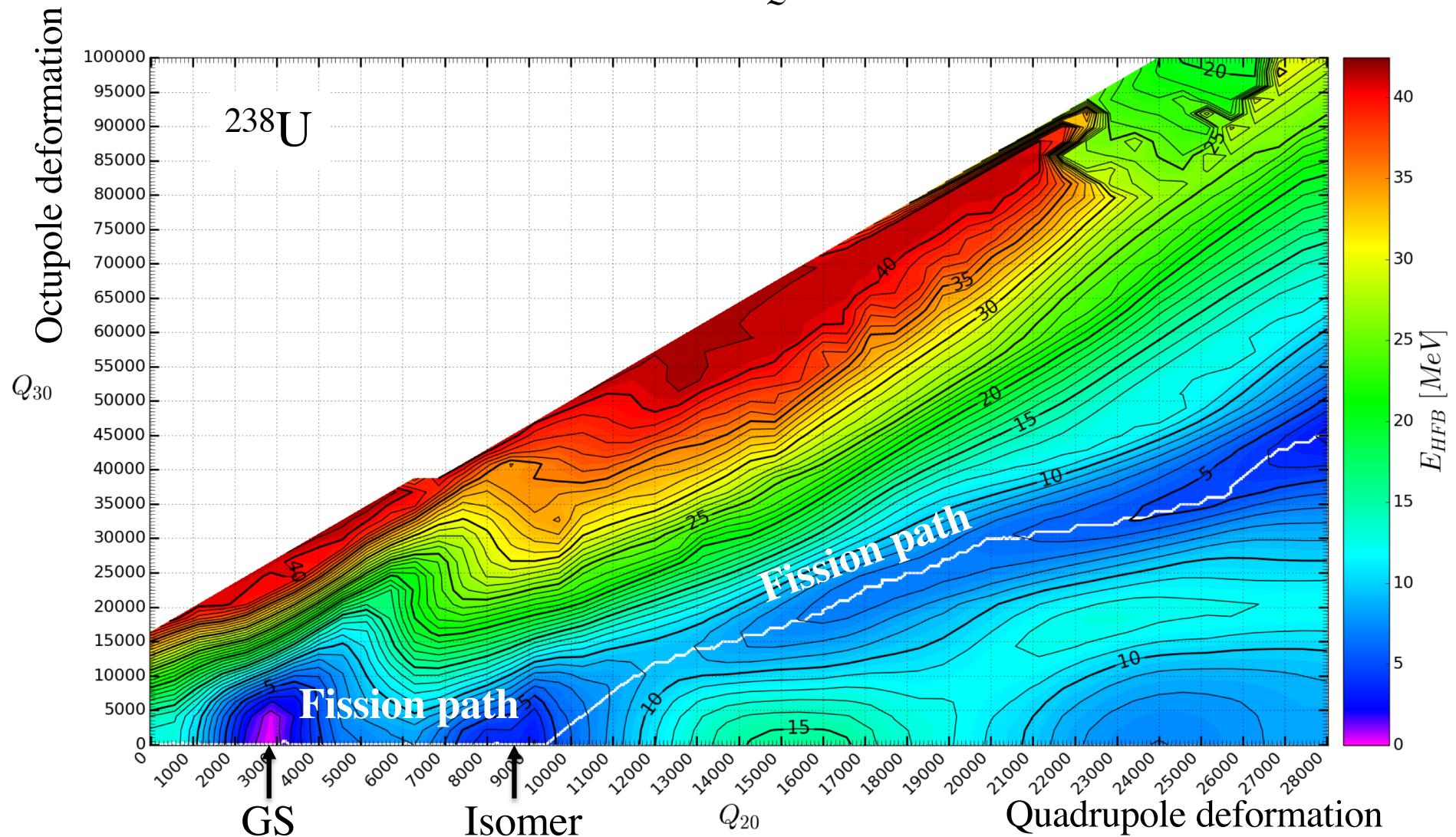
Increase production of

- Lanthanides & 3<sup>rd</sup> peak
- Actinides & chronometers

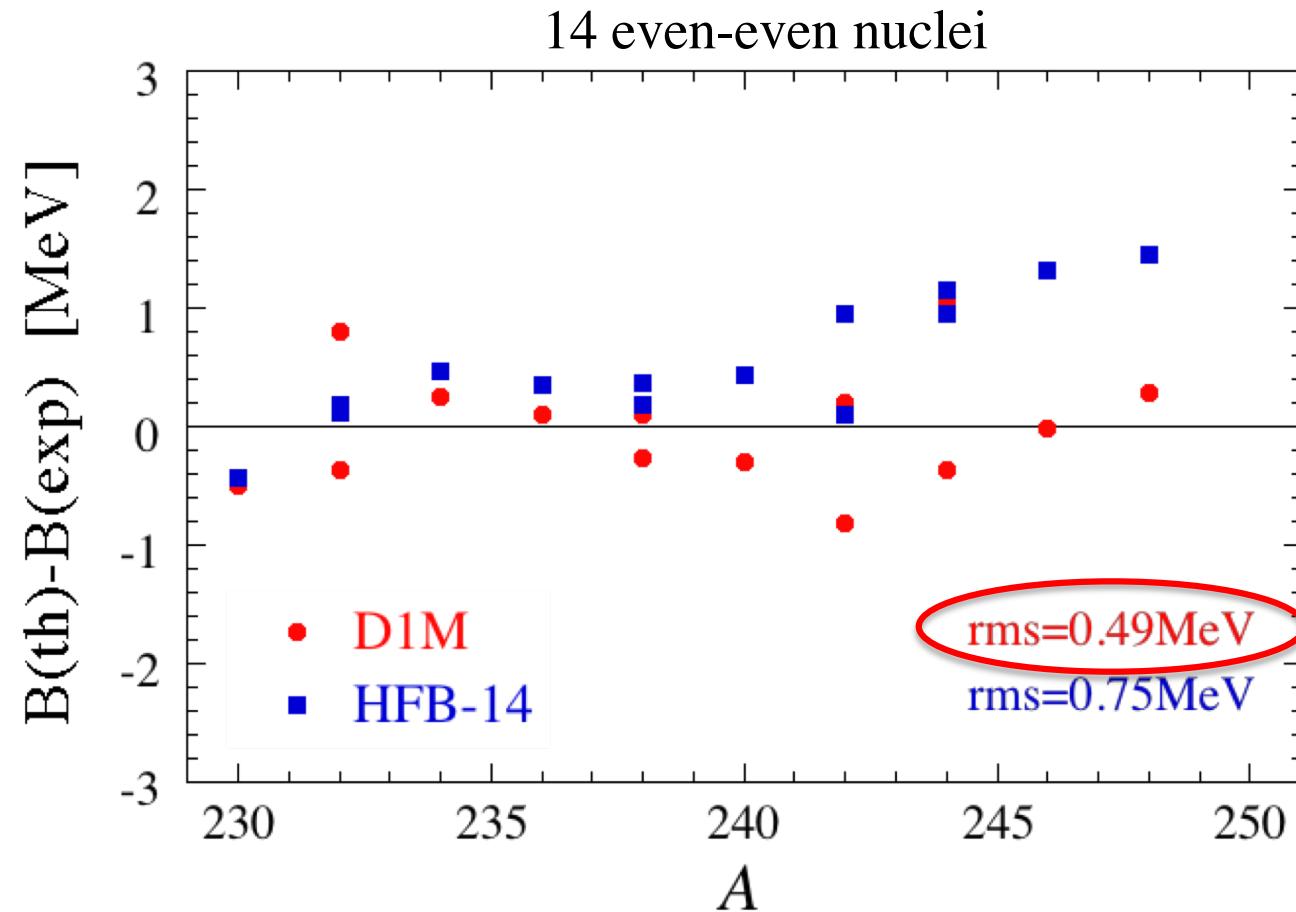
# Calculation of the Gogny-HFB fission paths

Potential energy surface determined with D1M interaction from

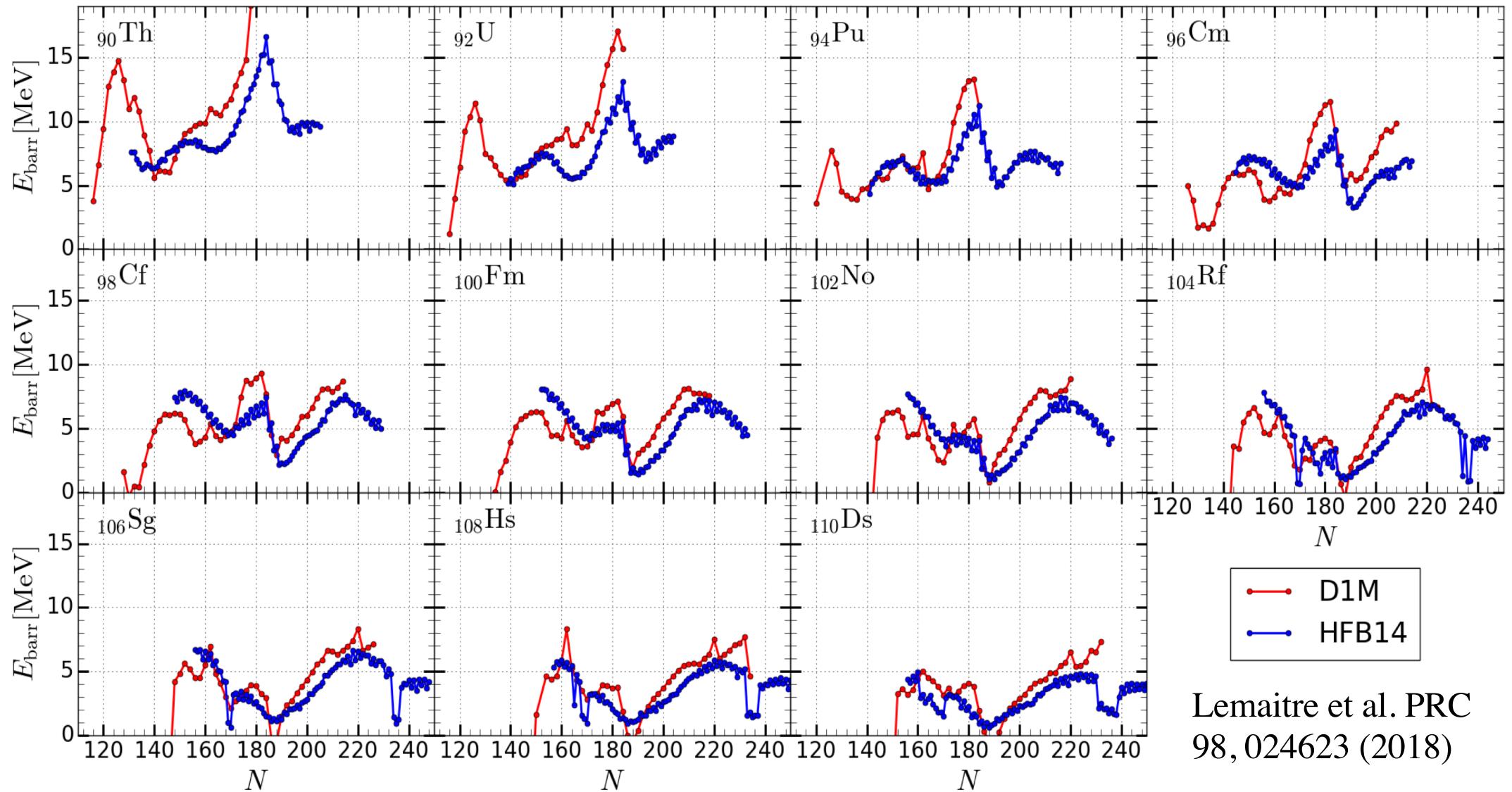
$$E(\beta_2, \beta_3) = E_{HFB} + \Delta E_{Quad} + \Delta E_{triax} + \Delta E_{ZPE} \quad (\text{N. Dubray 2016})$$



# Comparison of the *primary* barrier with empirical barriers (RIPL-3)



# Comparison for 424 fission barriers of e-e $90 \leq Z \leq 110$ nuclei



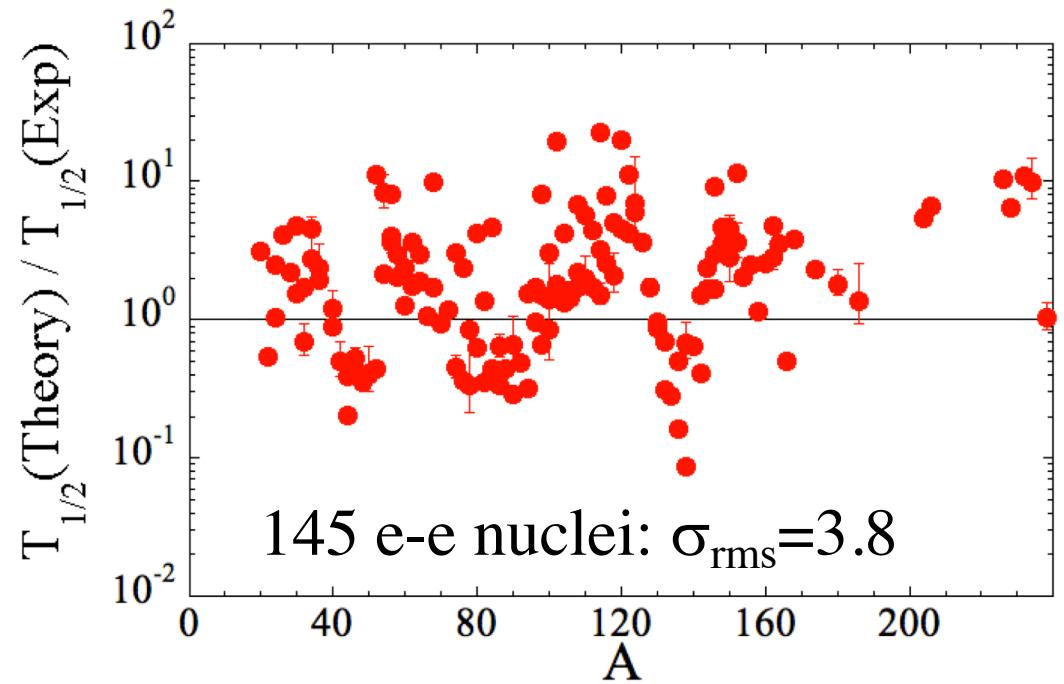
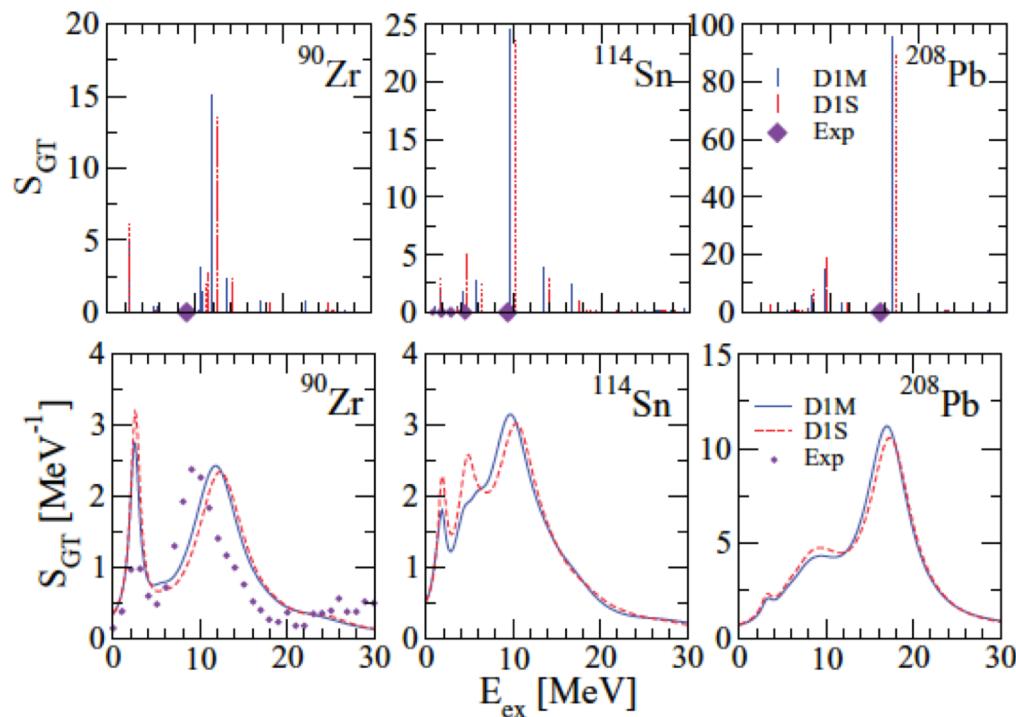
Lemaitre et al. PRC  
98, 024623 (2018)



PES for odd nuclei under progress

## $\beta$ -decay rates

Calculation of the GT  $\beta$ -strength function within the axially deformed HFB+QRPA approach with D1M Gogny force (Martini, Peru, SG 2014)



$$\frac{\ln 2}{T_{1/2}} = \frac{(g_A/g_V)_{\text{eff}}^2}{D} \int_0^{Q_\beta} f_0(Z, A, Q_\beta - E_{\text{ex}}) S_{\text{GT}}(E_{\text{ex}}) dE_{\text{ex}}$$



Inclusion of odd nuclei under progress

# Conclusions: still many open questions

1. How reliable are the present  $\beta$ -decay models (def, FF, odd-A,...) ?
2. How reliable are the present mass models ?
3. How well can we describe fission processes and FFD distributions ?
4. How reliable are NLD and PSF models for n-rich nuclei ?
5. What is the neutron absorption by n-rich nuclei (OMP) ?
6. What is the direct capture (and PE) contribution to the n-capture rates for exotic n-rich nuclei ?
7. What is the impact of the Resolved Resonance Region of n-rich nuclei on n-capture rates ?
8. How can we properly treat nuclear UNCERTAINTIES in *r*-process simulations ?

... Still a lot of nuclear physics questions to answer to ...