

Systematic analysis of light heavy-ion scatterings based on microscopic structure and reaction models

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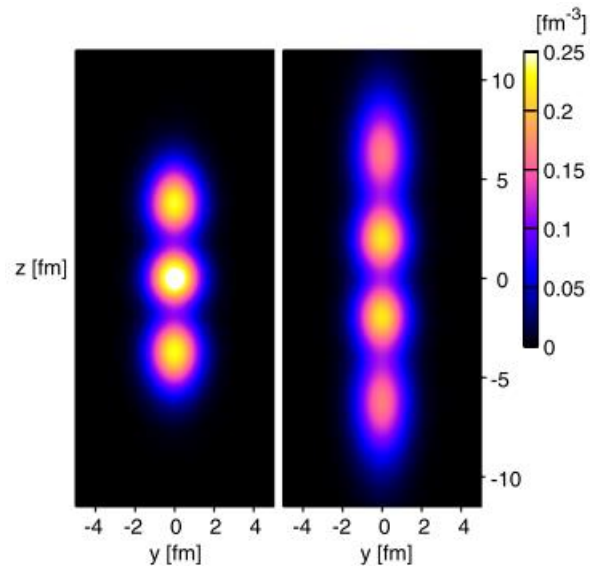
N. Itagaki (YITP, Kyoto U.)

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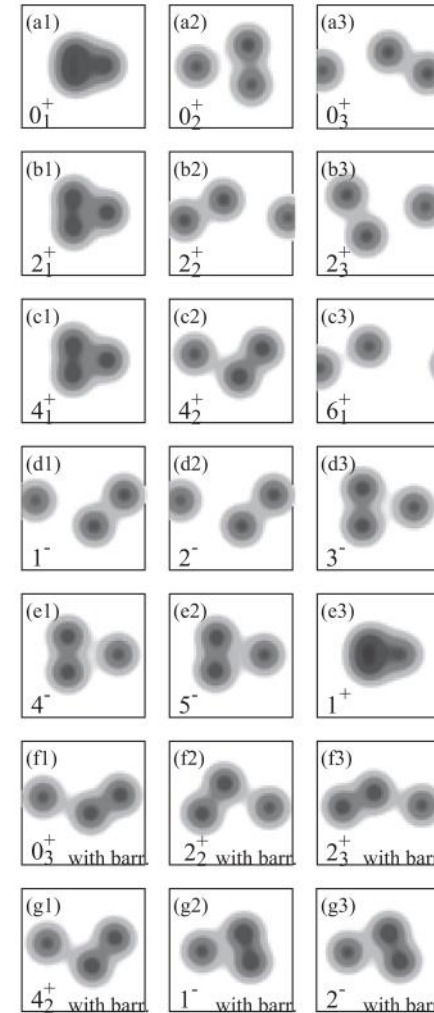
- **Introduction**
Exotic structures in light heavy-ions
- **Formalism**
based on cluster and folding models
- **Result**
Li isotopes & $A = 10$ nuclei
- **Summary**

Introduction ~Exotic structures in stable light heavy-ions~

⇒ The cluster structure is one of hot topics in the light heavy-ions

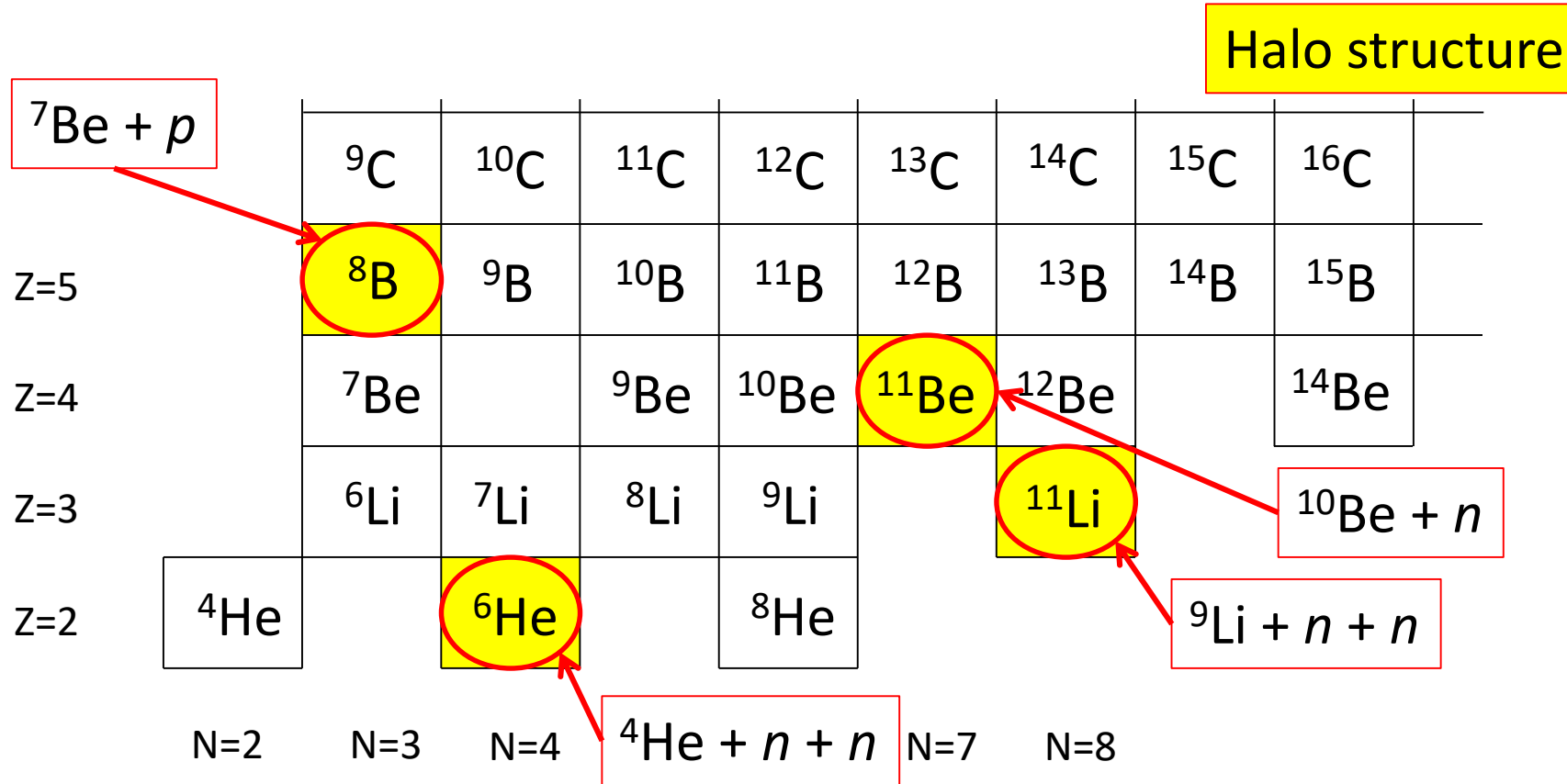


T. Suhara et al., PRL112, 062501 (2014)



Y. Kanada-En'yo, PTP117, 655 (2007)

Introduction ~Exotic structures in unstable light heavy-ions~

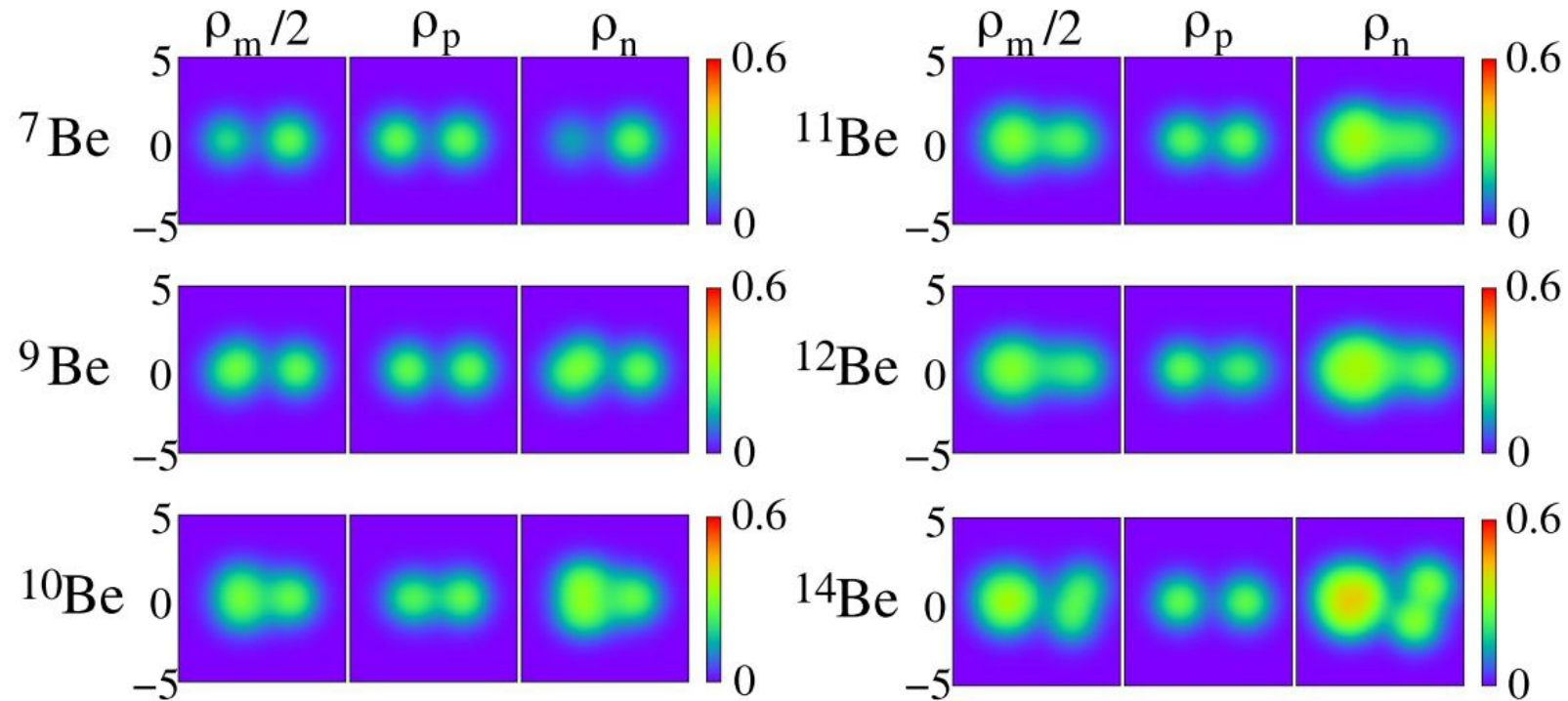


Experiment

I. Tanihata et. al., (Phys. Rev. Lett. 55 (1985) 2676)
T. Minamisono et. al., (Phys. Rev. Lett. 69 (1992) 2058)
I. Tanihata, (J. Phys. G: Nucl. Part. Phys. 22 (1996) 157)

Introduction ~Exotic structures in unstable light heavy-ions~

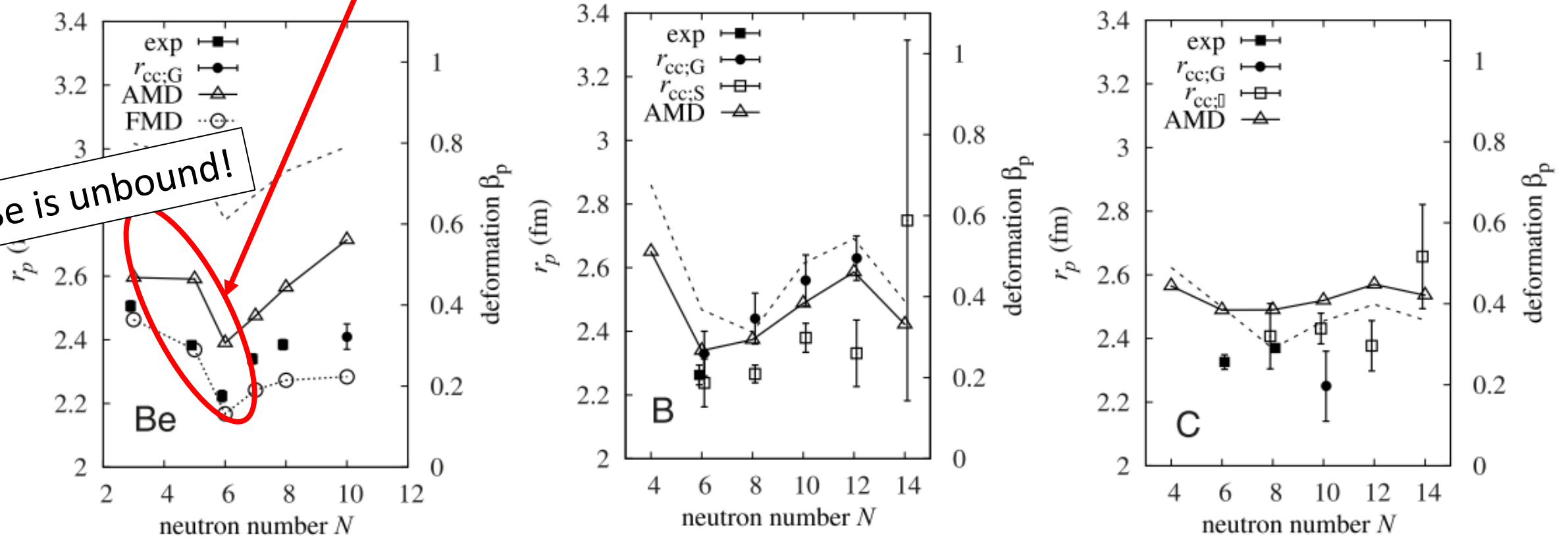
⇒ The cluster feature in unstable nuclei is also hot topics!



Introduction ~Exotic structures in unstable light heavy-ions~

By glue-like effect of the valence neutron(s)

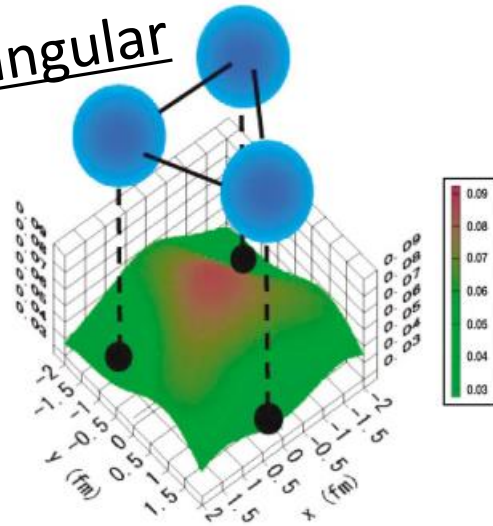
${}^8\text{Be}$ is unbound!



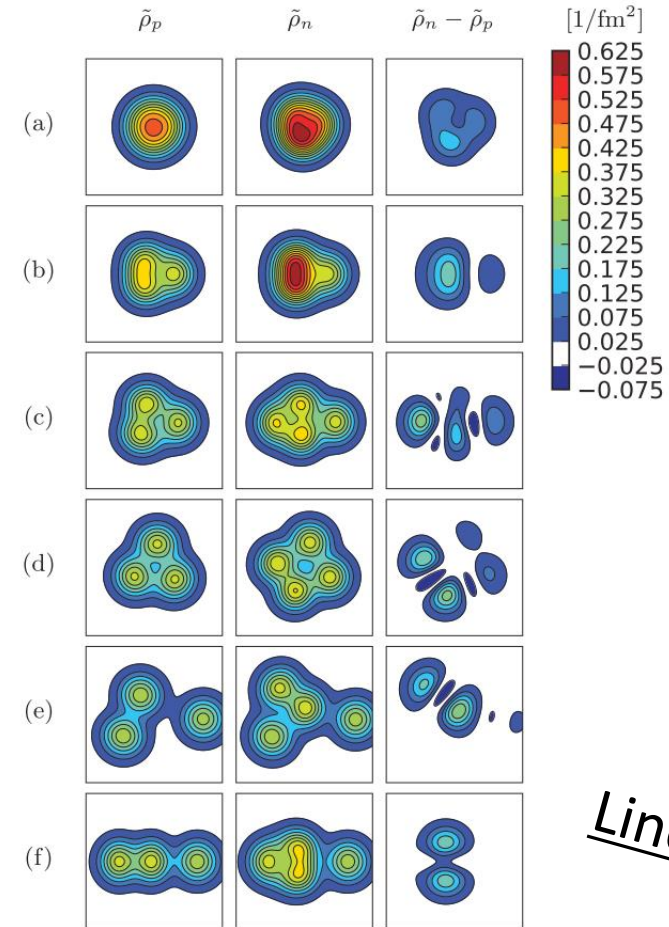
Introduction ~Exotic structures in unstable light heavy-ions~

⇒ The excess neutrons play a glue-like role between alpha-clusters in ^{14}C .

Equilateral-Triangular



N. Itagaki et al., PRL92, 142501 (2004)



Linear-chain

T. Suhara and Y. Kanada-En'yo, PRC82, 044301 (2010)

Formalism (Outline)

To investigate such exotic features in light heavy-ions, we have combined the microscopic **structure** and **reaction** models.

“Structure”

Cluster model + Stochastic multi-configuration mixing



- Binding energy
- Radius
- Transition strength

Transition density



“Reaction”

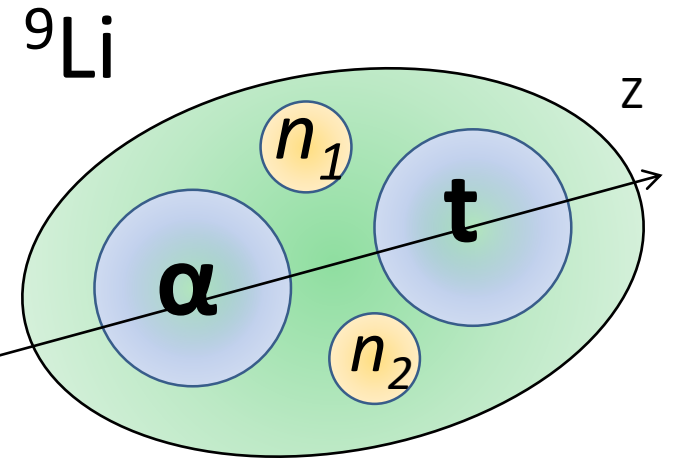
Microscopic Coupled Channel (MCC)



- Cross section

Structure part (${}^9\text{Li}$ case)

Cluster model + Stochastic multi-configuration mixing



$$\Phi^{J^\pi M} = \sum_K \sum_i c_{i,K} \Psi_i^{J^\pi MK}$$

$$\Psi_i^{J^\pi MK} = P^\pi P^{JMK} [A\{\phi_\alpha(\mathbf{r}_1 \mathbf{r}_2 \mathbf{r}_3 \mathbf{r}_4, \mathbf{R}_1) \phi_t(\mathbf{r}_5 \mathbf{r}_6 \mathbf{r}_7, \mathbf{R}_2) \phi_{n_1}(\mathbf{r}_8, \mathbf{R}_3) \phi_{n_2}(\mathbf{r}_9, \mathbf{R}_4)\}]$$

Randomly fixed

$$\phi(\mathbf{r}_j, \mathbf{R}_k) = \exp\left[-\nu\left(\mathbf{r}_j - \frac{\mathbf{R}_k}{\sqrt{\nu}}\right)^2\right] \chi_j$$

χ_j : spin and isospin parts

Hamiltonian

$$\hat{H} = \sum_{i=1}^A \hat{t}_i - \hat{T}_{c.m.} + \sum_{i>j}^A \hat{v}_{ij}$$

$$\nu = 0.235$$

$$M = 0.60$$

$$B = H = 0.08$$

Effective NN interaction

Central : Volkov No.2 (W=1-M)

LS : G3RS ($V_{LS}=2000\text{MeV}$)

Reaction part (Heavy-ion scattering)

Microscopic Coupled Channel (MCC)

$$\left[T_R + U_{\alpha\alpha}(\mathbf{R}) - E_\alpha \right] \chi_\alpha(\mathbf{R}) = - \sum_{\beta \neq \alpha}^N U_{\alpha\beta}(\mathbf{R}) \chi_\beta(\mathbf{R})$$

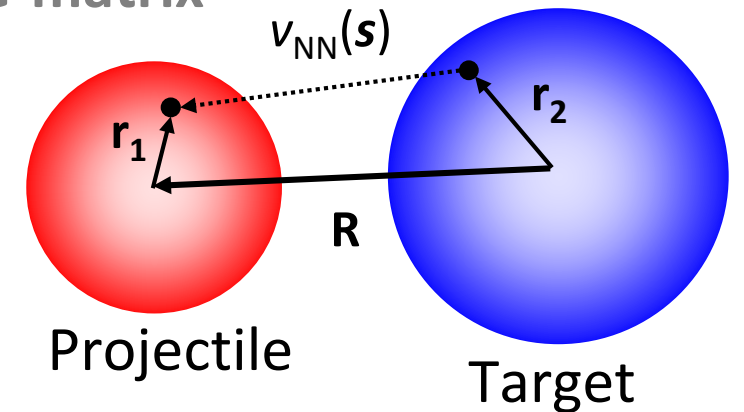
The diagonal and coupling potentials are derived from microscopic view point.

$$U_{\alpha\beta}(\mathbf{R}) = \int \rho_{ik}^{(P)}(\mathbf{r}_1) \rho_{jl}^{(T)}(\mathbf{r}_2) v_{NN}(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2$$

transition density Complex G-matrix

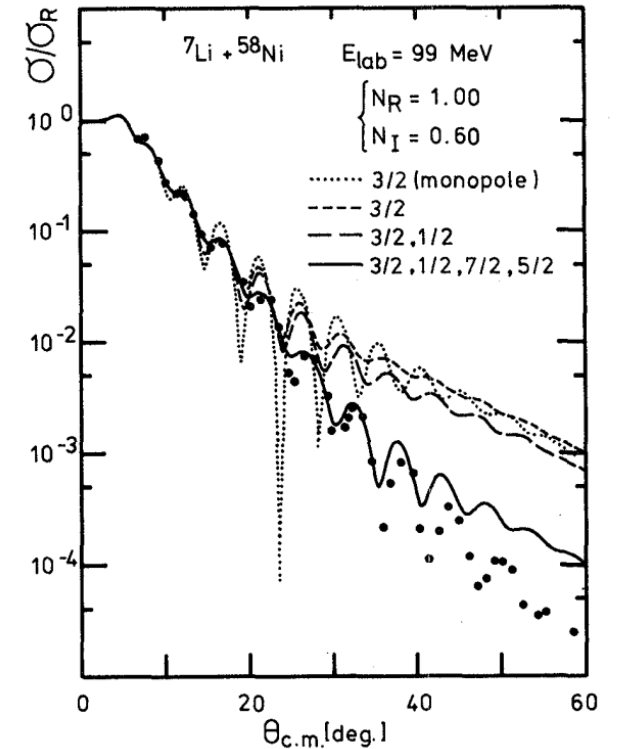
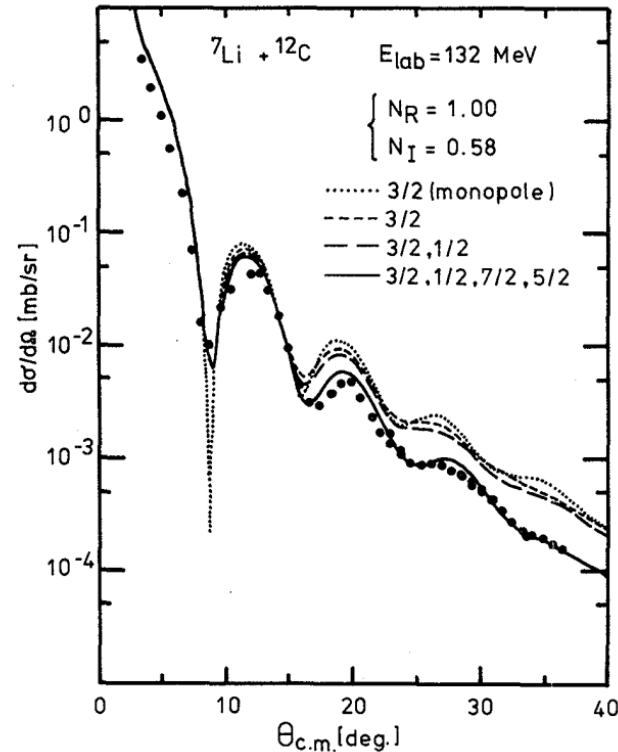
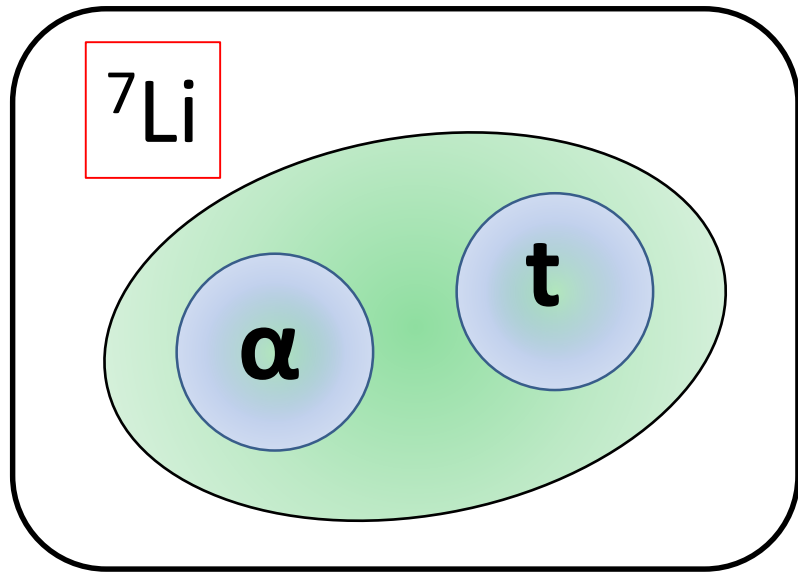
Transition density

$$\rho_{ik}(\mathbf{r}) = \langle \varphi_i(\xi) | \sum_i \delta(\mathbf{r}_i - \mathbf{r}) | \varphi_k(\xi) \rangle$$



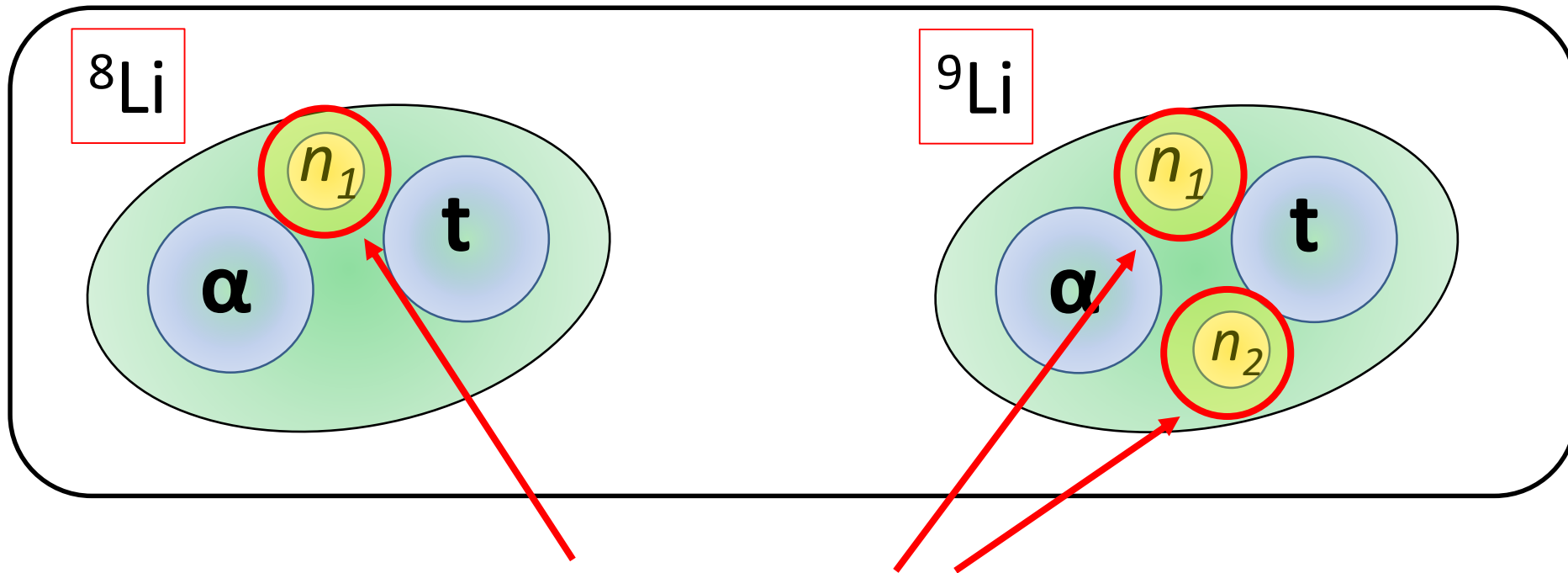
Results \sim Li isotopes \sim

Glue like role of neutron in Li isotopes



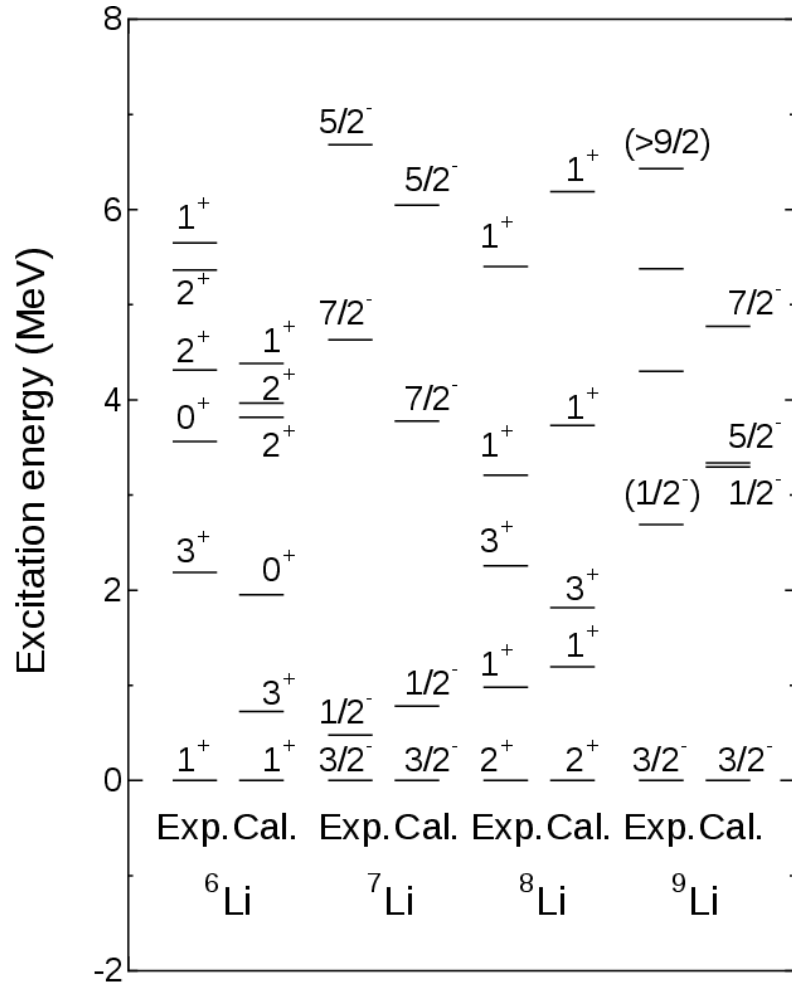
Y. Sakuragi, M. Yahiro and M. Kamimura, PTP89, 136 (1986)

Results *~Li isotopes~*



It is anticipated that the excess neutron(s) plays **a glue-like role** to attract α and t clusters in the ^8Li & ^9Li nuclei.

Comparison with data for excitation energy



Proton radius becomes smaller and smaller.

		Point proton (fm)	Point neutron (fm)	Point matter (fm)
	^6Li	2.54(3)	2.54(3)	2.54(3)
	^7Li	2.43(3)	2.54(3)	2.50(3)
	^8Li	2.41(3)	2.57(3)	2.51(3)
	^9Li	2.30(2)	2.50(2)	2.43(3)
Ref. [52]	^6Li	2.517(30)		
	^8Li	2.299(32)		
	^9Li	2.217(35)		
Calc.	^6Li	2.662 [2.523]	2.556	2.539
	^7Li	2.636 [2.495]	2.604	2.558
	^8Li	2.530 [2.379]	2.615	2.529
	^9Li	2.393 [2.237]	2.562	2.445

[1] I. Tanihata et al., PRL55, 2676 (1985)
 [52] R. Sanchez et al., PRL96, 033002 (2006)

T. Furumoto, T. Suhara, and N. Itagaki, PRC97, 044602 (2018)

Transition strength ($B(IS2)$) of Li isotope

6Li

$B(IS2)$ (g.s. $\rightarrow 3_1$)	57.13 fm ⁴
(g.s. $\rightarrow 2_1$)	54.78 fm ⁴
(g.s. $\rightarrow 1_2$)	33.56 fm ⁴

7Li

$B(IS2)$ (g.s. $\rightarrow 1/2_1$)	54.47 fm ⁴
(g.s. $\rightarrow 5/2_1$)	22.06 fm ⁴
(g.s. $\rightarrow 7/2_1$)	96.78 fm ⁴

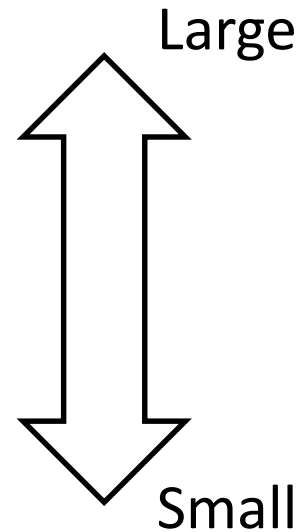
8Li

$B(IS2)$ (g.s. $\rightarrow 1_1$)	16.24 fm ⁴
(g.s. $\rightarrow 3_1$)	47.96 fm ⁴
(g.s. $\rightarrow 2_1$)	7.384 fm ⁴
(g.s. $\rightarrow 2_2$)	15.68 fm ⁴

9Li

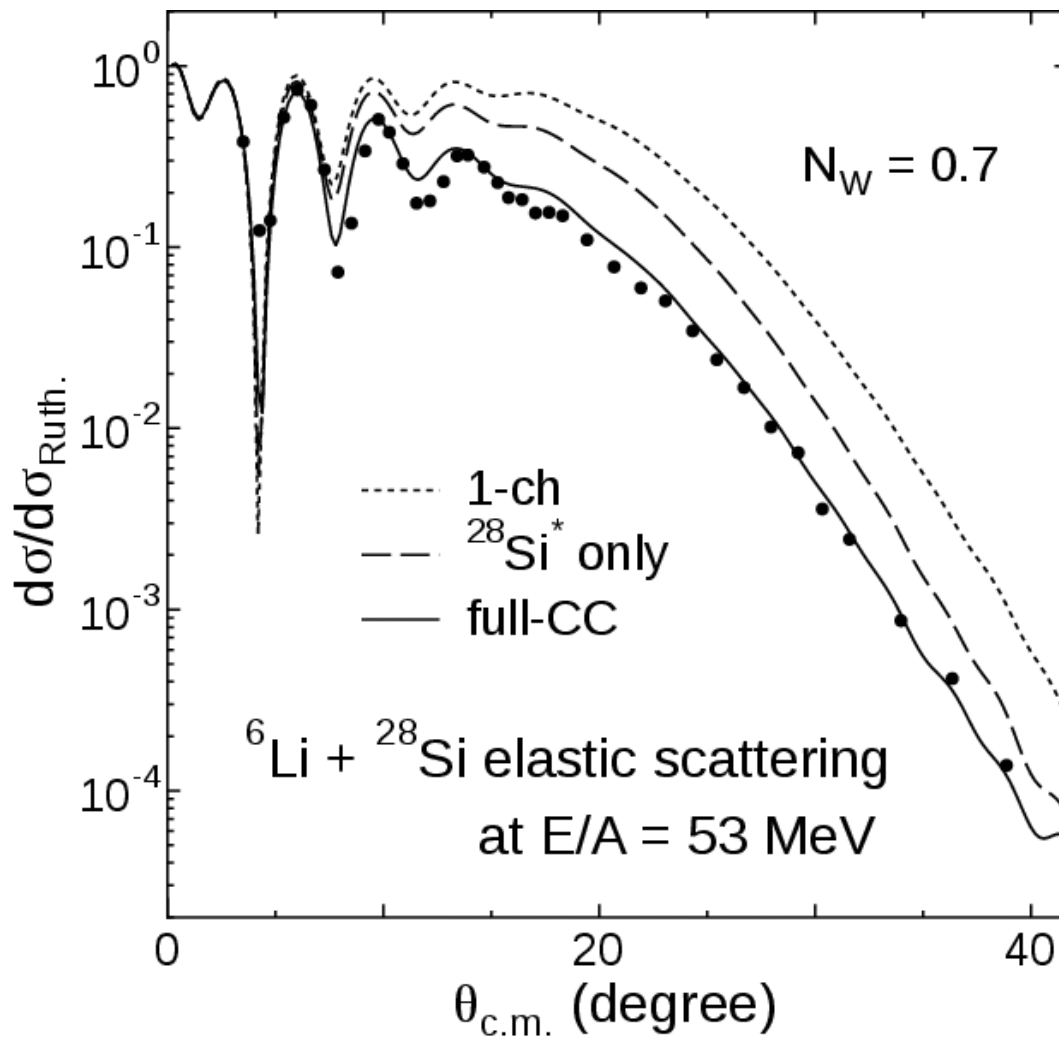
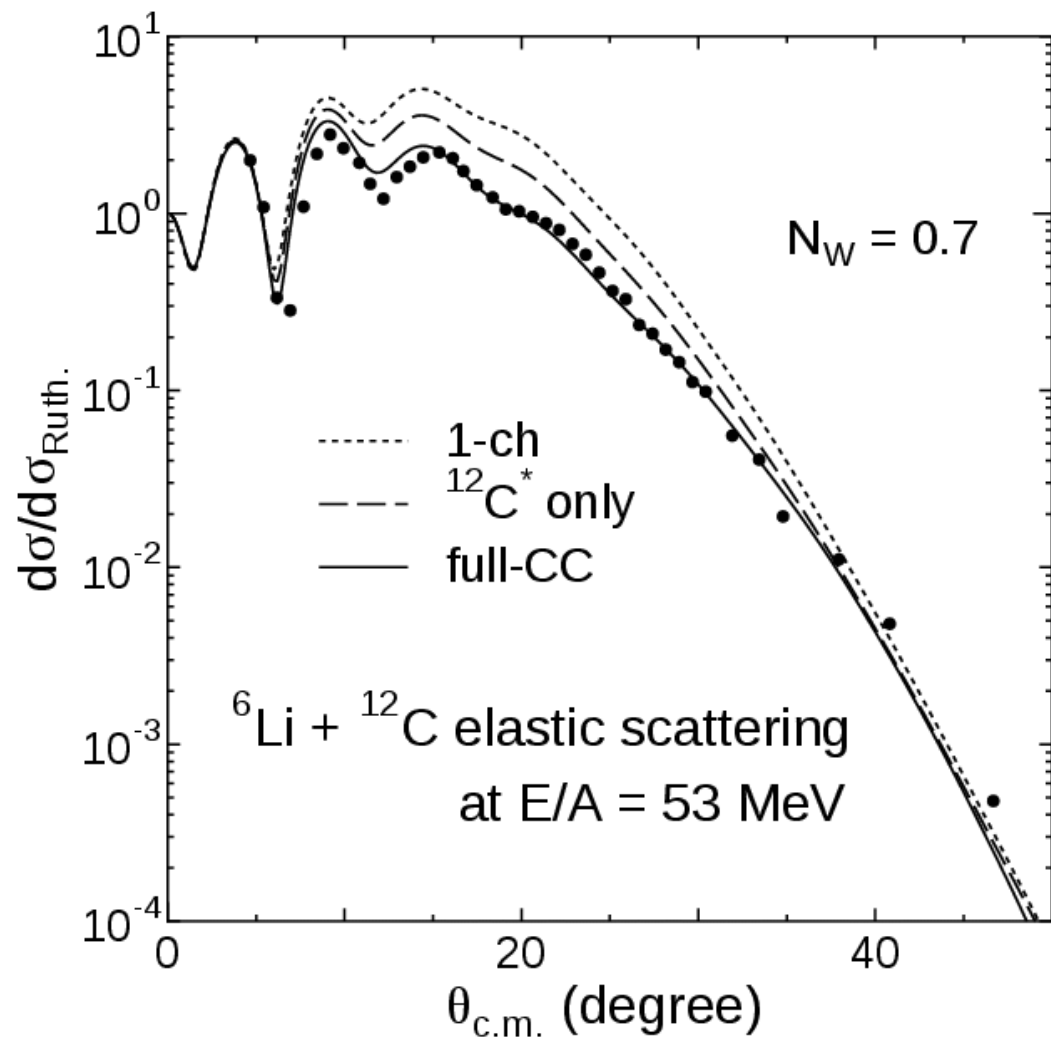
$B(IS2)$ (g.s. $\rightarrow 1/2_1$)	19.26 fm ⁴
(g.s. $\rightarrow 5/2_1$)	19.19 fm ⁴
(g.s. $\rightarrow 7/2_1$)	31.55 fm ⁴

By shrinking the density,
the transition strength becomes small.

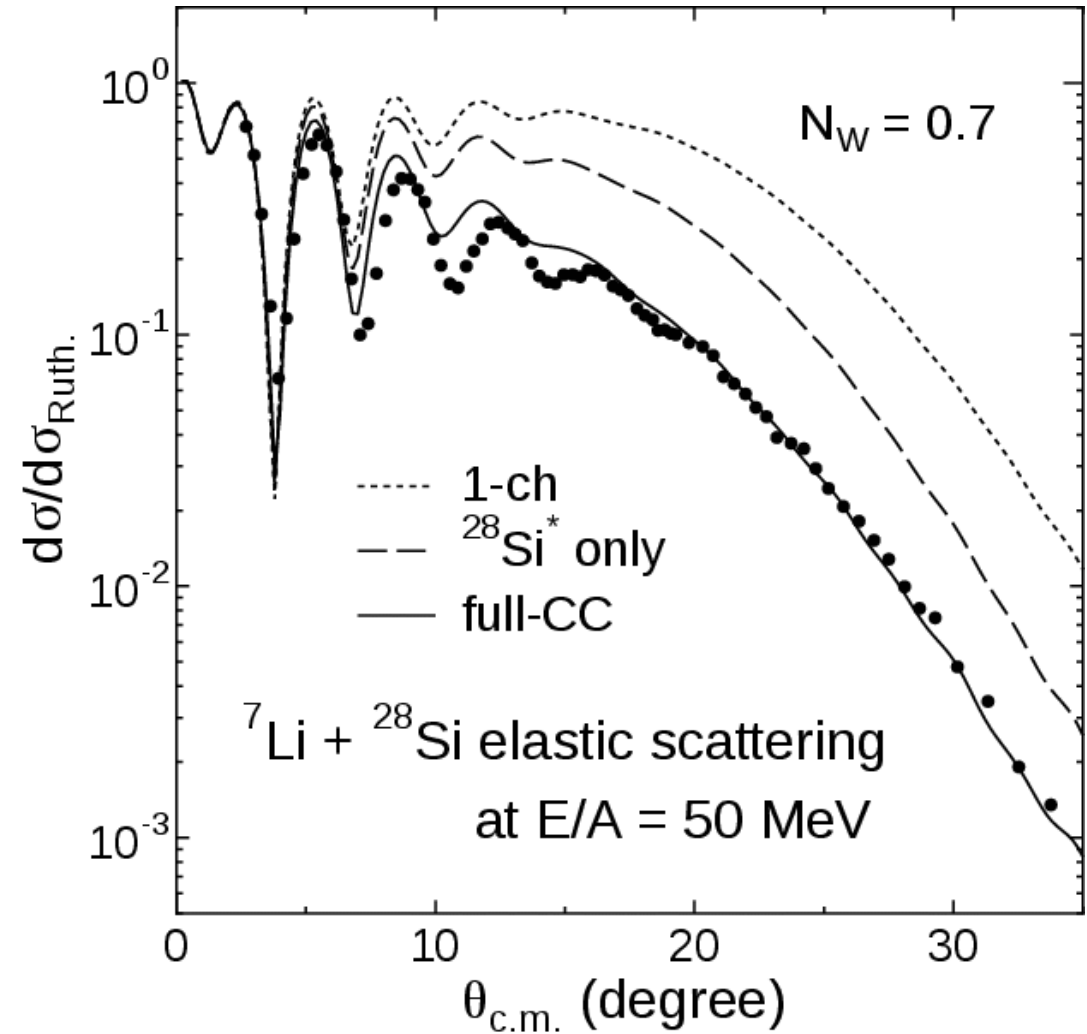
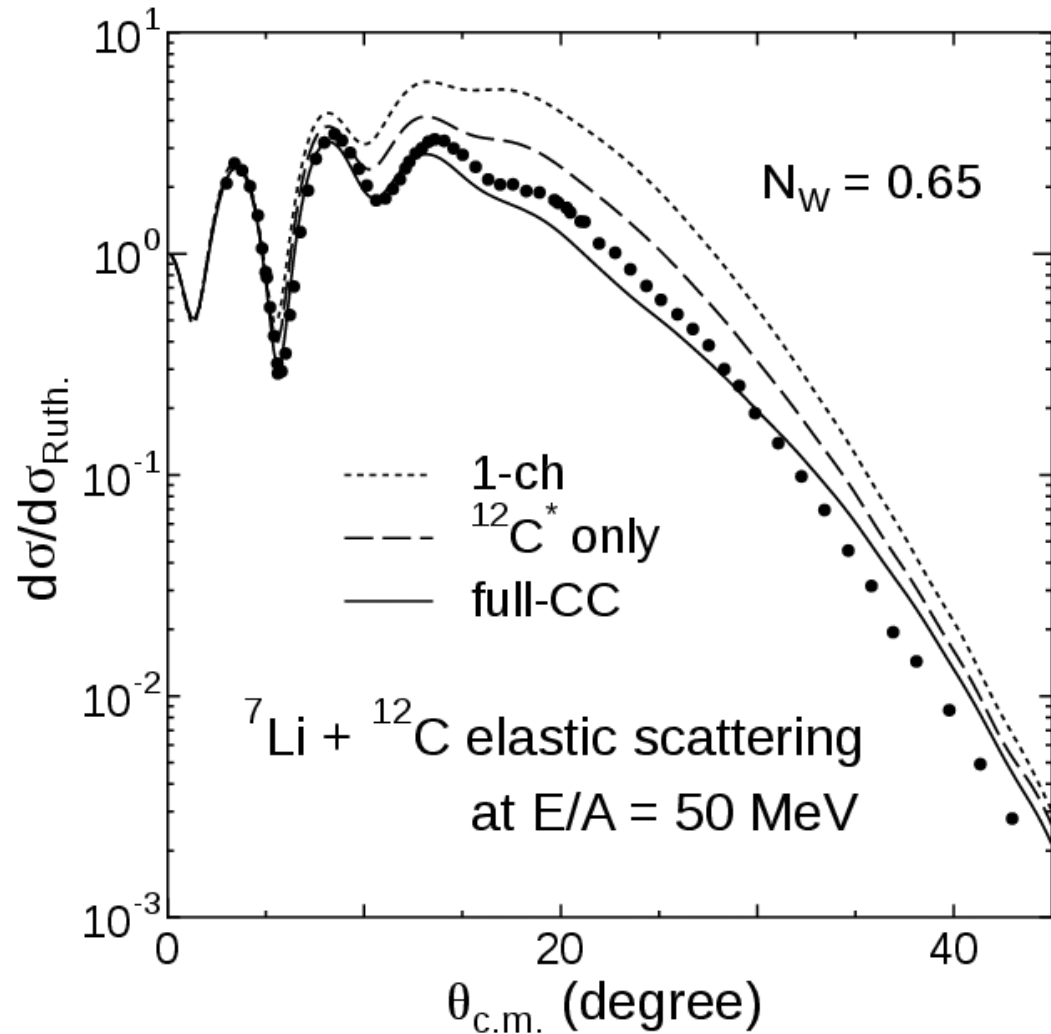


⇒ glue-like role by valence neutrons

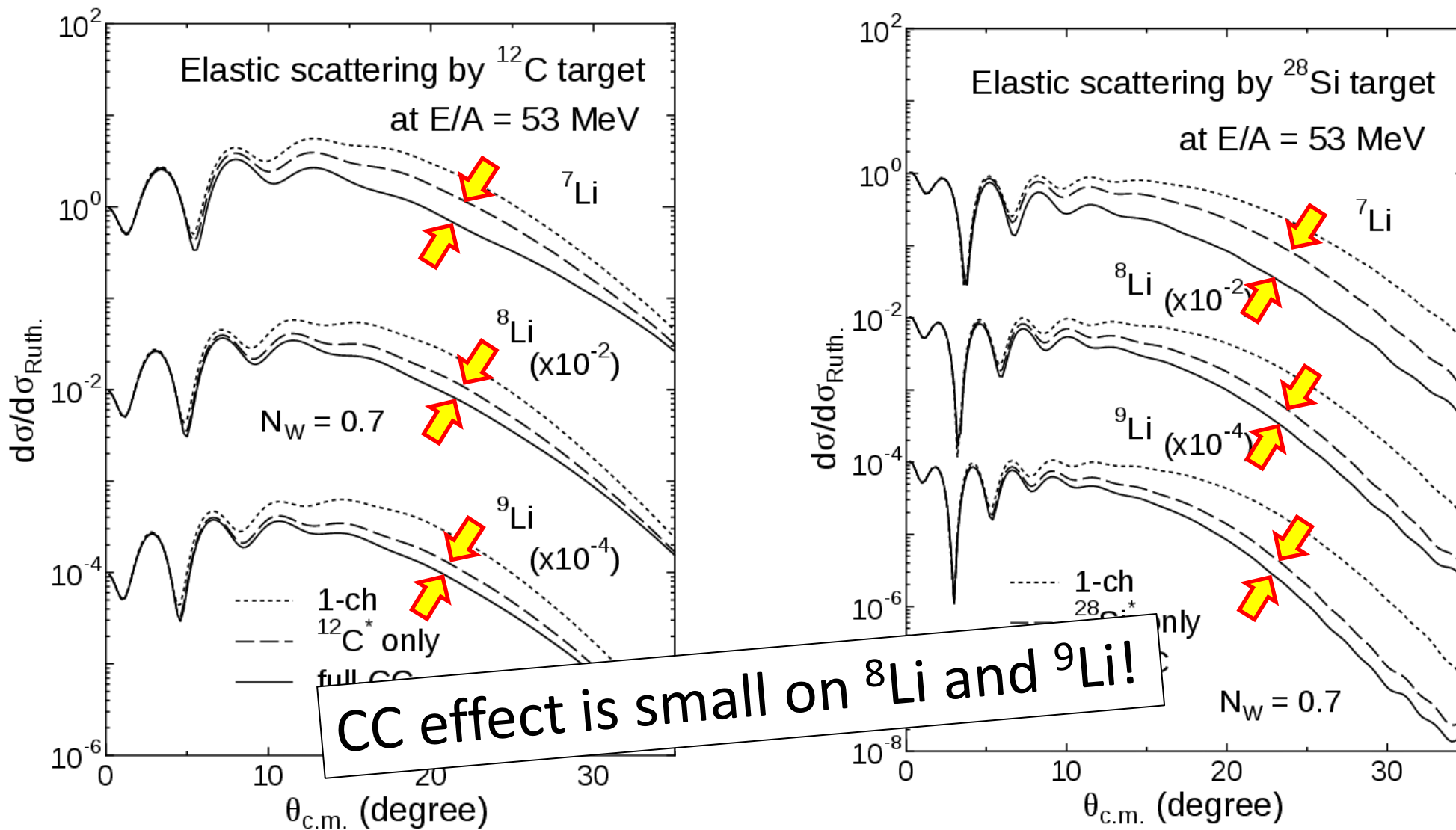
${}^6\text{Li}$ elastic scattering by ${}^{12}\text{C}$ and ${}^{28}\text{Si}$ at $E/A = 53$ MeV



${}^7\text{Li}$ elastic scattering by ${}^{12}\text{C}$ and ${}^{28}\text{Si}$ at $E/A = 50$ MeV



$7\text{-}^9\text{Li}$ elastic scatterings by ^{12}C and ^{28}Si at $E/A = 53$ MeV

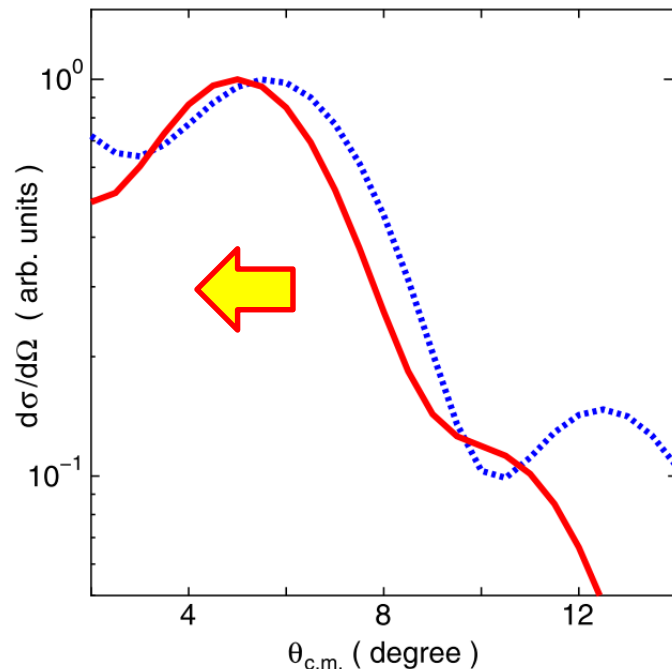


Is it true that the inelastic cross section tells the radius of the excited state?



M. Ito (Phys. Rev. C 97 (2018) 044608)

having an extended radius. Thus, I conclude that the shrinkage in the 2_2^+ cross section originates from the spatially extended structure of the Hoyle rotational 2_2^+ state if the difference in the excitation energies, that is, $E_x = 4.44$ MeV for 2_1^+ and $E_x = 10.3$ MeV for 2_2^+ , can be safely neglected.



To compare the two excited states

TABLE I. Difference in spatial size of the $2_{1,2}^+$ states. ΔR_m represents the difference in the nuclear radius $\Delta R_m = R_m(2_2^+) - R_m(2_1^+)$ with root mean squared radius R_m , while ΔR_{trd} represents the difference in the peak position of the transition density, that is, $\Delta R_{\text{trd}} = R_{\text{trd}}(2_2^+) - R_{\text{trd}}(2_1^+)$. Δa shows the difference in the diffraction radius, which is derived from the differential cross section at $E_\alpha = 386$ MeV. See text for details.

ΔR_m	ΔR_{trd}	Δa
1.6 fm	1.0 fm	0.6 fm

a is inelastic black-sphere radius.

Is it true that the inelastic cross section tells the radius of the excited state?

➔ We apply our models to investigate such mechanisms and systematics.

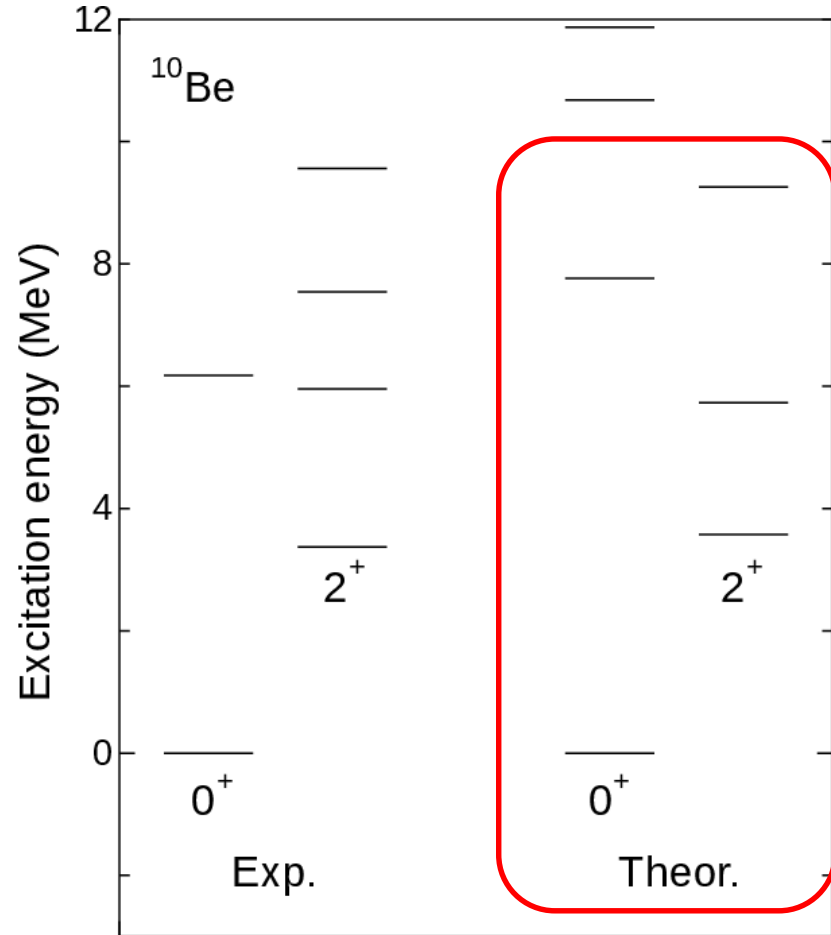
In this test, we choose the well-known nucleus.

→ ^{10}Be nucleus.

- $2\alpha + 2n$ clusters system
- Two neutrons occupy the π orbit in the ground state
but the σ orbit in the 2nd 0^+ state
- Low lying 2^+ states are rotational band members of the 0^+ states

Results $\sim A = 10$ nuclei

➔ Excitation energy and r.m.s. radius



^{10}Be

- 1st 0^+ : 2.526 fm
- 2nd 0^+ : 3.137 fm

- 1st 2^+ : 2.558 fm
- 2nd 2^+ : 2.672 fm
- 3rd 2^+ : 3.158 fm

^{10}Be elastic & inelastic scatterings by ^{12}C target at $E/A = 59.4$ MeV

→ We test radii of the excited states for the ^{10}Be nucleus.

^{10}Be

1st 0⁺ : 2.526 fm

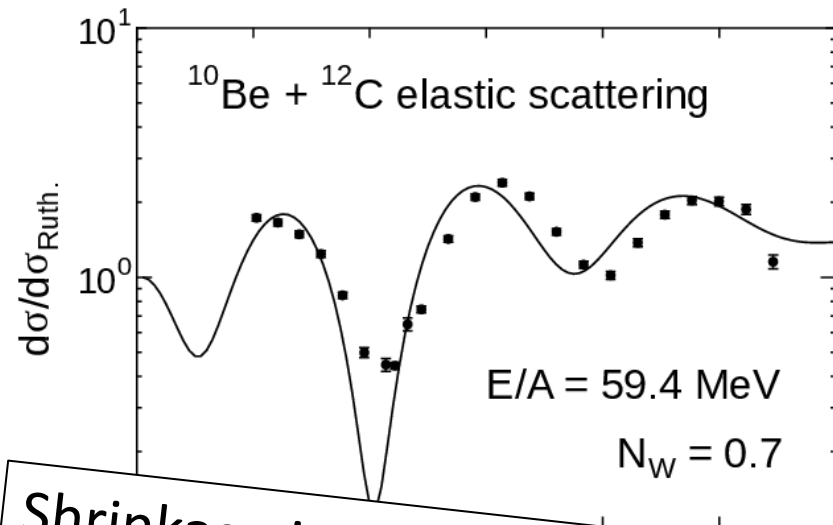
2nd 0⁺ : 3.137 fm

1st 2⁺ : 2.558 fm

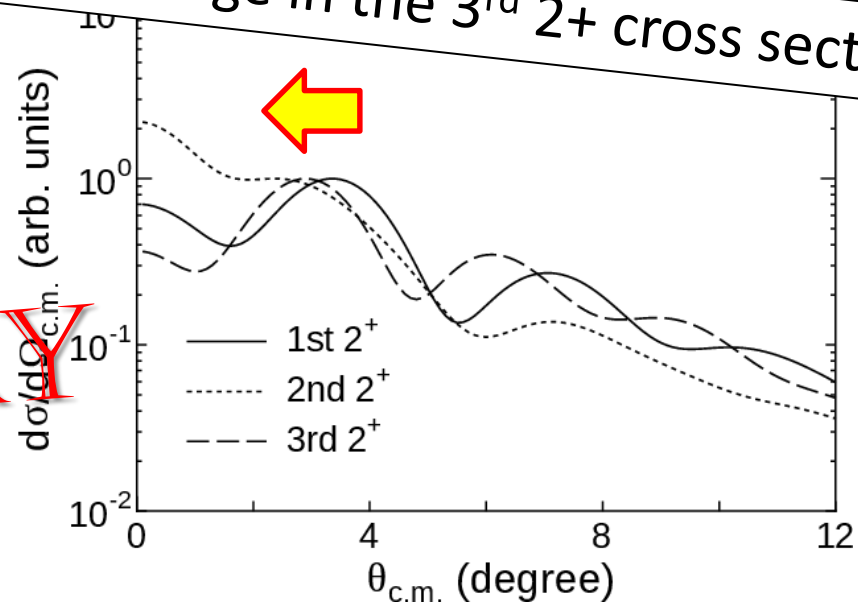
2nd 2⁺ : 2.672 fm

3rd 2⁺ : 3.158 fm

VERY PRELIMINARY



Shrinkage in the 3rd 2⁺ cross section



Summary

- **We have combined microscopic structure and reaction models**

⇒ the structure and reaction parts

are based on the cluster and folding models, respectively.

- **Application to Li isotope scatterings**

⇒ give the feature of glue-like role in Li isotopes

from the structure and reaction viewpoints.

- **Application to ^{10}Be nucleus**

⇒ The size of the excited states for ^{10}Be is investigated based on the previous study.