

Second proton-neutron random phase approximation studied by the Lipkin-Meshkov-Glick (LMG) model in SU(4) group

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First Tsukuba-CCS-RIKEN joint workshop on microscopic theories of nuclear structure and dynamics (12-16 Dec.)

This work was supported



by JSPS KAKENHI Grant No. 25400303.

Background ~ GT Quenching Problem ~

Gamow-Teller (GT) Strength up to E*=30 MeV



Two causes which might produce the quenching

- 1. Δ -h configuration
- 2. multi particle-multi hole configuration

Background ~ β-delayed Neutron ~

<u>Theory</u> can reproduce β Half-lives ...

FRDM+QRPA, RHB+QRPA, HFB+QRPA

β-delayed Neutron Branching Ratio



Main Problem causing the difference may be "Spreading Width."



Background ~ Theoretical Works ~

- 1p1h RPA and TDA with Skyrme force
 <u>96-99%</u> of sum-rule found in E*<20 MeV</p>
- Ф 2p2h TDA (STDA) with Skyrme FM, PRC 93, 044319 (2016)
 <u>80-88%</u> of Sum-rule in E*<20 MeV for 48Ca
 </p>
- Second RPA with G-matrix obtained by Bonn potential (HM3A)
 <u>70-75%</u> of Sum-rule in E*< ~20 MeV for 48Ca & 90Zr S. Drozdz, et al., PLB166, 18 (1986)</p>
- Perturbative calculation
 G.F. Bertsch and I. Hamamoto, PRC26, 1323 (1982).
 About <u>50 %</u> of Sum-rule for 90Zr.

Does really proton-neutron STDA and SRPA work well?

To confirm reliability of proton-neutron STDA & SRPA We adopt the Lipkin-Meshkov-Glick (LMG) Model (enable to compare the model with exact solution)

Several Examples:

1) proton-neutron QRPA in SO(5) group, J. G. Hirsch et al., PRC56, 199 (1997)

2) proton-neutron RPA in SU(2) x SU(2) group, S. Stoica et al., PRC64, 017303 (2001)

3) SRPA in SU(3) group, D. Gambacurta, PRC73, 024319 (2006) -OOOO -OOOO neutron

-OOOO -OOOO proton neutron

Formalism

pnSRPA and pnSTDA are Investigated by LMG model in <u>SU(4) group</u>



- #0 & #1 states are fully occupied by neutron and proton.
- #0 & #2 states are assigned to neutron low and upper states.
- #1 & #3 states are assigned to proton low and upper states, respectively.

Model Hamiltonian

Only $\tau_{+}\tau_{-}, \tau_{-}\tau_{+}$ channels are considered in two-body interaction $H = \sum_{i} e_{i}K_{ii} + \frac{V_{0}(K_{21}K_{12} + K_{30}K_{03}) + V_{1}(2K_{21}K_{30} + 2K_{12}K_{30})}{p-p} + \frac{V_{2}(K_{30}K_{23} + K_{21}K_{32} + K_{30}K_{01} + K_{21}K_{10} + c.c)}{3p-1h \& 1p-3h} + \frac{V_{3}(K_{23}K_{32})}{4p}$

--- Model parameters ---

 $V_0 = 0.2\epsilon/N$ $V_1 = \kappa\epsilon/N$ $V_2 = -\kappa\epsilon/(5N)$ $V_3 = \kappa\epsilon/(25N)$ $N = N_n + N_p$ We will see the excitation energy and transition strength as a function of κ

Exact Wave Function

$$\begin{split} |\Phi\rangle &= \sum_{\alpha,\beta,\gamma,\delta} C_{\alpha\beta\gamma\delta} |N_n - \beta - \gamma, N_p - \alpha - \delta, \alpha + \gamma, \beta + \delta \rangle \\ &= \sum_{\alpha,\beta,\gamma,\delta} C_{\alpha\beta\gamma\delta} \left(K_{21}\right)^{\alpha} \left(K_{30}\right)^{\beta} \left(K_{20}\right)^{\gamma} \left(K_{31}\right)^{\delta} |HF\rangle \\ &= \sum_{\alpha,\beta,\gamma,\delta} C_{\alpha\beta\gamma\delta} \left(K_{21}\right)^{\alpha} \left(K_{30}\right)^{\beta} \left(K_{20}\right)^{\gamma} \left(K_{31}\right)^{\delta} |HF\rangle \\ &= C_{\alpha\beta\gamma\delta} \text{ are derived from diagonalization of } H \text{ with } |\Phi\rangle \quad 7 \end{split}$$

Formalism



RPA ground state $\begin{array}{c} \downarrow \\ Q | RPA \rangle = 0 \\ | RPA \rangle \Rightarrow | HF \rangle
\end{array}$

RPA Equation

D. Gambacurta, PRC73, 024319 (2006)

 $\begin{pmatrix} A & B \\ B^* & A^* \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix} = \hbar \omega \begin{pmatrix} G & 0 \\ 0 & -G^* \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix}$

proton neutron SRPA

SRPA Equation

$$Q^{\dagger} = \frac{X\Theta^{+} - Y\Theta^{-}}{\sqrt{\langle HF|[\Theta^{-},\Theta^{+}]|HF\rangle}} + \frac{X\theta^{+} - Y\theta^{-}}{\sqrt{\langle HF|[\theta^{-},\theta^{+}]|HF\rangle}} \qquad \begin{pmatrix} \mathcal{A} & \mathcal{B} \\ \mathcal{B}^{*} & \mathcal{A}^{*} \end{pmatrix} \begin{pmatrix} \mathcal{X} \\ \mathcal{Y} \end{pmatrix} = \hbar\omega \begin{pmatrix} \mathcal{G} & 0 \\ 0 & -\mathcal{G}^{*} \end{pmatrix} \begin{pmatrix} \mathcal{X} \\ \mathcal{Y} \end{pmatrix}$$

Transition Operator
$$\mathcal{A} = \begin{pmatrix} A & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \quad \mathcal{B} = \begin{pmatrix} B & 0 \\ 0 & 0 \end{pmatrix}$$

 β^{-} transition

 β^+ transition

$$M^{-} = \chi \sum_{mm'} a_{3m}^{\dagger} a_{0m'}$$
 $M^{+} = \chi \sum_{mm'} a_{2m}^{\dagger} a_{1m'}$

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Result: Excitation Energy



At small κ, RPA, SRPA, STDA are in a good agreement with the exact.
 For 1st Excitation energy, RPA & SRPA are close.
 For 2nd Excitation energy, SRPA & STDA are close.

Result: Transition Strength of β-



1) SRPA takes into account ground state correlation to some extent, which STDA doesn't.

2) The same difference is obtained between RPA and TDA \rightarrow



3) In the charge exchange channel, the ground state correlation generally doesn't play a significant role. Namely, we usually see results at small κ in a realistic nucleus (see Next Slide)

4) But we have used zero-range force in STDA. We might have used a large κ ?? \leftarrow Further analysis is in progress.

Result: 120Sn Fermi Transision



If ground state significantly differs from HF state, TDA shows different distribution from RPA.

Summary

Investigate proton-neutron SRPA & STDA with LMG Model in SU(4) group

1) At small κ, SRPA & STDA work well for excited state But, STDA is now accurate enough for transition strength

This might be one of the factors producing difference SRPA and STDA calculations on the GT quenching effect.

2) SRPA seems to work well as compared the exact solution. Higher correlation is required ? or we cannot exclude a contribution of Δ -h coupling to explain the quenching problem.