Non-perturbative renormalization of B_K and quark mass in SF scheme with quenched domain-wall QCD Yousuke Nakamura, Yusuke Taniguchi

for CP-PACS Collaboration

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Plan of talk

- **Definiton of** *B_K* and its renormalization factor.
- **Strategy to derive RGI** \widehat{B}_K
- SF scheme (definition of non-perturbative renormalization factor) Alpha collaboration '05
- Simulation results
 - **• RGI** \hat{B}_K
 - **Scaling behavior of SSF at** $g^2(L) = 3.480$
- Renormalization of quark mass

Summary

Definition of *B_K* **and its renormalization factor**

J The kaon *B* parameter: B_K

$$B_K \equiv \frac{\langle \bar{K^0} | \bar{s} \gamma^{\mu} (1 - \gamma_5) d \bar{s} \gamma^{\mu} (1 - \gamma_5) d | K^0 \rangle}{\frac{8}{3} \langle \bar{K^0} | \bar{s} \gamma^{\mu} (1 - \gamma_5) d | 0 \rangle \langle 0 | \bar{s} \gamma^{\mu} (1 - \gamma_5) d | K^0 \rangle}$$

CP-PACS obtained B_K with quenched domain-wall fermion 2001 \Rightarrow We need non-perturbative renormalization for B_K

- Renormalization factor for B_K $Z_{B_K}(a\mu) = \frac{Z_{VV+AA}(g_0, a\mu)}{Z_A^2(g_0)} = \frac{Z_{VA+AV}(g_0, a\mu)}{Z_V^2(g_0)}$
 - **DWF** has a good chiral symmetry even on lattice
 - $Z_V(g_0) = Z_A(g_0)$ has been showen by CP-PACS 2003

$$Z_{VV+AA}(g_0, a\mu) = Z_{VA+AV}(g_0, a\mu)$$

 $\cdot \Rightarrow$ this relation will be checked later

Our strategy



- **D** Three steps for RGI \hat{B}_K via SF scheme.
 - I. $B_K^{(0)} \Rightarrow B_K^{(SF)}(a\mu_{\min})$ at hadronic scale suppress lattice artifact $O(a\mu)$: $a\mu \ll 1$
 - **II.** Non-pert. RG running : $B_K^{(SF)}(\mu_{\min}) \Rightarrow B_K^{(SF)}(\mu_{\max})$
 - **III.** Perturbative running to RGI: $B_K^{(SF)}(\mu_{max}) \Rightarrow \hat{B}_K$
- II. III. are reguralization independent. Possible to evaluate with cheap Wilson fermion. Already calculated by Alpha collab. for O_{VA+AV} '05 ⇒ O_{VA+AV} has no mixing problem even for wilson fermion.
 Our target: I. Z_{VA+AV}(g₀, aμ_{min}) with DWF.

Set up for SF scheme Alpha '05





$$\mathcal{F}^{\pm}_{[\Gamma_A,\Gamma_B,\Gamma_C]}(x_0) = \frac{1}{L^3} \langle \mathcal{O}_{21}[\Gamma_A] \mathcal{O}_{45}[\Gamma_B] \mathcal{O}^{\pm}_{VA+AV}(x) \mathcal{O}'_{53}[\Gamma_C] \rangle$$

Orbifolding construction for SF with DWF Y.T '04,'06

In order to cancel divergence due to boundary operator

boundary-boundary correlation function

$$f_1 = -\frac{1}{2L^6} \langle O'_{12}[\gamma_5] O_{21}[\gamma_5] \rangle, \quad k_1 = -\frac{1}{2L^6} \sum_k \langle O'_{12}[\gamma_k] O_{21}[\gamma_k] \rangle$$

Three specific cases



Renormalization condition

 $Z_{VA+AV;s}(g_0, a\mu)h_s^{\pm}(x_0; g_0) = h_s^{\text{cont}}(x_0, g_0)^{\pm}|_{g_0=0}, \quad s = 1, 3, 7$

Renormalization factor

$$Z_{VA+AV;s}^{\pm} = \frac{h_s^{\pm}(x_0;g_0)|_{g_0=0}}{h_s^{\pm}(x_0;g_0)}$$

Simulation parameters

- Same parameters as in previous CP-PACS calculation of B_K Phys.Rev. D64(2001)114506
 - **Domain-wall fermion with** M = 1.8 and $N_5 = 16$
 - Iwasaki gauge action at β = 2.6, β = 2.9, β = 3.2(new,not in the paper)
 corresponding to a⁻¹ ~ 2, a⁻¹ ~ 3, a⁻¹ ~ 4 GeV
- **Renormalization factor** $Z_{VA+AV}(\beta, a\mu)$

•
$$1/\mu_{\min} = 2L_{\max}$$
 Def. of L_{\max} : $\bar{g}_{SF}^2(L_{\max}) = 3.480$

 Z_{VA+AV} at $\beta = 2.6, \beta = 2.9, \beta = 3.2$

Fine tuning of β to give $aN_T = aN_L = 2L_{max} = 1.498r_0$

$L_{\rm max}/a$	6	8	10	12	14	16	18
β	2.446	2.6339	2.7873	2.9175	3.03133	3.0313	3.2254
# of conf	1000	1000	1000	1018	716	556	564

Simulation results

• Numerical result for Z_V and $Z_{VA+AV;s}$



We take the value at $x_0 = N_T/2$ (Alpha collaboration '05)

$$Z_{B_K}(a\mu_{\min}) = \frac{Z_{VA+AV;s}(x_0 = 16/2, a\mu_{\min})}{Z_V^2(x_0 = 16/2, a\mu_{\min})} = 1.2708(53)$$

Strategy to derive RGI \widehat{B}_K

J Lattice bare $B_K^{(0)}$ with DWF

 \Rightarrow **Renormalization group invariant (RGI)** \hat{B}_K

•
$$\mathcal{Z}_{B_K}(g_0) = \underbrace{Z_{VA+AV}^{\text{PT}}(\infty, \mu_{\max})}_{\text{III}} \underbrace{Z_{VA+AV}^{\text{NP}}(\mu_{\max}, \mu_{\min})}_{\text{II}} \underbrace{Z_{B_K}^{\text{NP}}(g_0, a\mu_{\min})}_{\text{I}}$$

Result for RGI Z_{B_K} and \hat{B}_K (quench)



• Fitting form $\hat{Z}_{B_K}(\beta) = a_1 + b_1(\beta - 3) + c_1(\beta - 3)^2$ interpolated at $\beta = 2.6, \beta = 2.9, \beta = 3.2$

 $\hat{B}_{K} = 0.773$ (7)(preliminary) : Scheme 1 $\hat{B}_{K} = 0.760$ (7)(preliminary) : Scheme 3 $\hat{B}_{K} = 0.778$ (7)(preliminary) : Scheme 7 $\hat{B}_{K} = 0.786(31)$: RBC $\hat{B}_{K} = 0.789(46)$: Alpha

- **Error is smaller**
 - Scaling behaviour is better
 - Error of each data is smalle
 - The smallest lattice spacing is finer

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Conversion of RGI B_K to that in \overline{MS} scheme

Solution Conversion from RGI to \overline{MS} NDR at $\mu = 2$ GeV by NLO PT

$$B_{K}^{\overline{\text{MS}}}(\text{NDR},\mu) = \left[\frac{g_{\overline{\text{MS}}}^{2}(\mu)}{4\pi}\right]^{\frac{\gamma_{O}^{+}}{2b_{0}}} \exp\left[\int_{0}^{g_{\overline{\text{MS}}}(\mu)} dg\left(\frac{\gamma^{+}}{\beta(g)} - \frac{\gamma_{0}^{+}}{b_{0}g}\right)\right] \hat{B}_{K}$$

Gauge coupling

$$\Lambda_{\overline{\text{MS}}} = \mu (b_0 g_{\overline{\text{MS}}}^2)^{-\frac{b_1}{2b_0^2}} \exp\left[-\frac{1}{2b_0 g_{\overline{\text{MS}}}^2}\right] \exp\left[-\int_0^{g_{\overline{\text{MS}}}} dg \left(\frac{1}{\beta_{\overline{\text{MS}}}(0)} + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right]$$

• $\Lambda_{\overline{\text{MS}}} = 0.586(48)/r_0$ Necco&Sommer '02

Comparison of our result for $B_K(\overline{MS}, 2\text{GeV})$ with previous ones



 $B_{K}^{\overline{MS}}(NDR, 2GeV) = 0.557(5)(_{-10}^{+4}) \text{ (preliminary)}$ $B_{K}^{\overline{MS}}(NDR, 2GeV) = 0.567(4) \quad : CP-PACS(perturbative ren.)$ $B_{K}^{\overline{MS}}(NDR, 2GeV) = 0.563(21) \quad : RBC$ $B_{K}^{\overline{MS}}(NDR, 2GeV) = 0.573(34) \quad : Alpha$

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Further checks

Scaling behavior of Step Scaling Function

$$\Sigma_{VA+AV;s}(u, L_{\max}/a) = \frac{Z_{VA+AV;s}(g_0, a/2L_{\max})}{Z_{VA+AV;s}(g_0, a/L_{\max})}$$

Chiral symmetry breaking effect in renormalization factor

 $Z_{VA+AV} = Z_{VV+AA}$



Boundary O(a) effect is large at M = 1.8





No *N*₅ dependence

Our solution

 $M_{\rm tadpole} \sim 1$

SOLUTION 1 : results at M = 1.4



Scaling violation is small

S Continuum limit is consistent with Alpha's

SOLUTION 2 : Tree-level improvement

 $Z_{VA+AV;s}(g_0, a\mu)h_s^{\pm}(x_0; g_0) = h_s(x_0, g_0)^{\pm}|_{g_0=0}^{\text{CONT}} \Rightarrow h_s(x_0, g_0)^{\pm}|_{\text{tree}}^{\text{Lattice}}$



- Scaling behavior is improved
- Continuum limit is consistent with Alpha's

Scaling behavior of SSF for Z_{B_K}

$$\Sigma_{B_K}(u, L_{\max}/a) = \frac{Z_{B_K;s}(g_0, a/2L_{\max})}{Z_{B_K;s}(g_0, a/L_{\max})}$$



- **Solution** Results at M = 1.4 and M = 1.8 are consistent
- O(a) error is patrly cancelled in Z_{B_K}



• WT Identity

•
$$\delta q_1 = -i\gamma_5 \tilde{q}_1, \quad \delta \zeta_1 = i\gamma_5 \tilde{\zeta}_1, \quad \delta \zeta'_1 = i\gamma_5 \tilde{\zeta}_1'$$

 $\langle O_{VA+AV}O[\zeta] \rangle_S = \langle O_{VV+AA}O'[\zeta] \rangle_S$
 $\Rightarrow Z_{VV+AA} = Z_{VA+AV}$

• For the domain-wall case

• $S \to S + Y : Y$ is the chiral symmetry breaking term $\langle O_{VA+AV}O[\zeta] \rangle_S = \langle O_{VV+AA}O'[\zeta] \rangle_{S+Y}$ $\neq \langle O_{VV+AA}O'[\zeta] \rangle_S$

> $\langle O_{VA+AV}O[\zeta] \rangle_S \Leftrightarrow \langle O_{VV+AA}O'[\zeta] \rangle_S$ We investigate if $Z_{VV+AA} = Z_{VA+AV}$ holds or not.

• **Result for** $Z_{VA+AV;1}$ **and** $Z_{VV+AA;1}$ ($N_5 = 16$)



• Good agreement between $Z_{VA+AV;1}$ and $Z_{VV+AA;1}$ • $Z_{B_K} = \frac{Z_{VA+AV}}{Z_V^2}$: justified

Quark mass renormalization

RGI quark mass

$$M = \overline{m}(\mu) \left(2b_0 \overline{g}^2(\mu)\right)^{-\frac{d_0}{2b_0}} \exp\left(-\int_0^{\overline{g}(\mu)} dg\left(\frac{\tau(g)}{\beta(g)} - \frac{d_0}{b_0g}\right)\right)$$

$$\mathcal{Z}_m(g_0) = \underbrace{Z_m^{\text{PT}}(\infty, \mu_{\text{max}})}_{\text{III}} \underbrace{Z_m^{\text{NP}}(\mu_{\text{max}}, \mu_{\text{min}})}_{\text{II}} \underbrace{Z_m^{\text{NP}}(g_0, a\mu_{\text{min}})}_{\text{II}}$$

- Lattice bare mass with DWF \Rightarrow RGI mass
- II, III: given by Alpha collab. for PCAC mass
- $Z_m(g_0, a\mu_{\min}) = 1/Z_P(g_0, a\mu_{\min})$ for DWF



MS scheme at
$$\mu = 2$$
 GeV

$$m^{\overline{\mathrm{MS}}}(\mu) = M \left(2b_0 \left(g^{\overline{\mathrm{MS}}}(\mu) \right)^2 \right)^{\frac{d_0}{2b_0}} \exp\left(\int_0^{g^{\overline{\mathrm{MS}}}(\mu)} dg \left(\frac{\tau^{\overline{\mathrm{MS}}}(g)}{\beta^{\overline{\mathrm{MS}}}(g)} - \frac{d_0}{b_0 g} \right) \right)$$

• $\Lambda_{\overline{\text{MS}}} = 0.586(48)/r_0$ as an input.



 $m_{u,d}(\mu = 2\text{GeV}) = 4.020(58) \text{ MeV}, \quad m_s(\mu = 2\text{GeV}) = 105.4(15) \text{ MeV}$

Summary

- **Solution Evaluation of NP renormalization factor for** B_K
 - Bare $B_K^{(0)}$ on lattice with DWF \Rightarrow RGI \hat{B}_K
 - Investigation of chiral breaking effect
- **P RGI** $\hat{B}_K = 0.773(7)$ (preliminary)
- $\square \overline{\text{MS}} B_K(\text{NDR}, 2\text{GeV}) = 0.557(5)(^{+4}_{-10}) \text{ (preliminary)}$
 - Consistent with previous results
- Check of scaling behavior of SSF
- Evaluation of NP renormalization factor for quark mass

 $m_{u,d}(\mu = 2\text{GeV}) = 4.020(58) \text{ MeV}, \quad m_s(\mu = 2\text{GeV}) = 105.4(15) \text{ MeV}$