

Resonances in Lattice QCD

External Review on CCS, 2014/02/19

Naruhito Ishizuka with PACS-CS ,
University of Tsukuba

Results of ρ meson decay by PACS-CS

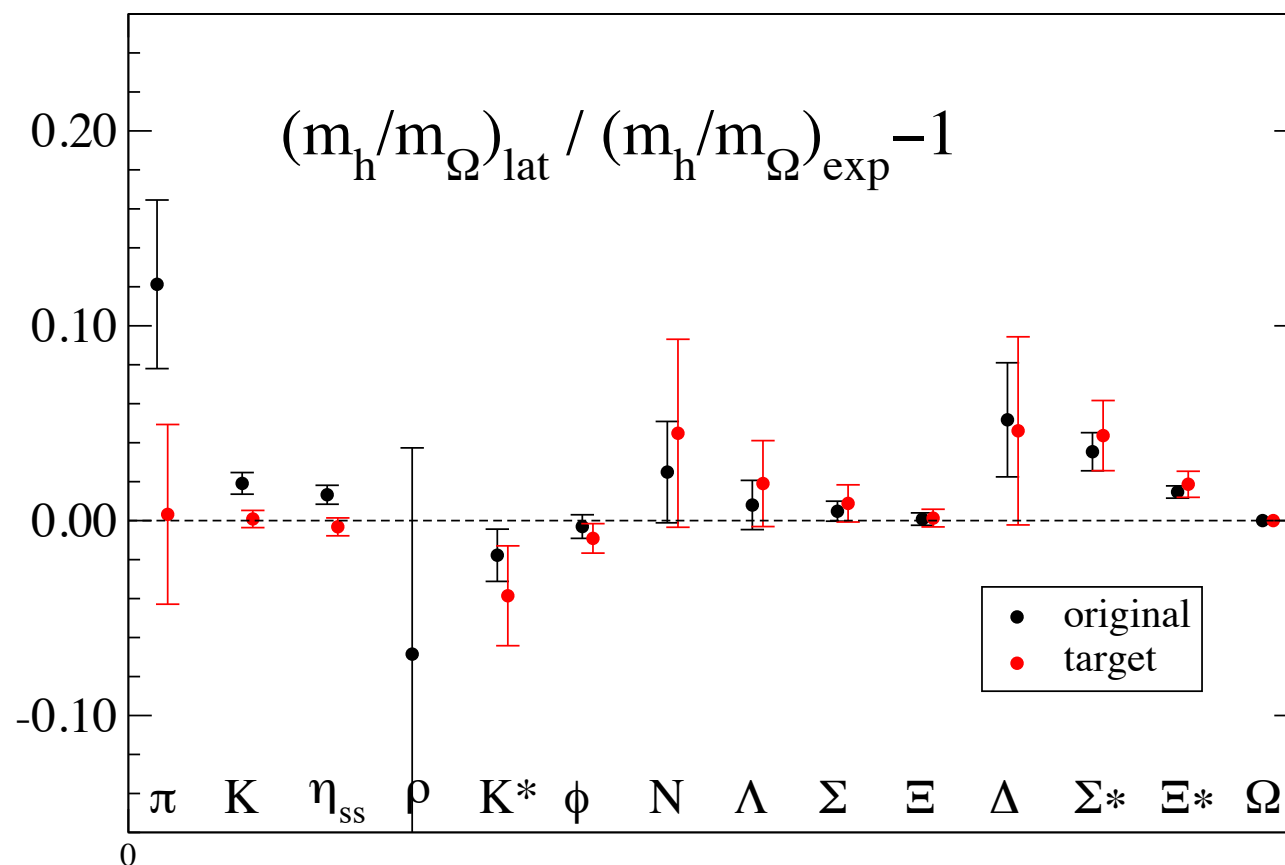
1. Introduction

Now, we can calculate hadron masses
at physical quark mass by lattice QCD.

But, it is only for stable particles

Hadron masses at physical quark mass

PACS-CS PRD81(2009)074503



: for unstable particle (ρ , K^* , Δ)
energies of ground states
on finite volume are plotted.

These are **not resonance masses**.

Calculations of resonance mass
and decay width of unstable particles still remain.

Comment on time correlation function

Calculation of mass of stable particle

$\mathcal{O}(t)$: Heisenberg op. on lattice

$$\begin{aligned} \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(t_0) | 0 \rangle \\ = \sum_H \left| \langle 0 | \mathcal{O} | H \rangle \right|^2 \cdot e^{-E_H \cdot (t-t_0)} \end{aligned}$$

$|H\rangle$: energy eigenstate
 E_H : energy eigvalue

for $t \gg t_0$

$$\sim \left| \langle 0 | \mathcal{O} | H_0 \rangle \right|^2 \cdot e^{-E_0 \cdot (t-t_0)}$$

➡ energy of ground state : E_0 by exp-fit

for ρ meson (resonance state)

($\rho \rightarrow \pi\pi$ by strong interaction)

$$\langle 0 | \rho^\dagger(t) \rho(0) | 0 \rangle = \sum_{\alpha} |\langle 0 | \rho | \alpha \rangle|^2 \cdot e^{-E_{\alpha} t}$$

: multi exp. form with $E_{\alpha} \in \mathbb{R}$

(: same form as for stable particle)

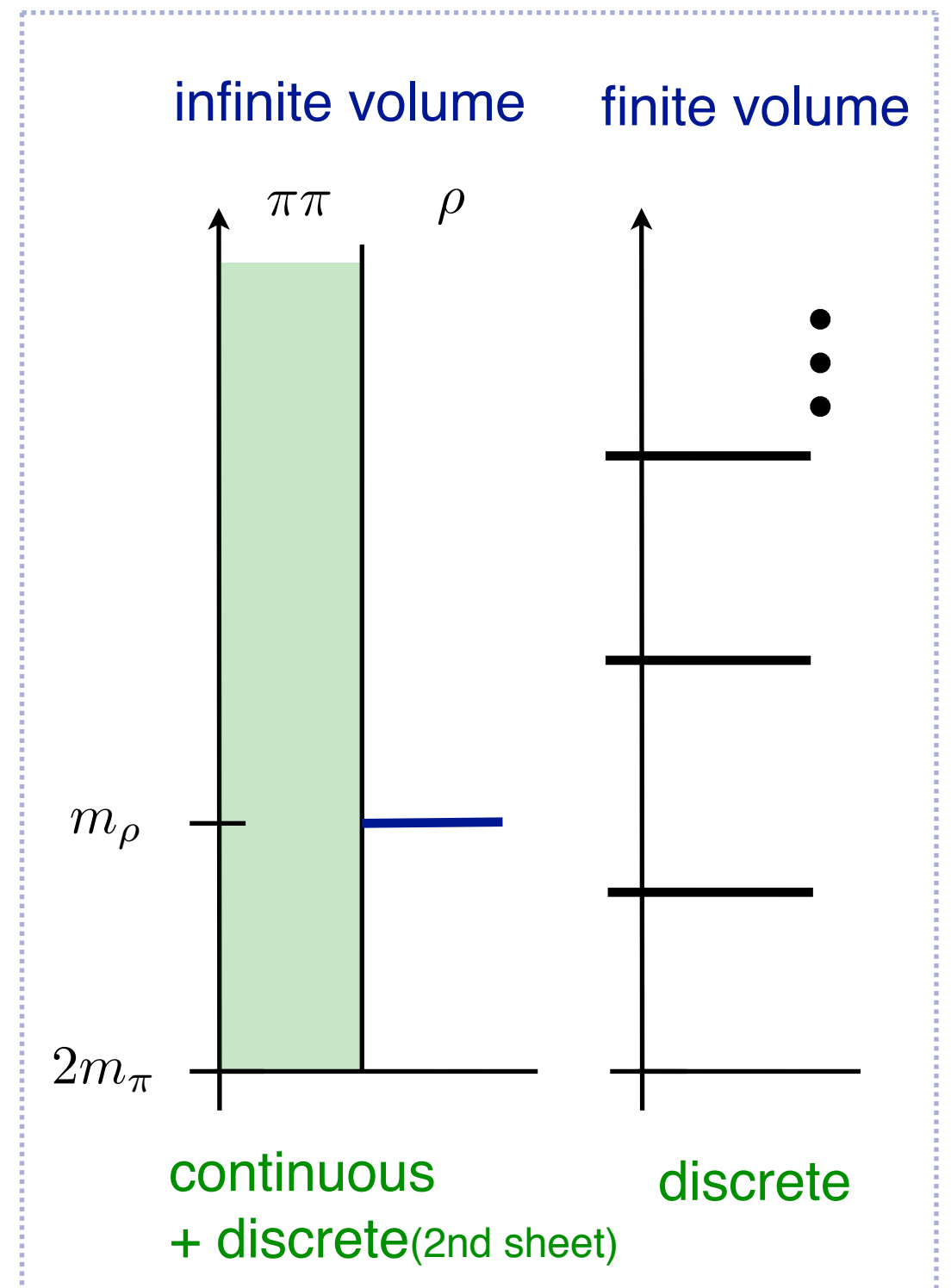
naive expectation :

for unstable particle

$$E = M + i\Gamma \quad (M, \Gamma \in \mathbb{R})$$

$$\text{time correlator} \sim e^{-Mt} \cdot e^{-i\Gamma t}$$

: ***This is only true
in infinite volume limit !***



Decay width can not be directly obtained
from correlation function on the lattice.

Approach from momentum space

ex) ρ meson decay

Lattice calc :

$$\langle 0 | \pi(p_1) \pi(p_2) \pi(p_3) \pi(p_4) | 0 \rangle \sim \sin \delta(k) e^{i\delta(k)}$$

$$\delta(k) \rightarrow \Gamma \quad \text{---} \quad ???$$

Wrong !!

Lattice QCD is formulated on the Euclidian space-time.

p_j : Euclid mom. on the lattice

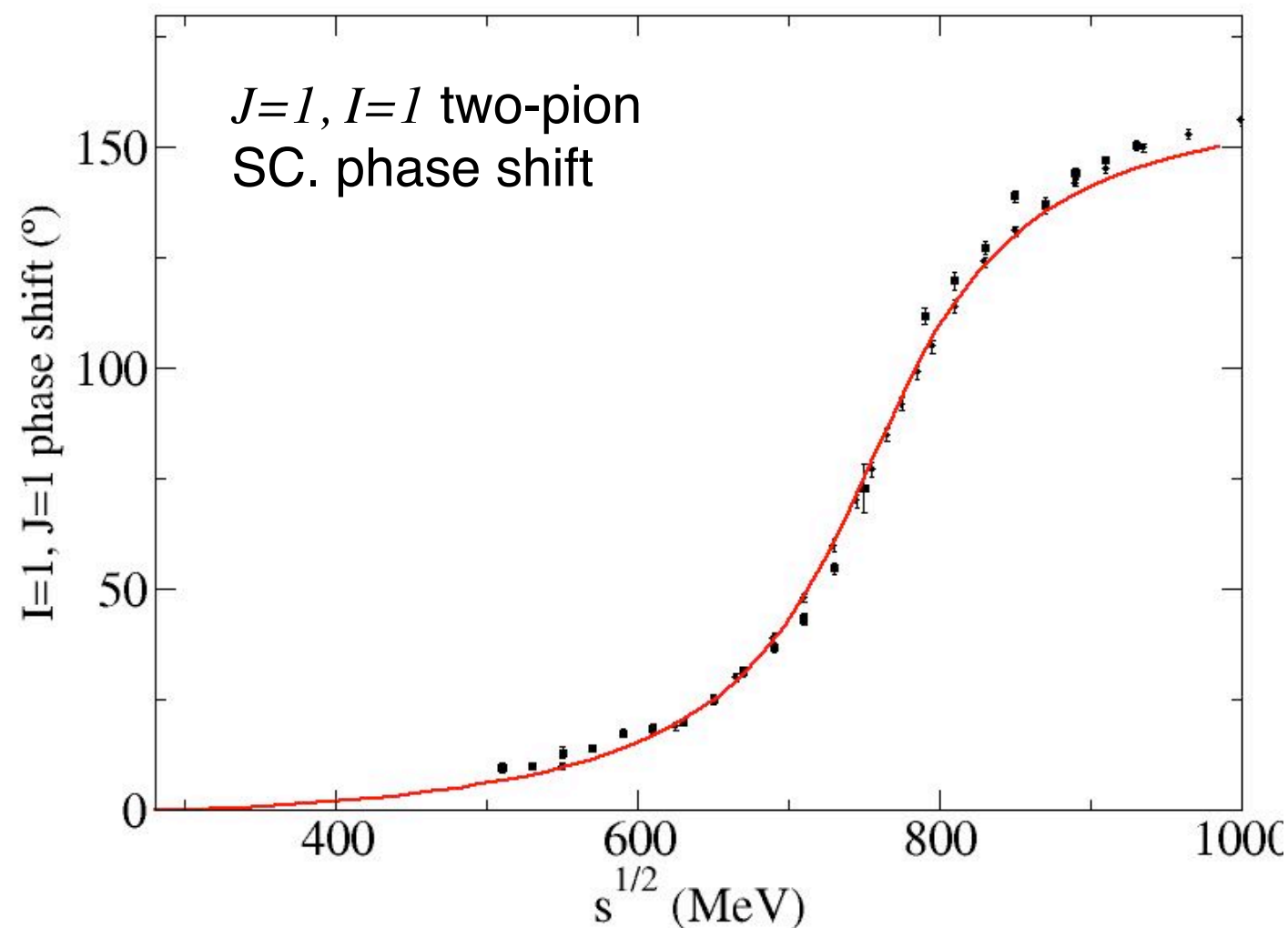
$\delta(k)$: un-physical

Is study of resonance from lattice QCD possible ?

YES !

*We can calculate scattering phase shift
by **finite size method**,
and obtain information of resonance from it.*

ex) ρ meson



➔ existence of ρ

$$m_{\rho} = 775.7^{+4.3}_{-3.3} \text{ MeV}$$

$$\Gamma_{\rho} = 135.5^{+8.0}_{-9.0} \text{ MeV}$$

This is also possible
in the lattice QCD !!

Ex) for $\pi\pi$ system

In $L \times L \times L$ periodic box (: lattice)

Energy of π :

$$E = \sqrt{m_\pi^2 + k^2} \quad k^2 = (2\pi/L)^2 \cdot n, \quad \underline{n \in \mathbb{Z}}$$

Energy of $\pi\pi$:

$$E = 2\sqrt{m_\pi^2 + k^2} \quad k^2 = (2\pi/L)^2 \cdot n, \quad \underline{n \notin \mathbb{Z}} \quad (: \text{discrete})$$

$$\frac{1}{\tan \delta_0(k)} = \frac{4\pi}{k} \cdot \frac{1}{L^3} \sum_{\mathbf{n} \in \mathbb{Z}^3} \frac{1}{p_n^2 - k^2} \quad (\mathbf{p}_n = \mathbf{n} \cdot (2\pi)/L)$$

: SC. phase shift in infinite volume

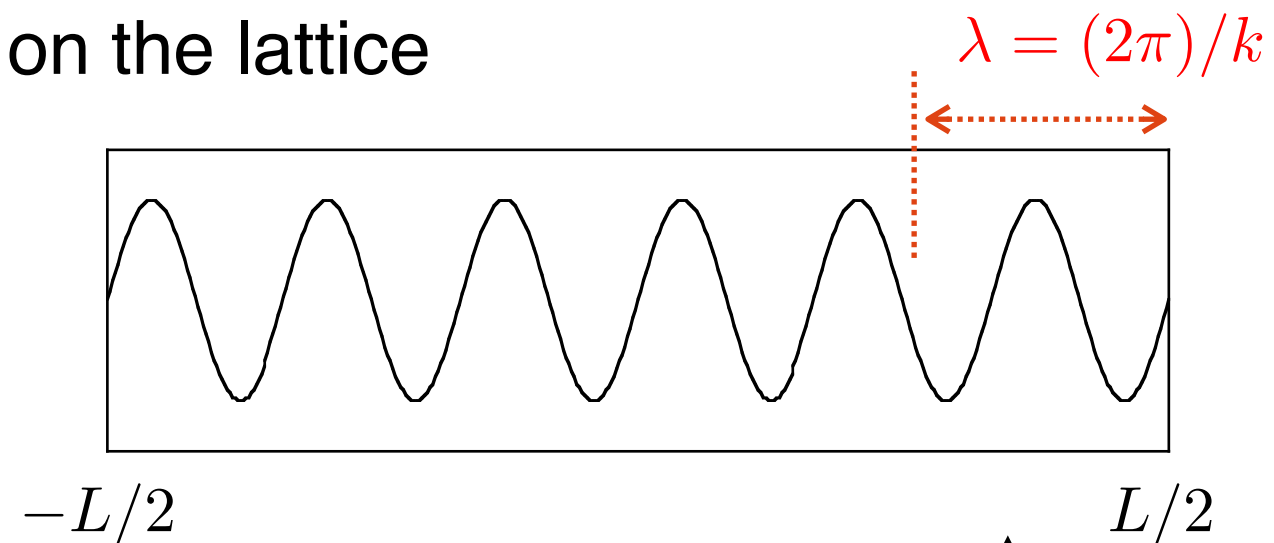
: Lüscher's formula

Energy of two-pion on the lattice $\implies$ SC. phase shift in infinite volume $\implies$ Resonance params.

Physical meaning of Lüscher's formula

for 1 dim case

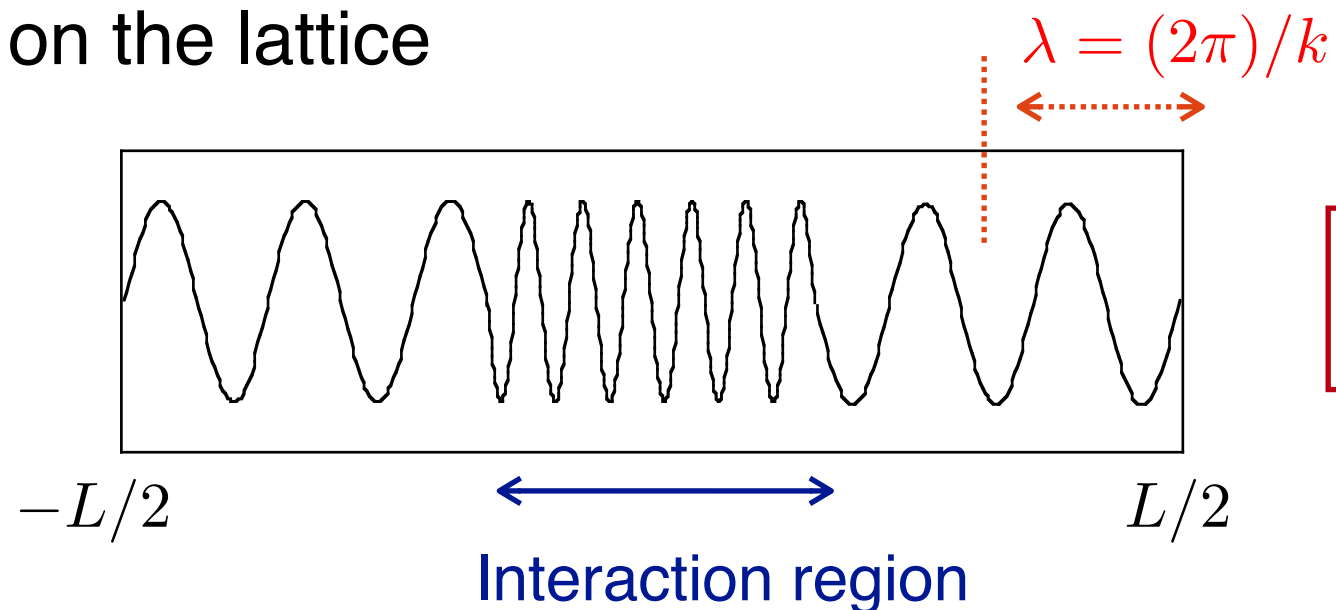
No interacting two-pion
on the lattice



$$kL = 2\pi \cdot n \quad (n \in \mathbb{Z})$$

phase diff. = SC. phase shift : $\delta(k)$

Interacting two-pion
on the lattice



$$kL + \delta(k) = 2\pi \cdot n \quad (n \in \mathbb{Z})$$

$\delta(k)$: phase shift

Other possible methods :

1. From time correlation function

Extraction the resonance parameter

from a time correlation function by using an effective theory.

: strongly depends on an effective theory.

Recently : model indep. method at $L = \text{huge}$

U.-G. Meissner et.al, NPB846(2011)1

2. From spectrum density V. Bernard et.al., JHEP 08(2008) 024

Calculations of energies on many lattice volumes
for very higher states are necessary.

: impossible in QCD !!

3. From “potential” extracted from BS-function

Solving Shrodinger eq.

HAL collaboration 2012

--> existence of resonance

Study of a resonance with these methods
has not been succeeded yet.

2. Previous works of ρ meson decay

$$\begin{aligned} N_f = 2 & : u, d \quad (m_u = m_d) \\ N_f = 2 + 1 & : u, d, s \quad (m_u = m_d < m_s) \end{aligned}$$

$$\mathcal{L} = g_{\rho\pi\pi} \cdot \epsilon^{abc} \rho_\mu^a \pi^b \overleftrightarrow{\partial}_\mu \pi^c$$

$$\text{exp. : } g_{\rho\pi\pi} = 5.878(22)$$

$$m_\rho = 775.5(0.4) \text{ MeV}$$

$$\Gamma = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{k_\rho^3}{m_\rho^2} = 146.4(1.1) \text{ MeV}$$

1) CP-PACS [PRD76\(07\)094506](#).

$$N_f = 2 \text{ (Wilson)}, a = 0.21 \text{ fm}, L = 2.5 \text{ fm}, m_\pi = 330 \text{ MeV}$$

$$g_{\rho\pi\pi} = 6.25(67)$$

2) M. Göckeler, R. Horsley, Y. Nakamura et.al. (QCDSF) [arXiv:0810.5337 \(Lat08\)](#)

$$N_f = 2 \text{ (Wilson)}, a = 0.072 - 0.084 \text{ fm}, m_\pi = 240 - 810 \text{ MeV}$$

$$g_{\rho\pi\pi} = 5.3_{-1.5}^{+2.1}$$

3) X. Feng, K. Jansen, D.B. Renner (ETMC) [arXiv:0910.4871 \(Lat09\)](#)

$$N_f = 2 \text{ (tmQCD)}, a = 0.086 \text{ fm}, L = 2.1 \text{ fm}, m_\pi = 391 \text{ MeV}$$

$$g_{\rho\pi\pi} = 6.16(48)$$

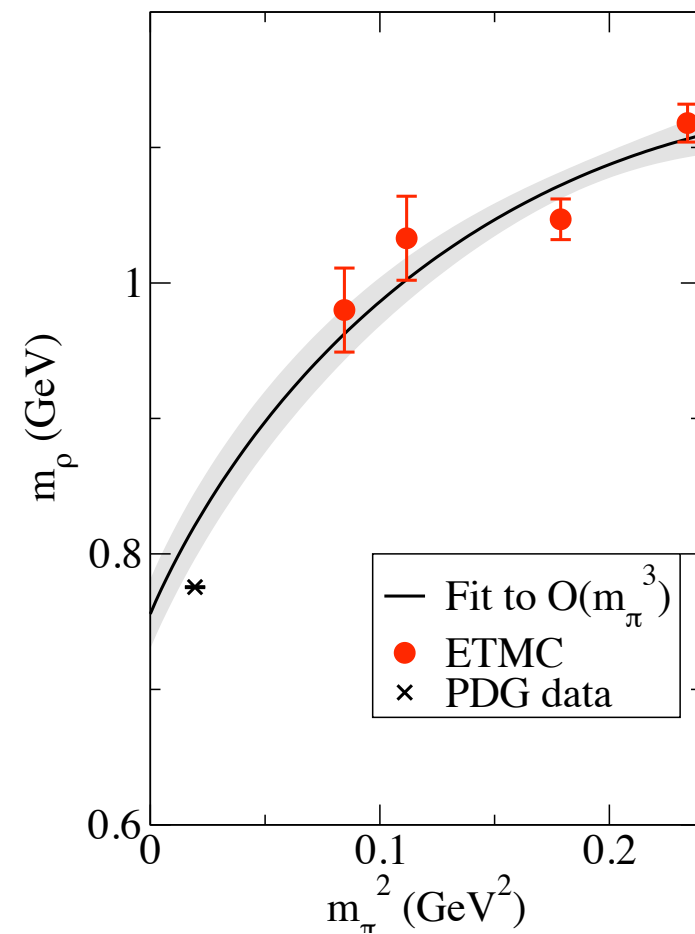
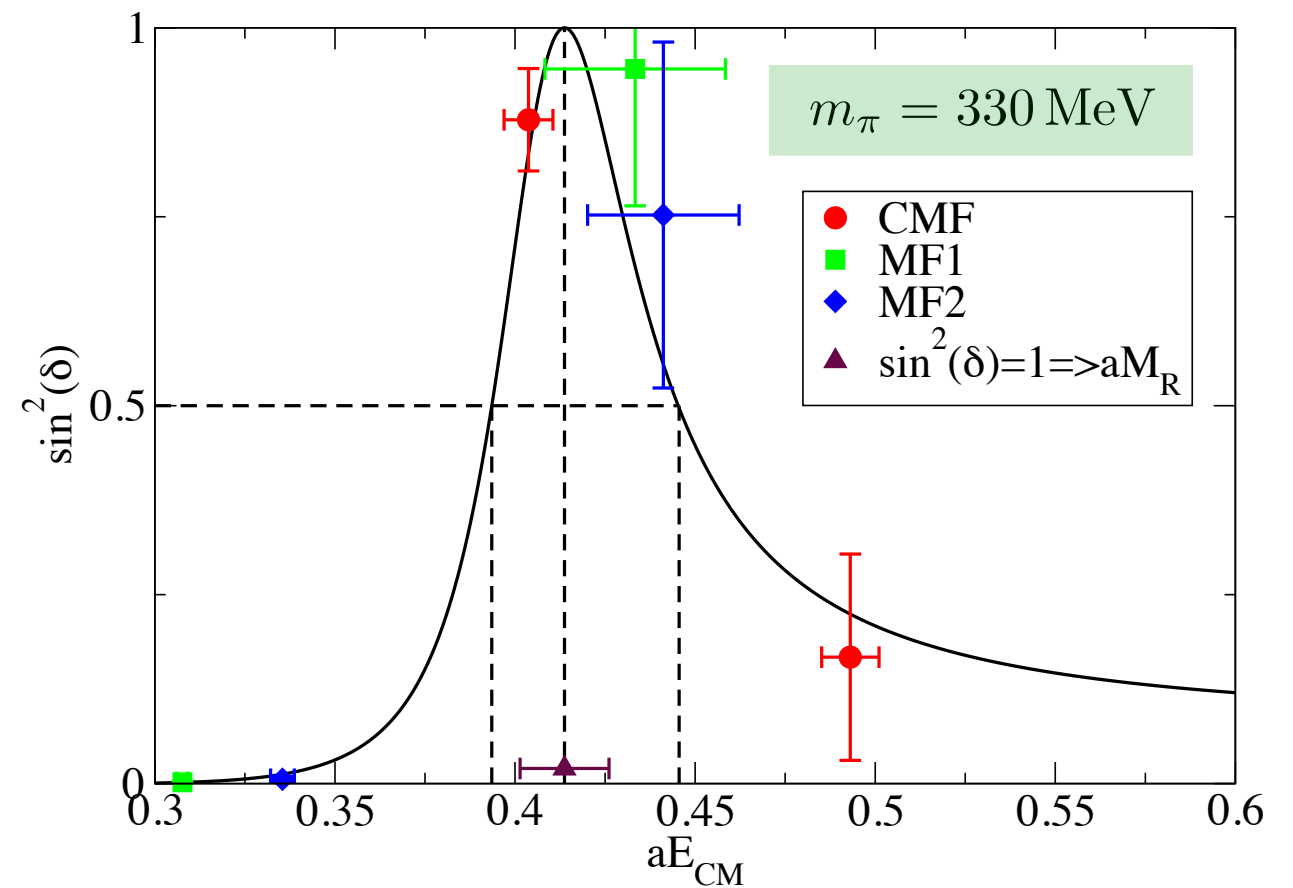
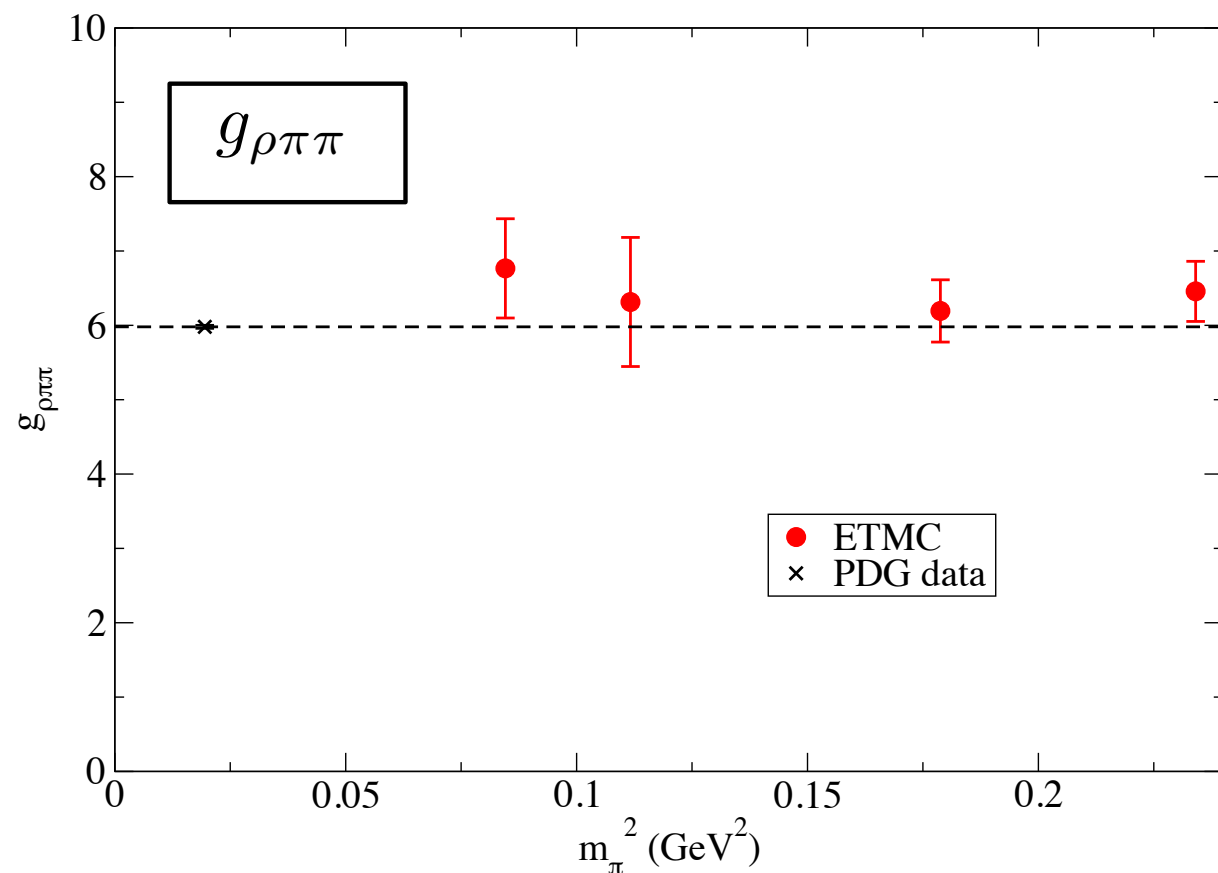
4) ETMC Lat10 (arXiv:1011.4288). PRD83(2011)094505.

$N_f = 2$ (tmQCD) $a = 0.079$ fm
(Iso-spin sym. is broken : $\rho^+ \neq \rho^0$)

$L = 2.5$ fm , $m_\pi = 290, 330$ MeV

$L = 1.9$ fm , $m_\pi = 420, 480$ MeV

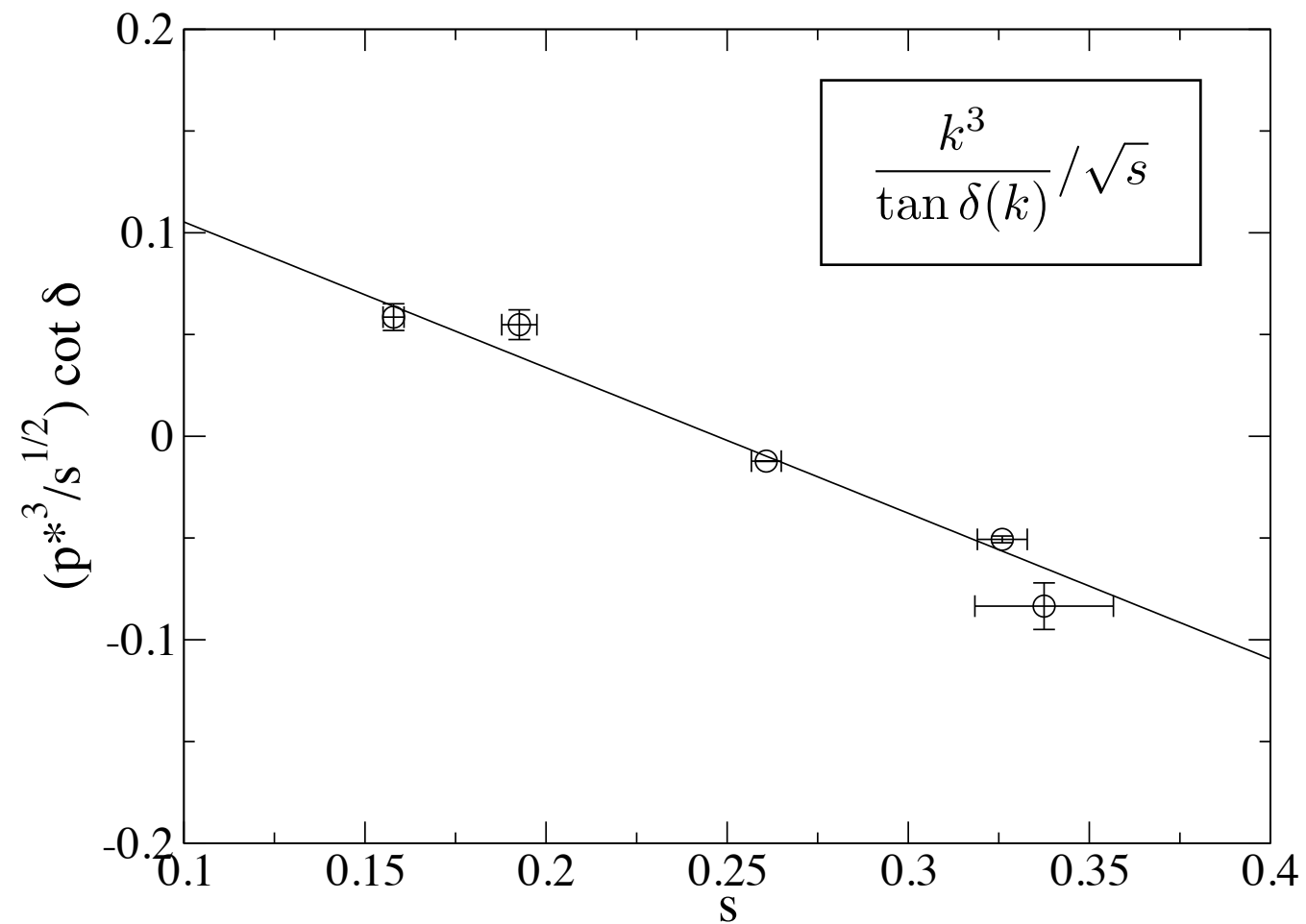
exp. : $g_{\rho\pi\pi} = 5.878(14)$
 $m_\rho = 775.49(34)$ MeV



5) C.B. Lang et.al. PRD84(2011)054503.

$N_f = 2$ (Wilson) $a = 0.124$ fm

$L = 2$ fm, $m_\pi = 266$ MeV



exp. : $g_{\rho\pi\pi} = 5.878(14)$
 $m_\rho = 775.49(34)$ MeV

$$\frac{k^3}{\tan \delta(k) \sqrt{s}} = \frac{6\pi}{g_{\rho\pi\pi}^2} \cdot (m_\rho^2 - s)$$

$g_{\rho\pi\pi} = 5.13(20)$
 $m_\rho = 792(7)(8)$ MeV
 at $m_\pi = 266$ MeV

2. Results of ρ meson decay (2009 - 2011)

PRD84(2011)094505.

Gauge conf. :

$N_f = 2 + 1$ (Wilson) , Iwasaki Gauge action

$a = 0.091$ fm , $L = 2.9$ fm

$m_\pi = 300, 410$ MeV

generated by PACS-CS col.
PRD79(2009)034503.

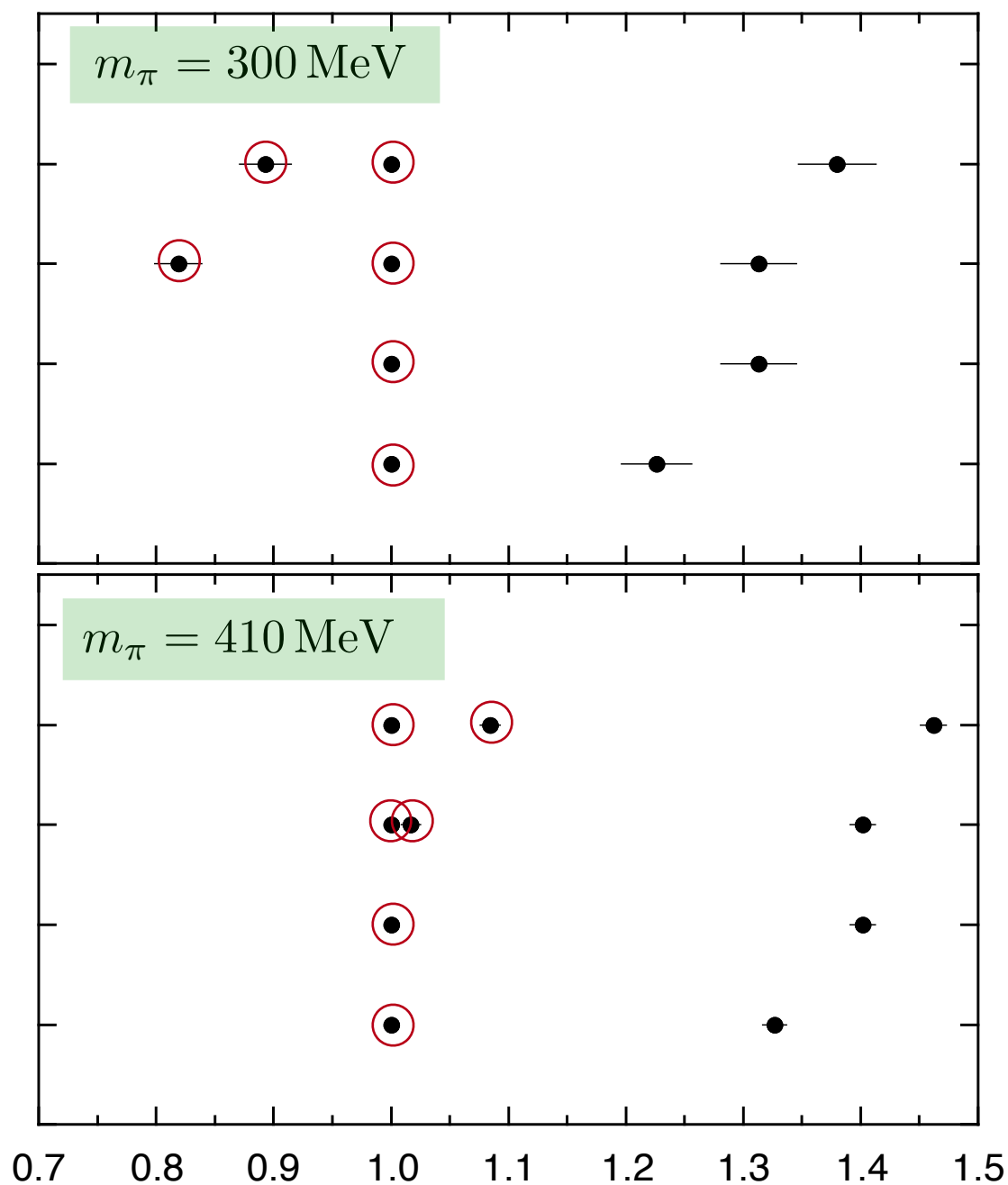
Computer : PACS-CS

Calc. points

We consider 4 irreps. : $(0, 0, 0) \mathbf{T}_1$, $(0, 0, 1) \mathbf{E}$, $(0, 0, 1) \mathbf{A}_2$, $(1, 1, 0) \mathbf{B}_1$

on 3 momentuma frames with $\mathbf{P}_{\text{tot}} = (0, 0, 0)$, $(0, 0, 1)$, $(1, 1, 0)$

$\sqrt{s}/m_\rho$ (ignoring particle int.)



	ground	1st. Ex.
$(1, 1, 0) \mathbf{B}_1$	$\pi(1, 1, 0)\pi(0, 0, 0)$	$\rho_{1+2}(1, 1, 0)$
$(0, 0, 1) \mathbf{A}_2$	$\pi(0, 0, 1)\pi(0, 0, 0)$	$\rho_3(0, 0, 1)$
$(0, 0, 1) \mathbf{E}$	$\rho_{1,2}(0, 0, 1)$	$\pi(1, 0, 1)\pi(-1, 0, 0)$
$(0, 0, 0) \mathbf{T}_1$	$\rho_{1,2,3}(0, 0, 0)$	$\pi(0, 0, 1)\pi(0, 0, -1)$

	ground	1st. Ex.
$(1, 1, 0) \mathbf{B}_1$	$\rho_{1+2}(1, 1, 0)$	$\pi(1, 1, 0)\pi(0, 0, 0)$
$(0, 0, 1) \mathbf{A}_2$	$\rho_3(0, 0, 1)$	$\pi(0, 0, 1)\pi(0, 0, 0)$
$(0, 0, 1) \mathbf{E}$	$\rho_{1,2}(0, 0, 1)$	$\pi(1, 0, 1)\pi(-1, 0, 0)$
$(0, 0, 0) \mathbf{T}_1$	$\rho_{1,2,3}(0, 0, 0)$	$\pi(0, 0, 1)\pi(0, 0, -1)$

Finite size formula

M. Lüscher, NPB354(1991)531.

K.Rummainen and S.Gottlib, NPB450(1995)397.

ETMC, arXiv:1011.5288.

Relation between $\delta(k)$ and E

$$\frac{1}{\tan \delta(k)} = \begin{cases} V_{00} & \text{for } (0, 0, 0) \mathbf{T}_1 \\ V_{00} - V_{20} & \text{for } (0, 0, 1) \mathbf{E} \\ V_{00} + 2 \cdot V_{20} & \text{for } (0, 0, 1) \mathbf{A}_2 \\ V_{00} - V_{20} + \sqrt{6} \cdot V_{22} & \text{for } (1, 1, 0) \mathbf{B}_1 \end{cases}$$

$$V_{lm}(k; \mathbf{P}_{\text{tot}}) = \frac{1}{\gamma q^{l+1}} \frac{1}{\pi^{3/2} \sqrt{2l+1}} e^{-im\pi/4} \cdot Z_{lm}(1; q; \mathbf{d}) \quad (: \text{Real})$$

$$\sqrt{E^2 - P_{\text{tot}}^2} = \sqrt{s} = 2\sqrt{m_\pi^2 + k^2}$$

$$q = kL/(2\pi) \quad \mathbf{d} = \mathbf{P}_{\text{tot}} L/(2\pi)$$

$$Z_{lm}(s; q; \mathbf{d}) = \sum_{\mathbf{r} \in D(\mathbf{d})} \mathcal{Y}_{lm}(\mathbf{r}) \cdot (|\mathbf{r}|^2 - q^2)^{-s}$$

$$D(\mathbf{d}) = \{ \mathbf{r} \mid \mathbf{r} = \hat{\gamma}^{-1}(\mathbf{n} + \mathbf{d}/2) , \mathbf{n} \in \mathbb{Z}^3 \}$$

$$\hat{\gamma}^{-1} : \text{Inv. Lorentz boost with } \gamma = E/\sqrt{s}$$

For $(0, 0, 1) \mathbf{A}_2$ and $(1, 1, 0) \mathbf{B}_1$

For $(0, 0, 0) \mathbf{T}_1$ and $(0, 0, 1) \mathbf{E}$
energies of ground state are extracted
from correlation function of ρ meson
as usual calculations of hadron masses.

Energies are extracted by variational method.

$$\mathcal{O}_1(t) = \rho(\mathbf{p}, t)$$

$$\mathcal{O}_2(t) = \pi^-(\mathbf{p}, t)\pi^+(\mathbf{0}, t) - \pi^+(\mathbf{p}, t)\pi^-(\mathbf{0}, t)$$

for $\mathbf{p} = (0, 0, 1), (1, 1, 0)$

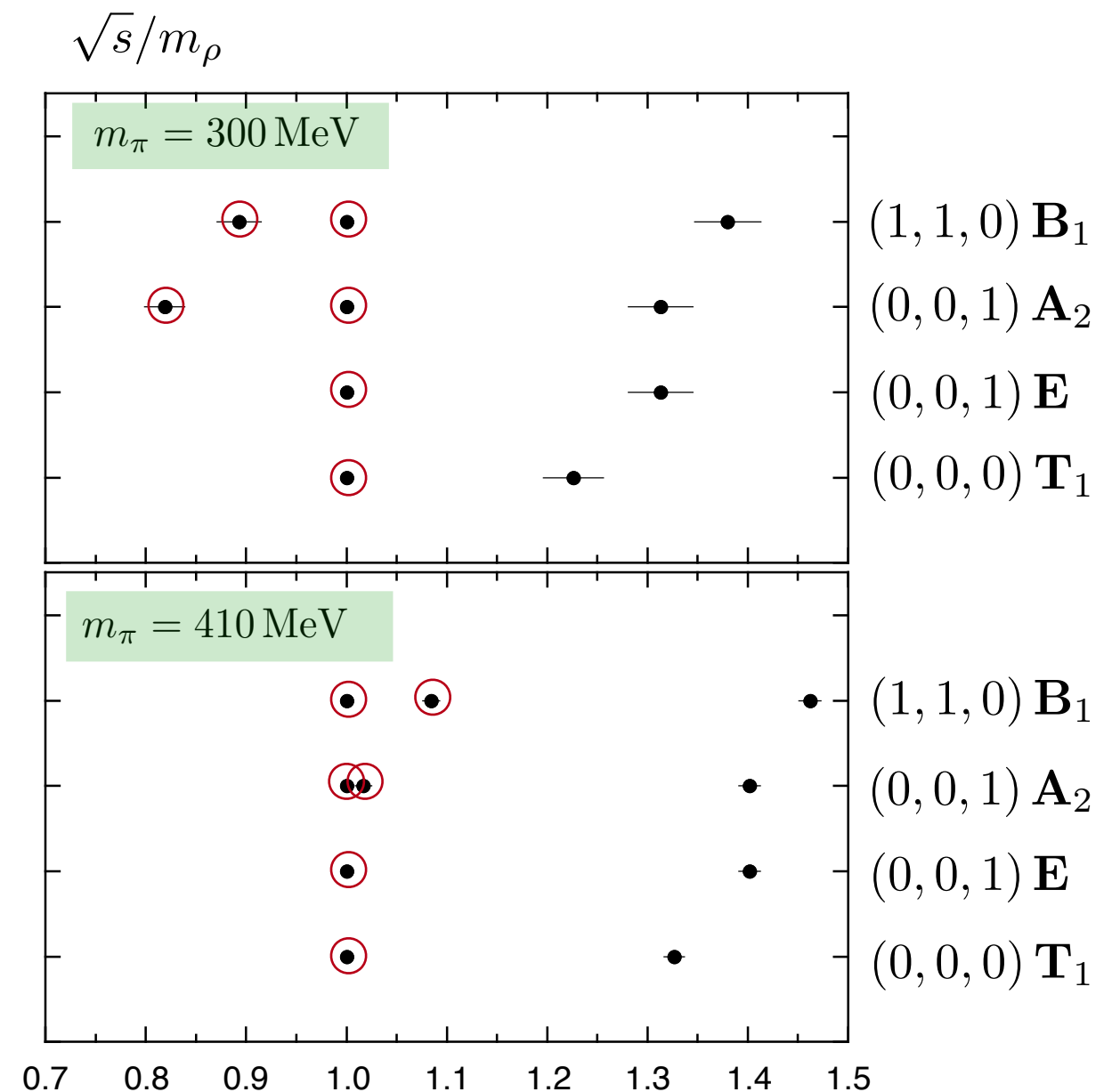
$$G_{ij}(t) = \langle \mathcal{O}_i^\dagger(t) \mathcal{O}_j(0) \rangle$$

(: 2×2 matrix)

Assuming 2 states dominant for $t \gg t_s$

$$\text{Ev}[G(t) G^{-1}(t_0)]_\alpha = e^{-W_\alpha \cdot (t-t_0)} \quad (\alpha = 1, 2)$$

W_α : energy eigenvalue



Calc. of $G(t)$

$$G_{\pi\pi\rightarrow\pi\pi}(t) =$$

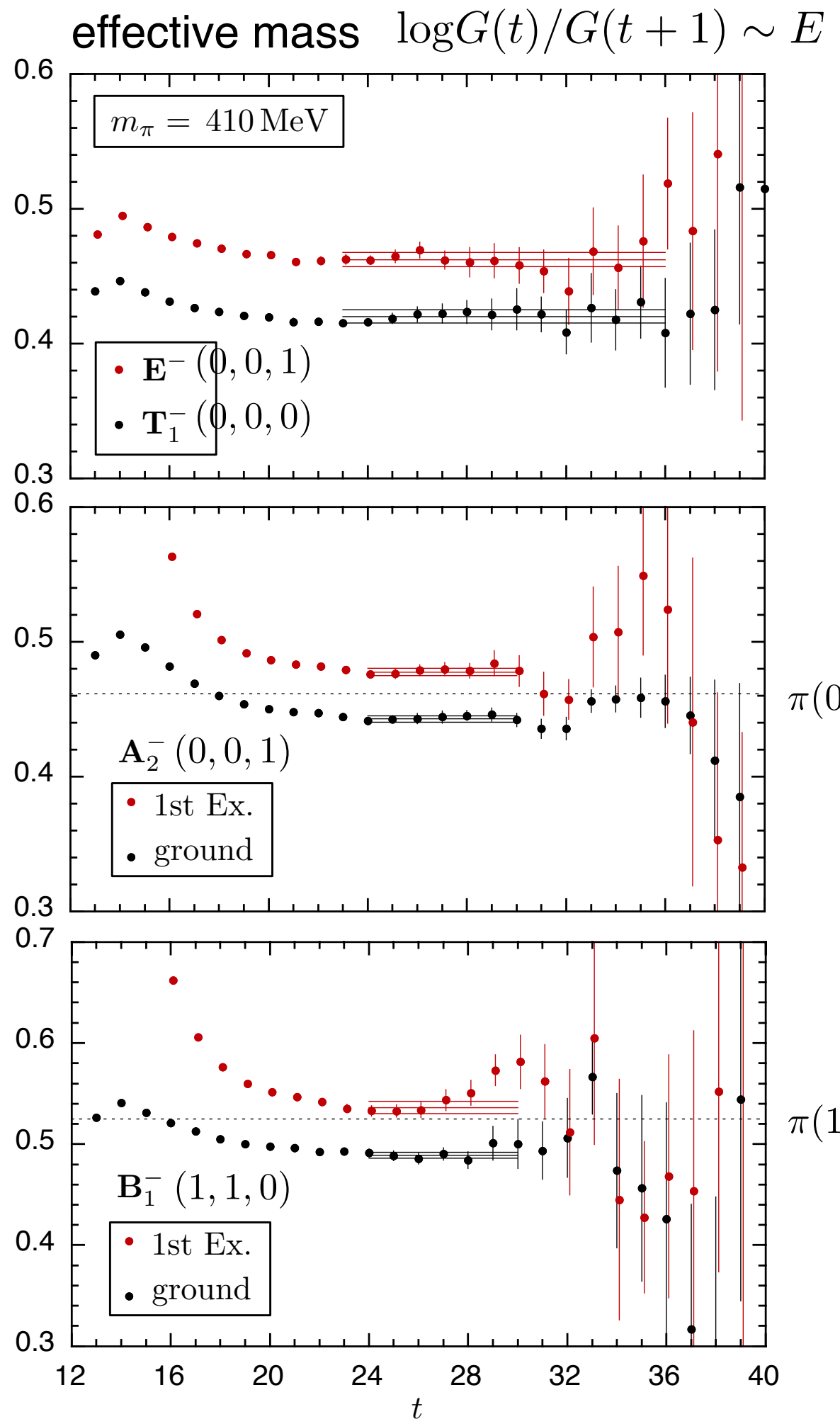
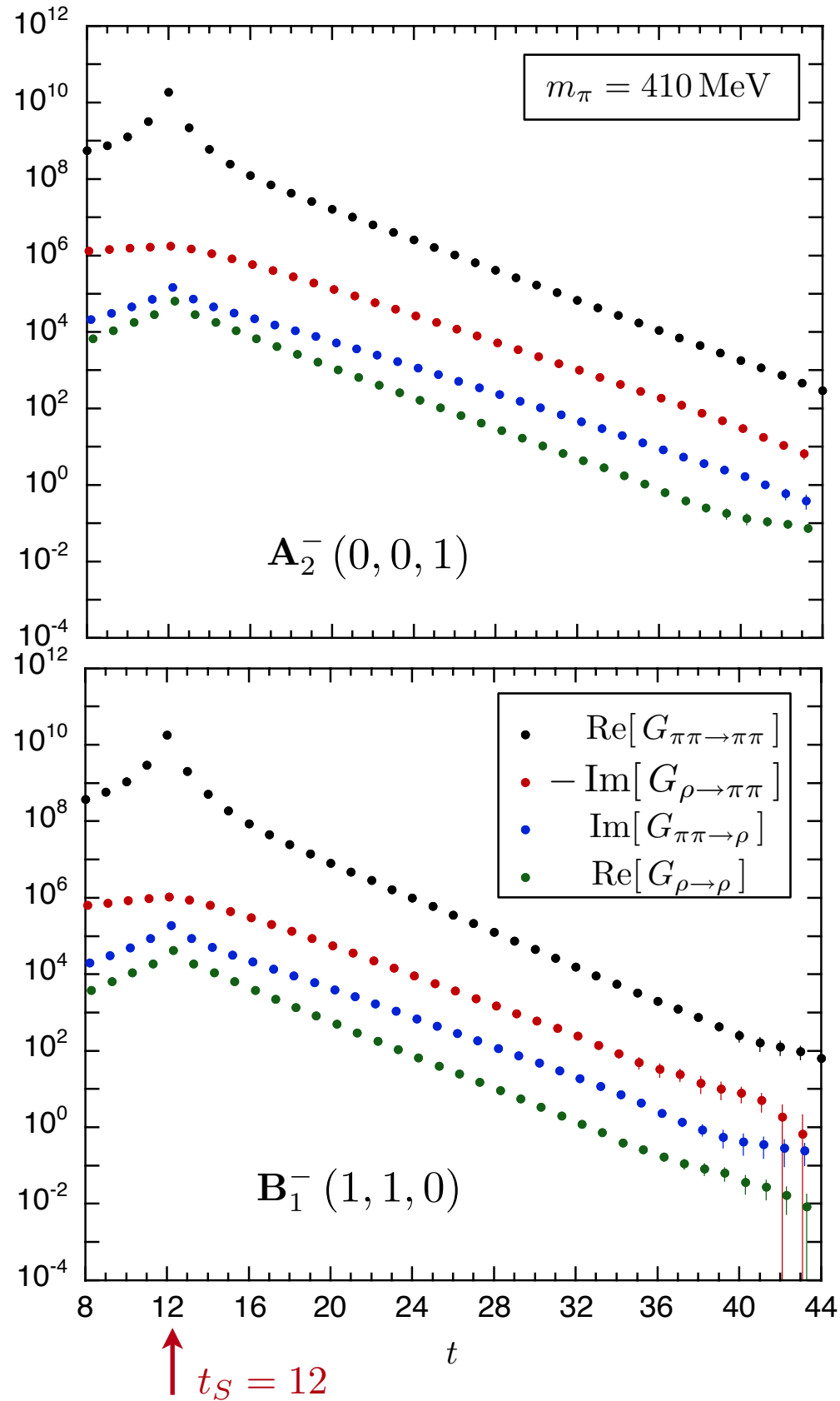
$$G_{\pi\pi\rightarrow\rho}(t) =$$

$$G_{\rho\rightarrow\pi\pi}(t) =$$

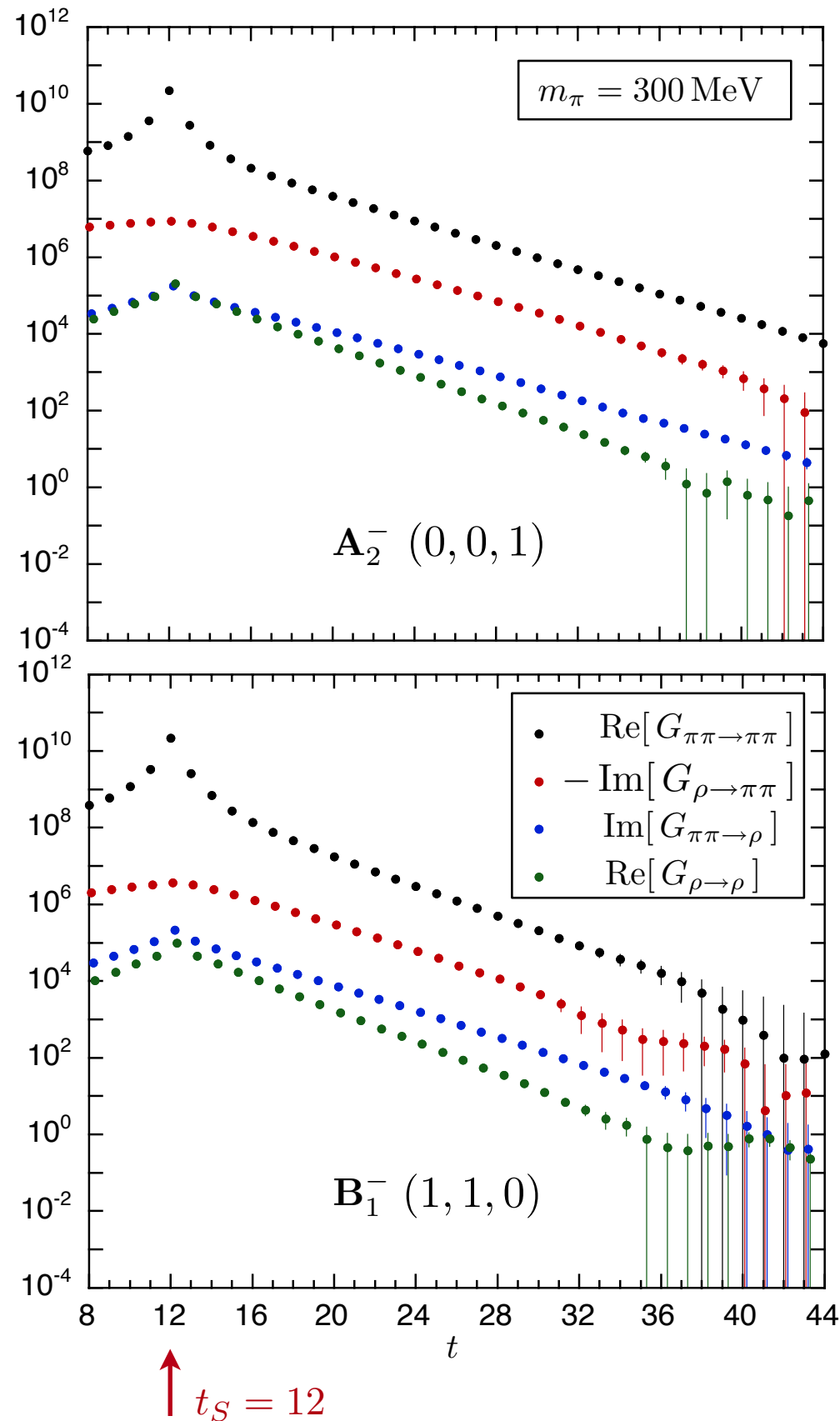
t

These are calculated
by the stochastic source and the source method.

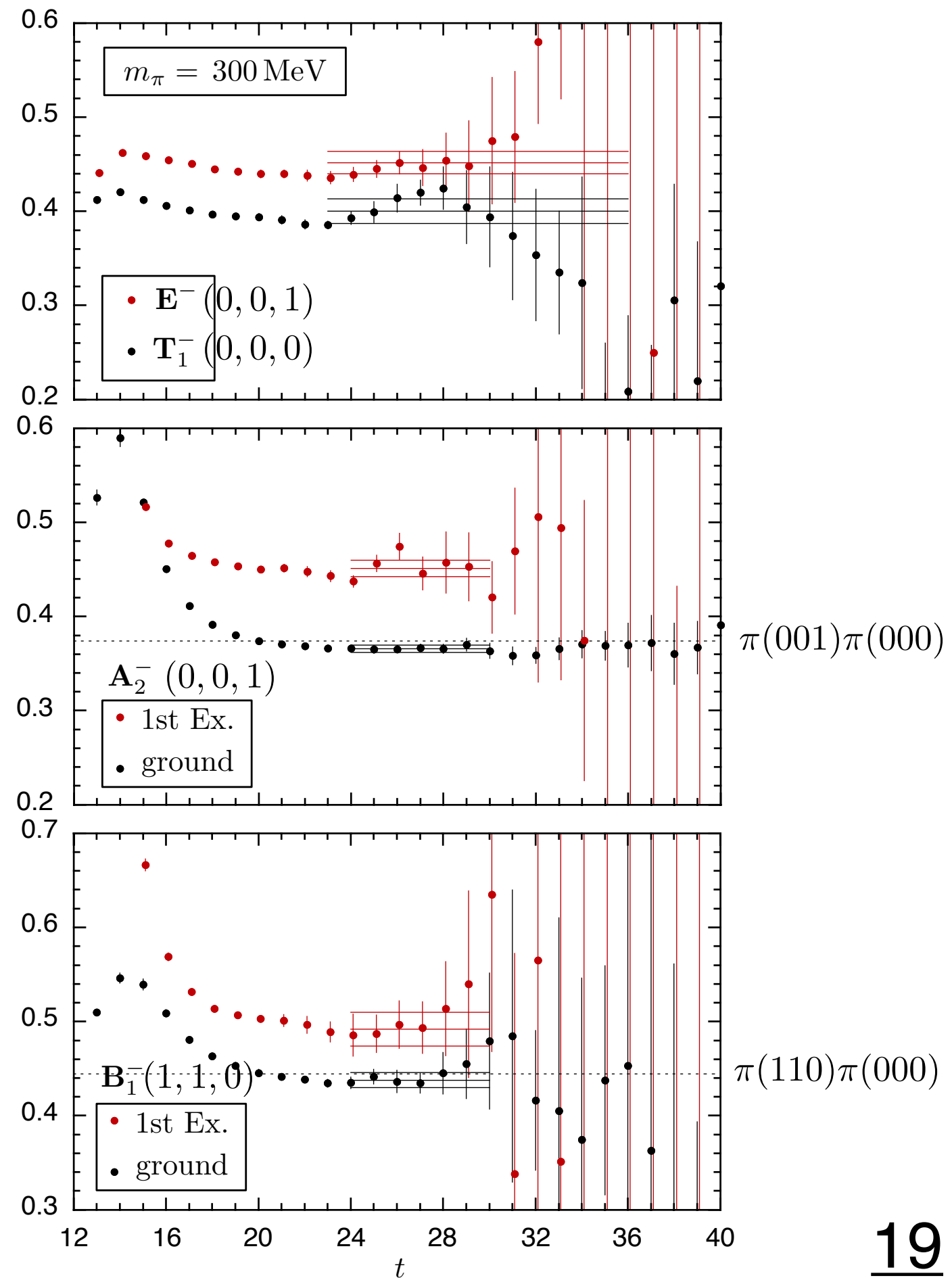
Results at $m_\pi = 410 \text{ MeV}$



Results at $m_\pi = 300 \text{ MeV}$



effective mass $\log G(t)/G(t+1) \sim E$



Results of SC. phase shift

$N_f = 2 + 1$ (Wilson) $a = 0.091$ fm

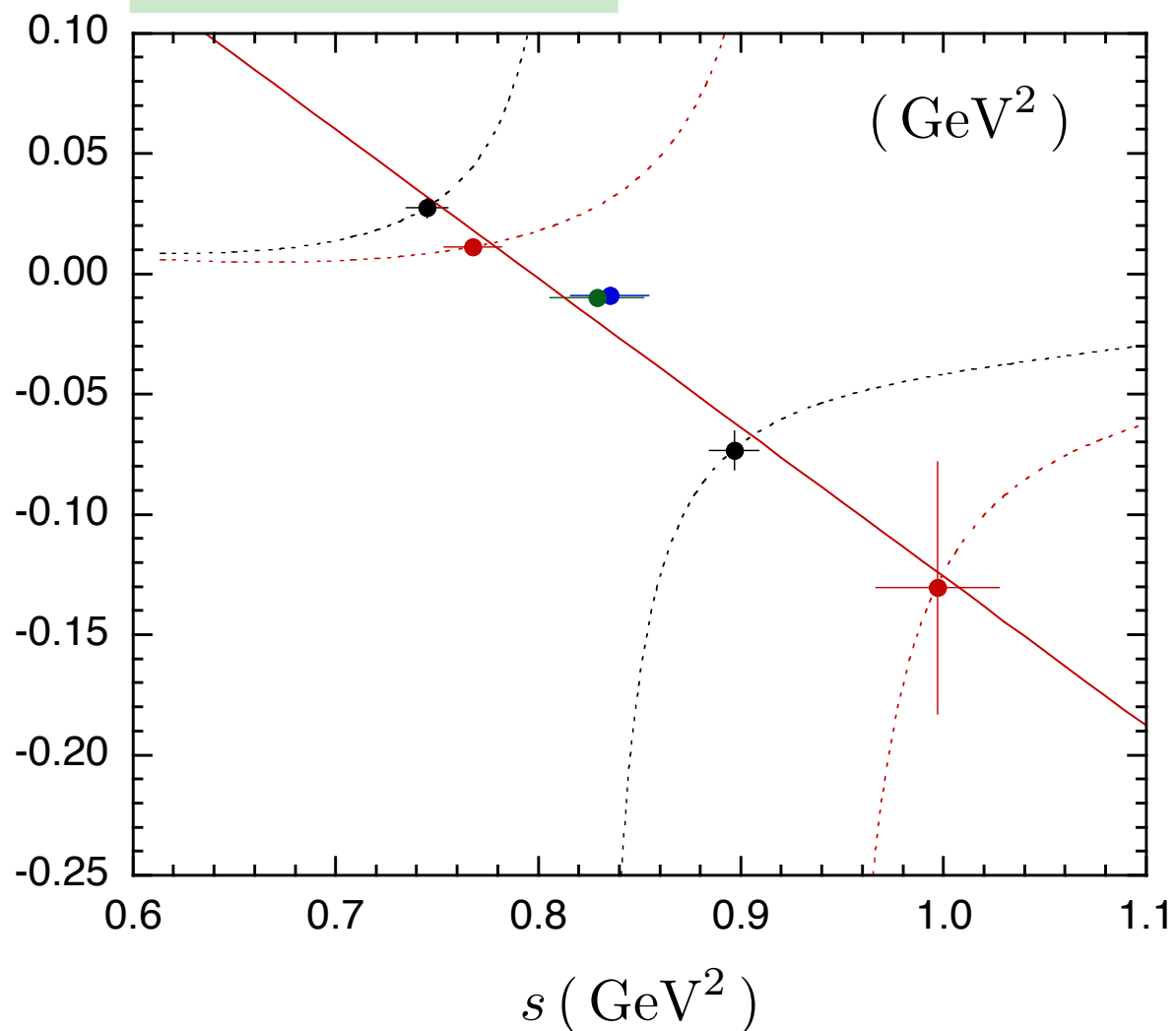
$L = 2.9$ fm , $m_\pi = 300, 410$ MeV

exp. : $g_{\rho\pi\pi} = 5.878(14)$

$m_\rho = 775.49(34)$ MeV

$$\frac{k^3}{\tan \delta(k)} / \sqrt{s} = \frac{6\pi}{g_{\rho\pi\pi}^2} \cdot (m_\rho^2 - s)$$

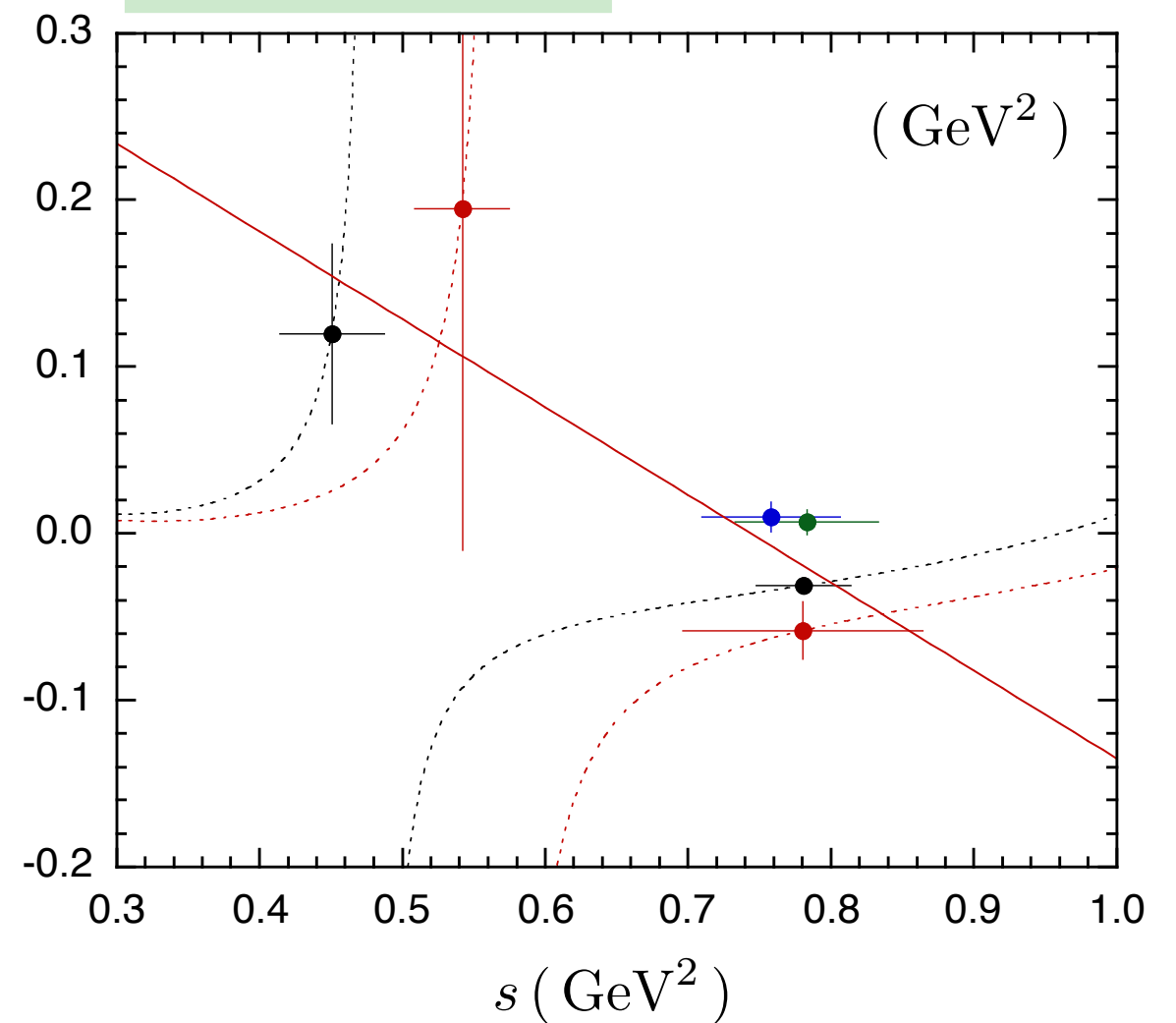
$m_\pi = 410$ MeV



$$g_{\rho\pi\pi} = 5.51 \pm 0.40$$

$$m_\rho = 892.8 \pm 5.5 \pm 13 \text{ MeV}$$

$m_\pi = 300$ MeV



$$g_{\rho\pi\pi} = 5.98 \pm 0.56$$

$$m_\rho = 863 \pm 23 \pm 12 \text{ MeV}$$

3. Comparison with other groups

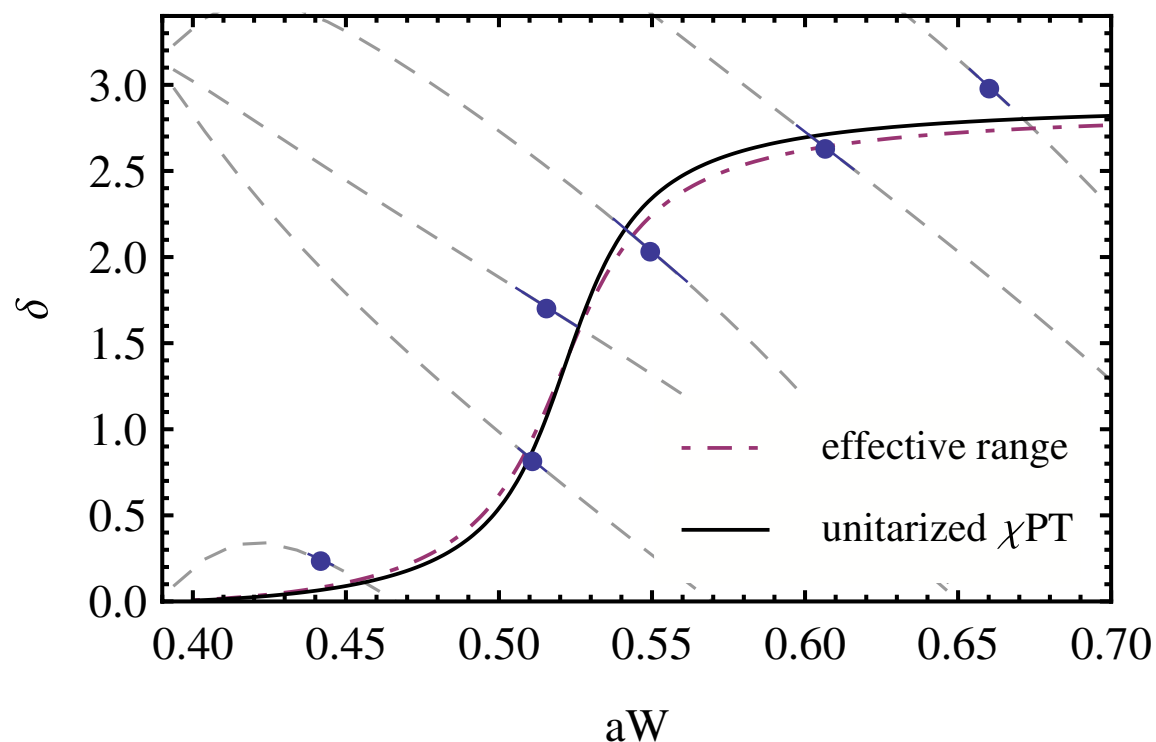
9) Pelissier et.al.

PRD87(2013)014503.

$N_f = 2$ (Wilson) $a = 0.1255$ fm

$m_\pi = 300$ MeV

$V = 24^2 \times \eta 24 \times 48$ ($\eta = 1.0, 1.25, 2.0$)



$$g_{\rho\pi\pi} = 6.67(42)$$

$$m_\rho = 827(3)(5) \text{ MeV}$$

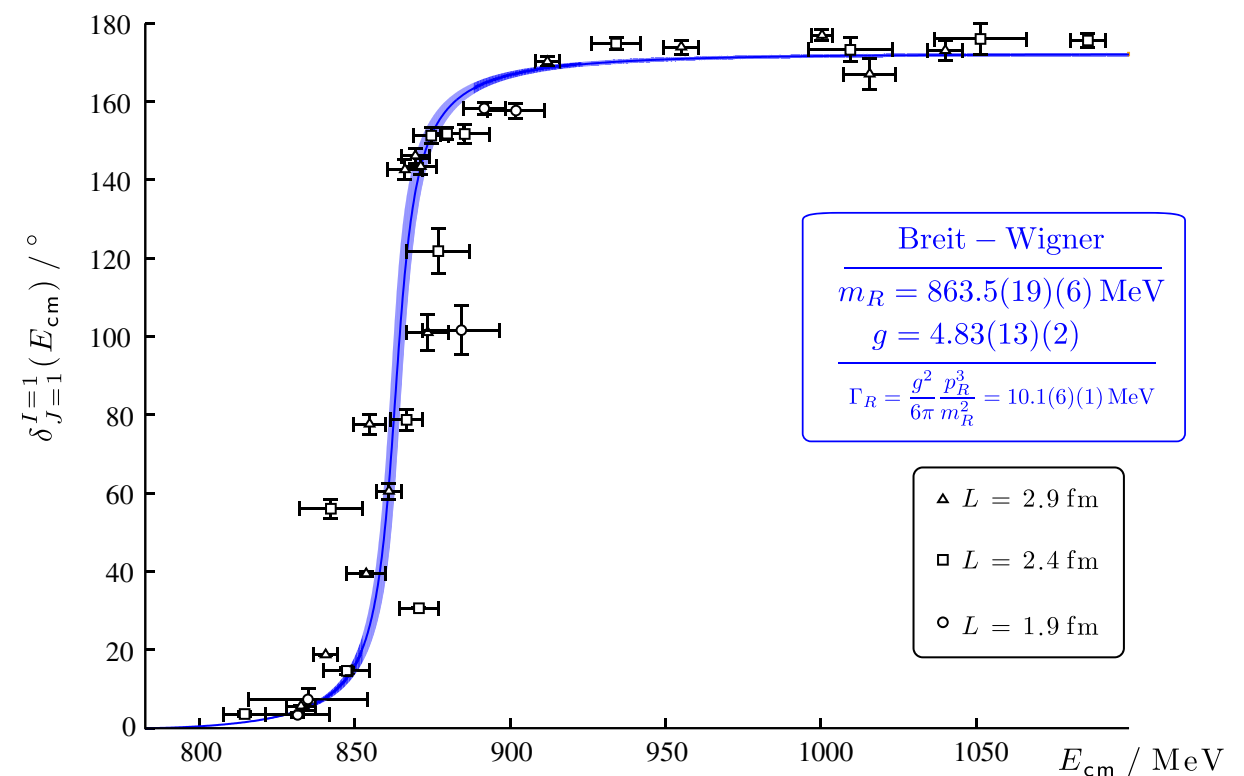
at $m_\pi = 300$ MeV

10) HSC Collaboration

PRD87(2013)034505.

$N_f = 2 + 1$ (Wilson) $a = 0.12$ fm

$L = 1.9, 2.4, 2.9$ fm, $m_\pi = 391$ MeV



$$g_{\rho\pi\pi} = 4.83(13)(2)$$

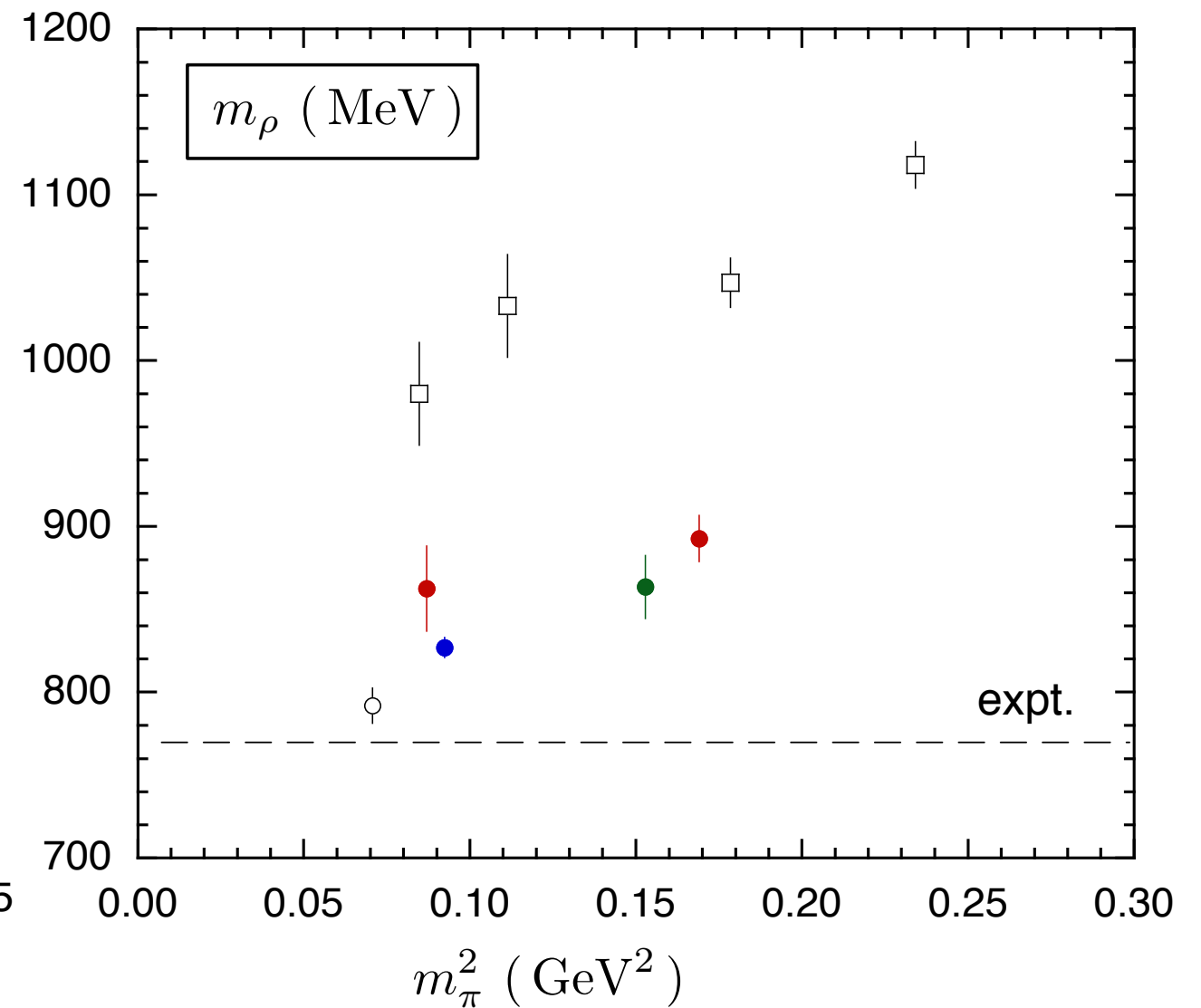
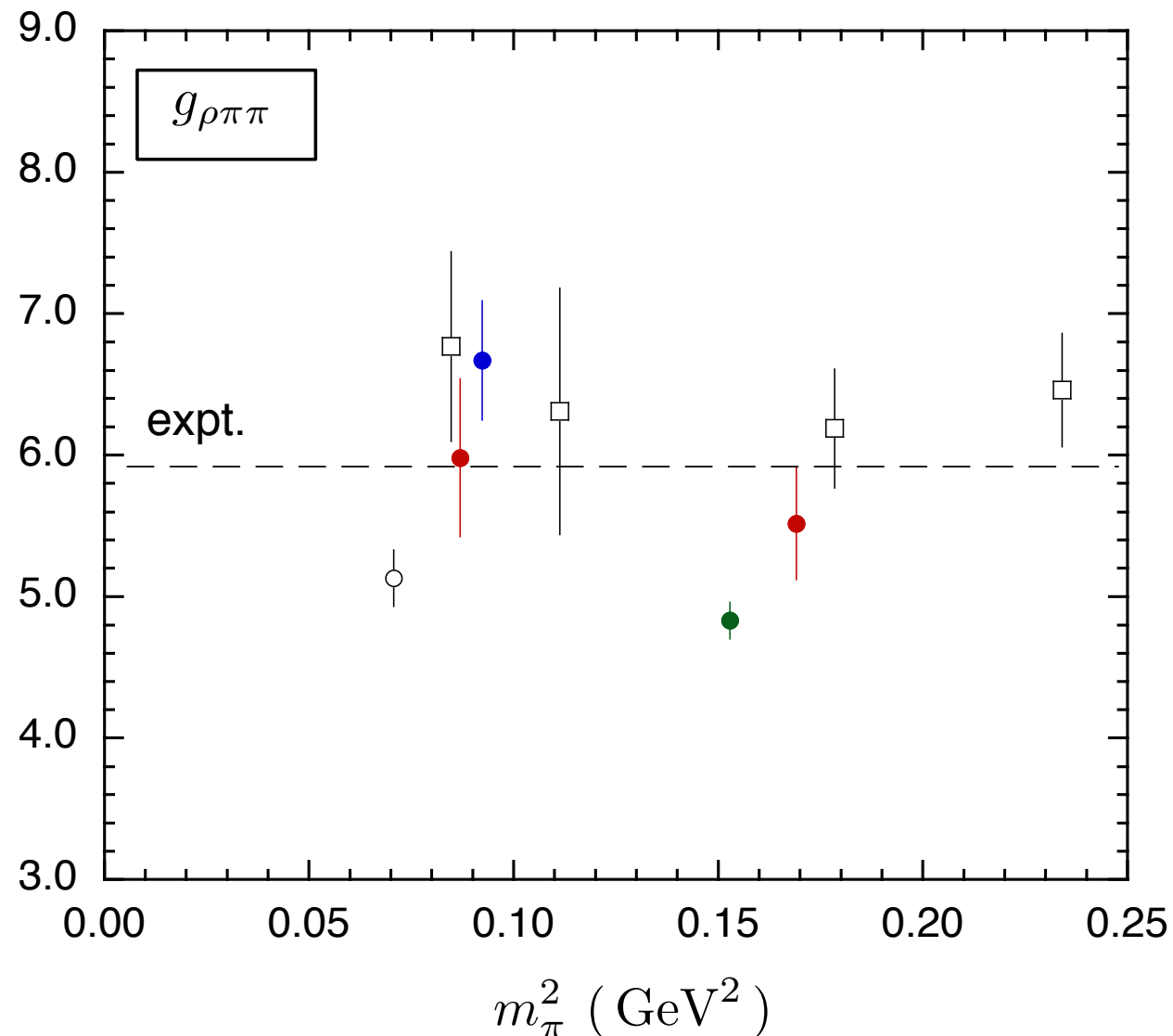
$$m_\rho = 863.5(19)(6) \text{ MeV}$$

at $m_\pi = 391$ MeV

Comparison

□ ETMC ○ Lang et.al. ● Pelissier et.al.
● PACS-CS ● HSC

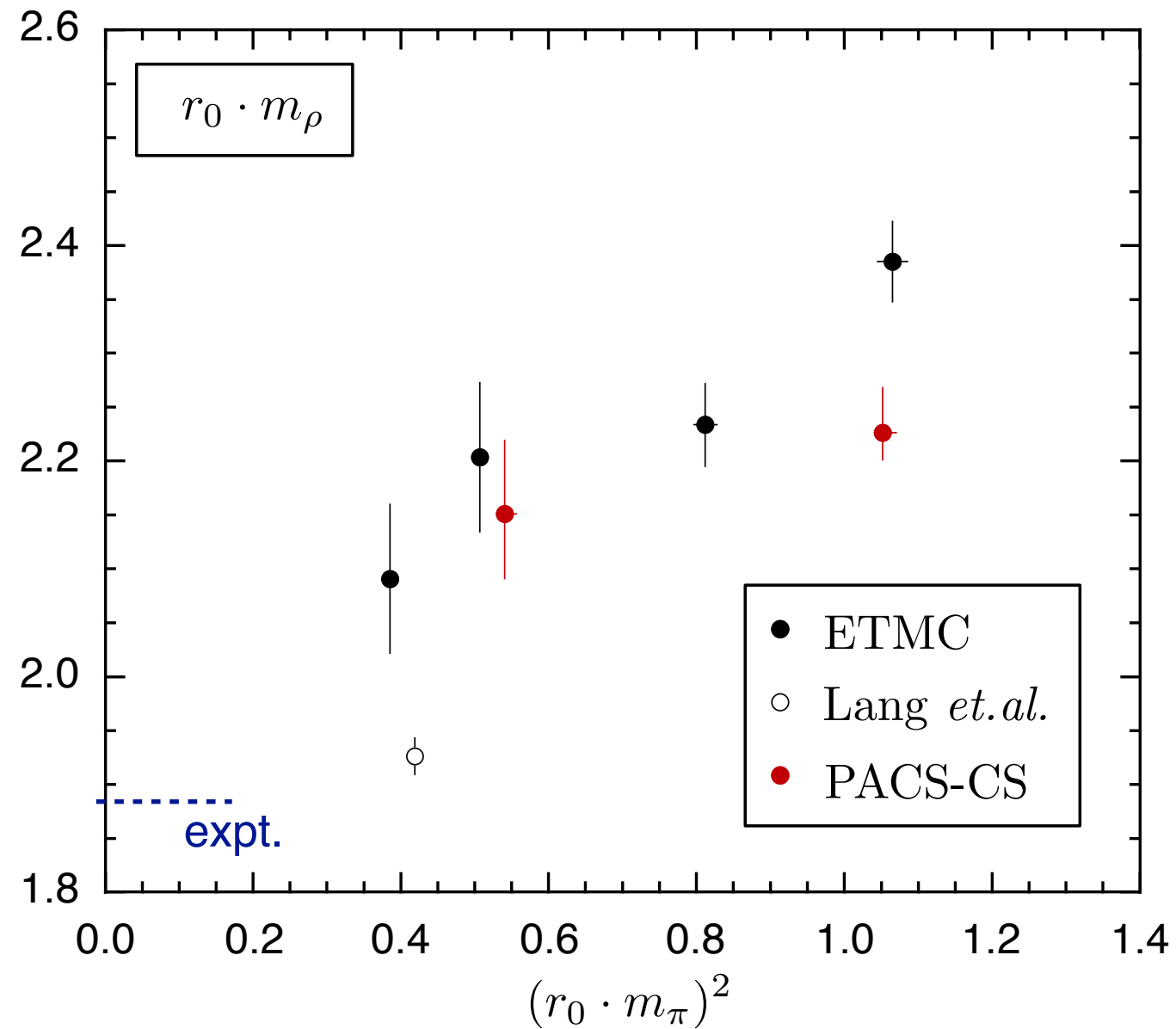
exp. : $g_{\rho\pi\pi} = 5.878(14)$
 $m_\rho = 775.49(34) \text{ MeV}$



We find discrepancies.

Precise calc. near physical quark mass and cont. limit needs !!

Comparison for dim. less values



Discrepancies still remain.

4. Summary

We calculate SC. phase shift for $I=1$ two-pion system.
We evaluate coupling constant and resonance energy from it.

exp. :

$$g_{\rho\pi\pi} = 5.878(22)$$

$$m_\rho = 775.5(0.4) \text{ MeV}$$

$$\Gamma = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{k_\rho^3}{m_\rho^2} = 146.4(1.1) \text{ MeV}$$

Our results :

- $N_f = 2 + 1$ (Wilson) $a = 0.091 \text{ fm}$
 $L = 2.9 \text{ fm}$, $m_\pi = 410 \text{ MeV}$

$$\begin{aligned} g_{\rho\pi\pi} &= 5.65 \pm 0.44 \\ m_\rho &= 892.7 \pm 5.2 \text{ MeV} \\ \Gamma_\rho &= 135 \pm 21 \text{ MeV} \\ &\text{(assuming : } g_{\rho\pi\pi} : \text{const.)} \end{aligned}$$

- $N_f = 2 + 1$ (Wilson) $a = 0.091 \text{ fm}$
 $L = 2.9 \text{ fm}$, $m_\pi = 300 \text{ MeV}$

$$\begin{aligned} g_{\rho\pi\pi} &= 5.95 \pm 0.57 \\ m_\rho &= 860 \pm 18 \text{ MeV} \\ \Gamma_\rho &= 150 \pm 28 \text{ MeV} \\ &\text{(assuming : } g_{\rho\pi\pi} : \text{const.)} \end{aligned}$$

We find :

Discrepancies among lattice calculations
for both resonance mass and effective constant.

Precise calc. near physical quark mass and cont. limit needs !!

In Feature works :

- We have to study at phys. point.

To set $\sqrt{s} \sim m_\rho$

for $m_\pi \rightarrow 0$

$\pi(\mathbf{p})\pi(-\mathbf{p})$ with $\mathbf{p}$: larger

Stat. Error. $\rightarrow$ large

1. We have to improve the operator of π with high momentum.

2. How extract energies of high excited state ?

3. Other methods ?

- We have to study other resonances.

K^*, Δ

a_0, K, σ

charm resonance ($X, Y, Z \dots$)

many works remain !!

$\sqrt{s}/m_\rho$ of $\pi\pi$ at physical point

for $\mathbf{P}_{\text{tot}} = 0$

on $L = 6$ fm

