Solving the Core-Cusp Problem of CDM Halos and the Origin of their Observational Laws

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Collaborators

Masao Mori, Tomoaki Ishiyama, Andreas Burkert, Yohei Miki, Taisuke Boku, Naohito Nakasato The Core-Cusp Problem in CDM Halos and Supernova Feedback Ogiya and Mori, 2011, ApJ, 736, L2 Ogiya and Mori, 2012, arXiv: 1206.5412 Ogiya et al., 2013, ACS, 6(3), 58-70 (Japanese)

What is the Core-Cusp Problem?



Numerical simulation is a powerful tool!!

Supernova Feedback



Mashchenko et al. (2008)

Gas Oscillation

- Cosmoligical N-body+SPH simulation
 - Supernova feedback etc.
 - Blue: gas, Yellow: star

– Gas

- Blown out (expansion)
- Fall back towards center
- Repeat many times
- Gas Mass-Loss
 - Feedback → Galactic winds (e.g., Mori+ 1997, 1999;

Mac Low & Ferrara 1999)

Motivation & Policy

- Understanding the dynamical response of DM halos to Gas Mass-Loss/Oscillation
- Previous studies (e.g, Navarro et al. 1996; Pontzen & Governato 2012)
 - Poor authenticities of numerical simulations
 - Too complicated simulations
- → Constructing an analytical model
- → N-body simulations of idealized models with sufficient authenticities

Effects of Mass-Loss

- The central cusp becomes flatter when mass-loss occur in a shorter timescale.
- But the central cusp still remains.
 - Counter-evidence against previous studies with poor authenticities



Recurring Change in Potential



1) Gas heating by supernovae

2) Gas expansion

- 3) Energy loss by radiative cooling
- 4) Contraction towards the center
- 5) Ignition of star formation again

Repetition of these processes Gas Oscillation

Change of potential \Rightarrow DM halo is affected gravitationally \Rightarrow Cusp-to-Core transformation?

Linear Analysis: Resonance Model

- Equilibrium system (0) + External force (ex) \Rightarrow Induced values (ind) •
- Focus on the particle group with $\rho_0 = const.$, $v_0 = const.$ •
- External force : $-\frac{\partial \Phi_{\text{ex}}}{\partial r} = \sum_{n} A_n \cos\left(kr - n\Omega t\right)$ ullet

(A: strength, k: wavenumber, Ω : frequency)

Linearized continuity eq. and Euler eq.

$$\rho_{ind}(t,r) = -\sum_{n} \underbrace{A_n \rho_0 k}_{(n\Omega - kv_0)^2} \quad v_{ind}(t,r) = -\sum_{n} \underbrace{A_n}_{n\Omega - kv_0} \{\sin(kr - n\Omega t) - \sin(kr - kv_0 t)\} \\ \times \{\sin(kr - n\Omega t) - \sin(kr - kv_0 t) + (n\Omega - kv_0)t\cos(kr - kv_0 t)\}$$
Resonance condition:

$$n\Omega \sim kv_0$$

$$\underbrace{I'Hôpital's rule}_{n\Omega \to kv_0} \rho_{ind,n} = \frac{A_n \rho_0 k}{2} t^2 \sin[k(r - v_0 t)] \quad \lim_{n\Omega \to kv_0} v_{ind,n} = A_n t\cos[k(r - v_0 t)]$$

Prediction of Core Scale

Resonance condition

$$T \approx t_d(r_{\text{core}}) = \sqrt{\frac{3\pi}{32G\overline{\rho}(r_{\text{core}})}}$$

Resonance occurs when the condition is satisfied

- ⇒ Efficient energy transfer
- ⇒ System expands
- ⇒ Cusp-to-Core transformation

<u>1. Mass profile of CDM halos</u>

$$\rho(r) = \frac{\rho_0 R_{DM}^3}{r^{\alpha} (r + R_{DM})^{3 - \alpha}} \ \text{(NFW: α=1)}$$

$$T_{\rm c}^2 \equiv \frac{\pi^2}{8G} \frac{R_{\rm DM}^3 c^{3-\alpha} {}_2 F_1[\alpha;-c]}{M_{\rm vir}} \quad \begin{array}{l} \mbox{\it F: Gauss's hyper-} \\ \mbox{geometric function} \end{array}$$

$$r_{\rm core} = R_{DM} \left(\frac{T}{T_c}\right)^{2/\alpha}$$

2. Inversion procedure

$$r_{\rm core} = t_{\rm d}^{-1}(T)$$



Numerical Model



DM halo (*N*-body system): NFW model (Navarro et al. 1997)

Tree code developed for GPU clusters (GO+ 2013)

Baryon (external potential): $\Phi_b(r,$ Hernquist potential (Hernquist 1990)

$$t) = -\frac{GM_b}{r + R_b(t)}$$
$$R_b(t) \propto \cos\left(2\pi t/T\right)$$

Property of DM halo

Number of particles N	16M, 128M
Softening parameter ϵ	0.004kpc
Virial mass of a DM halo $M_{ m vir}$	$10^9 M_{\odot}$
Virial radius of a DM halo $R_{\rm vir}$	10kpc
Scale radius of a DM halo $R_{\rm DM}$	2kpc
Total baryon mass $M_{ m b,tot}$	$1.7 \times 10^{8} M_{\odot}$

Oscillation period of the external force, T $T = 1,3,10 \tau$ $\tau \equiv 10 \text{Myr} \sim t_{d}(0.2 \text{kpc})$

Density Profile

The cusp-to-core transformation and resultant core scale depend on the oscillation period of the external force, T.



Fourier Spectrum of Radial Velocity



The Connection between the Cusp—to—Core Transformation and Observational Universalities of DM Halos Ogiya et al., accepted for publication in MNRAS Letters

Universalities of DM Halos



Assumption & Procedure



1. Halos form following an NFW profile with the $c(M_{200}, z)$ for given M_{200} and $z \rightarrow \rho_{\rm S}, r_{\rm S}$ <u>NFW profile</u> $\rho(r) = \frac{\rho_{\rm s} r_{\rm s}^3}{r(r+r_{\rm s})^2}$

2. NFW halos are transformed into Burkert halos

 $\rightarrow \rho_0, r_0$



Generation of the μ_{0D} Relation



Generation of the μ_{0D} Relation

GO, Mori, Ishiyama & Burkert, accepted



Summary

- Gas mass-loss is inefficient to flatten cusps.
- Resonances between DM particles and density waves of galactic gas plays a crucial role to flatten out the central cusp.
- Cusp-to-Core transformation naturally reproduces the μ_{0D} relation.

Estimation of transformation redshift